Decoherence and einselection in equilibrium in an adapted Caldeira Leggett model

Andreas Albrecht

Center for Quantum Mathematics and Physics (QMAP)

and Department of Physics UC Davis



Seminar University of Nottingham

April 4, 2019

Supported by the US Department of Energy

- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)

5. Conclusions

- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)
- 5. Conclusions

With A. Arrasmith

4/4/19

- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)

5. Conclusions

Comments on the arrow of time in cosmology (at board)

Q: What features of the universe are correlated with classicality?

- Arrow of time?
- Locality?
- Etc.

Q: What features of the universe are correlated with classicality?

- Arrow of time?
- Locality?
- Etc.

→ Explore the process of einselection in a toy model, relate to AoT, etc.



- Arrow of time?
- Locality?

Related to the emergence of

classical

Explore the process of einselection in a toy model, relate to AoT, etc.



- Arrow of time?
- Locality?

Related to the emergence of classical Traditionally connected with arrow of time

Explore the process of einselection in a toy model, relate to AoT, etc.

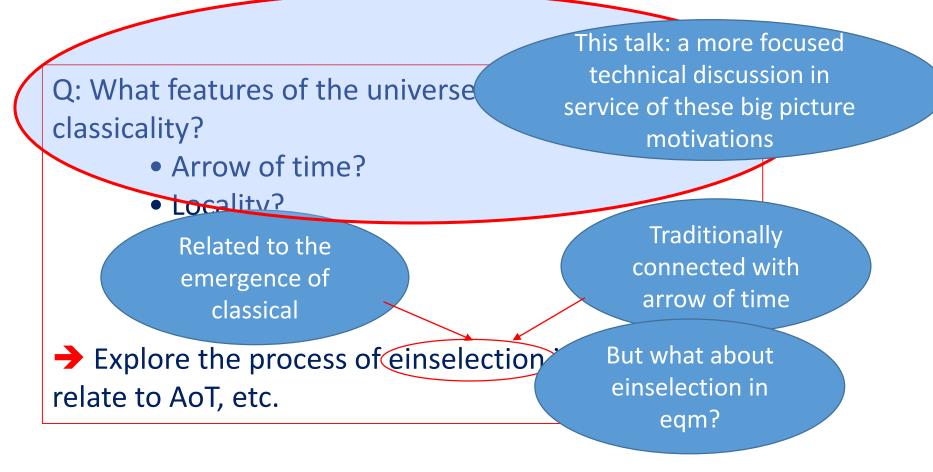
Q: What features of the universe are correlated with classicality?

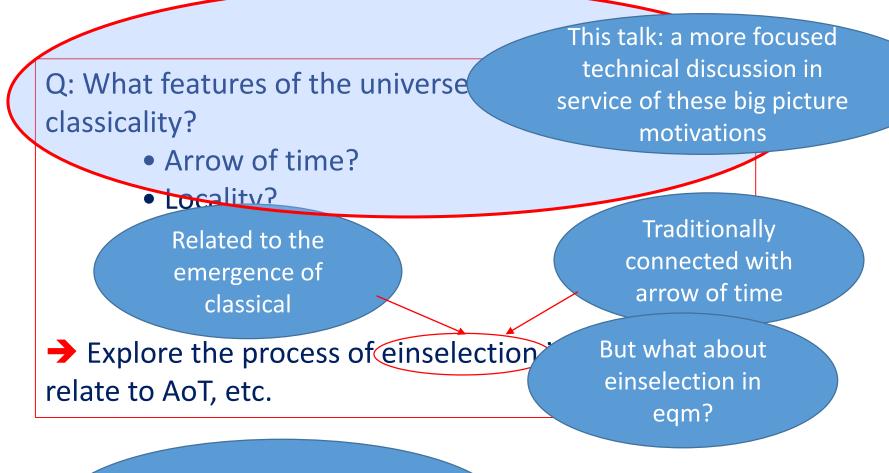
- Arrow of time?
- Locality?

Related to the emergence of classical

Explore the process of einselection relate to AoT, etc. Traditionally connected with arrow of time

But what about einselection in eqm?





If you are handed a theory, what are the classical degrees of freedom?

- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)

5. Conclusions

- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)

5. Conclusions

The preference of special "pointer" states of a system due to interactions with the environment

- Stability of pointer states
- Destruction of non-pointer states (including "Schrödinger cat" superpositions of pointer states)
- Pure non-pointer states → mixtures of pointer states via entanglement with the environment.

The preference of special "pointer" states of a system due to interactions with the environment

- Stability of pointer states
- Destruction of non-pointer states (including "Schrödinger cat" superpositions of pointer states)
- Pure non-pointer states → mixtures of pointer states via entanglement with the environment.

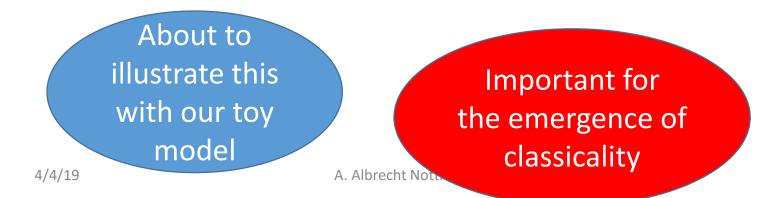
The preference of special "pointer" states of a system due to interactions with the environment

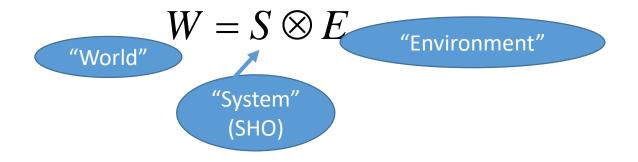
- Stability of pointer states
- Destruction of non-pointer states (including "Schrödinger cat" superpositions of pointer states)
- Pure non-pointer states → mixtures of pointer states via entanglement with the environment.



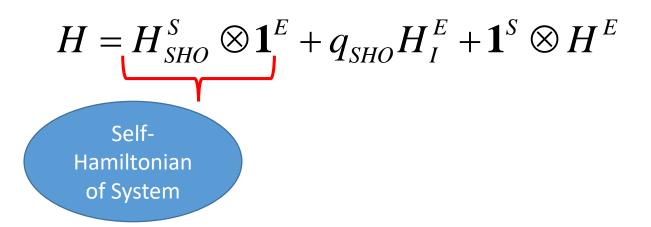
The preference of special "pointer" states of a system due to interactions with the environment

- Stability of pointer states
- Destruction of non-pointer states (including "Schrödinger cat" superpositions of pointer states)
- Pure non-pointer states → mixtures of pointer states via entanglement with the environment.

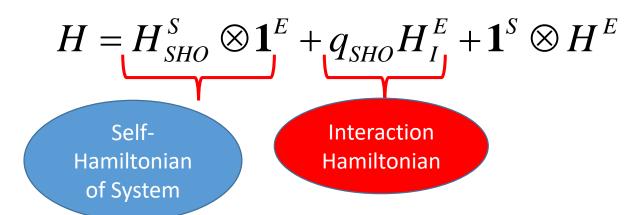




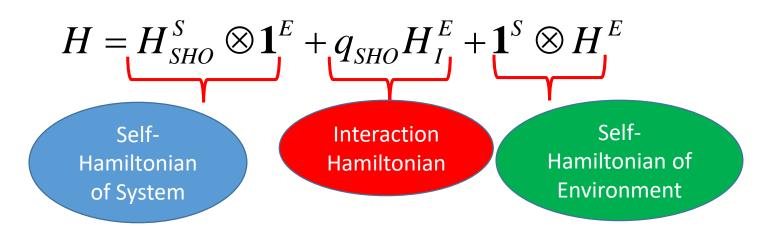
$$W = S \otimes E$$



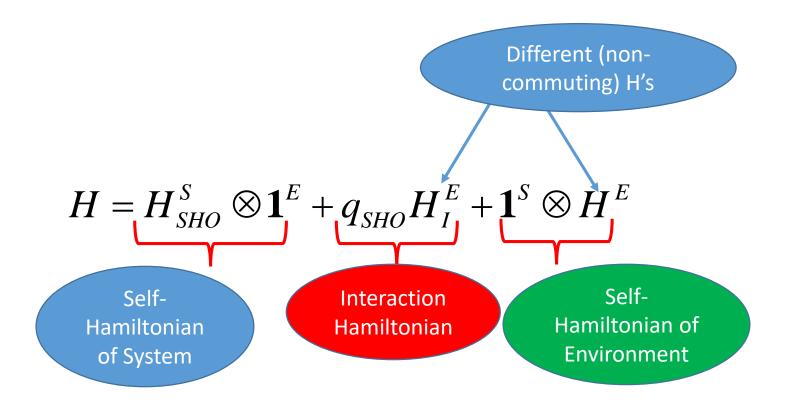
$$W = S \otimes E$$

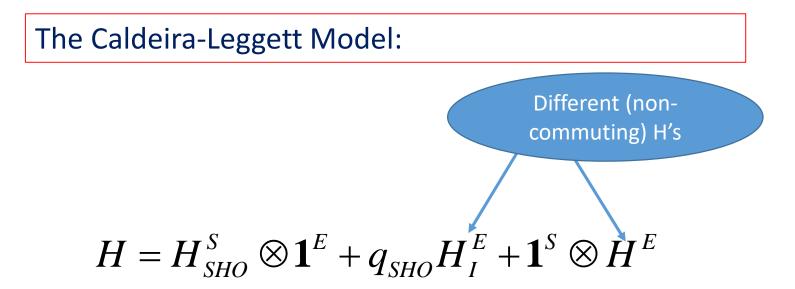


$$W = S \otimes E$$



$$W = S \otimes E$$





<u>CL:</u>

- i) Model *E* as an infinite set of SHOs with different frequencies
- ii) Take special (order of) limits and parameter choices to get an <u>(irreversible) stochastic equation</u> that describes this (unitary) evolution under certain conditions (including AoT)

iii) Demonstrate einselection etc. (CL and others)

The Caldeira-Leggett Model:

$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I_{s}}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

$$\frac{\text{Adapted CL:}}{\text{i) Model } E \text{ as finite system}}$$

$$\text{ii) Solve full unitary evolution in all regimes (numerical)}$$

- iii) Demonstrate Einselection under certain conditions (AoT)
- iv) Explore scope of einselection (eqm?)

Introducing the toy model

- No interaction case $(H_I^E = 0)$
- Model SHO with d=30 Hilbert space

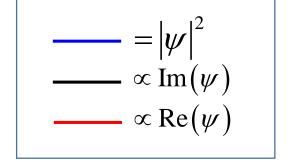
$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$

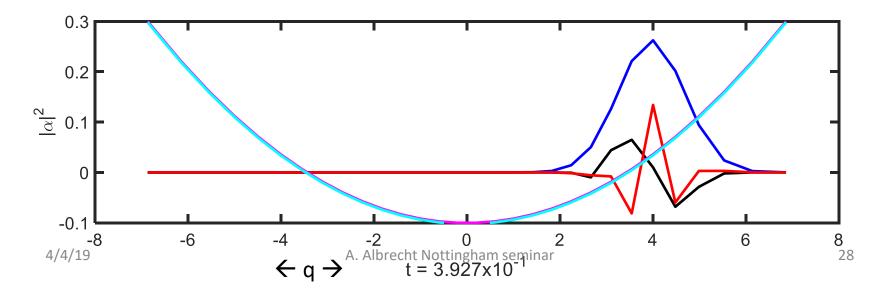
Introducing the toy model

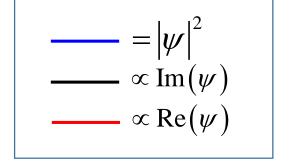
- No interaction case $(H_I^E = 0)$
- Model SHO with d=30 Hilbert space

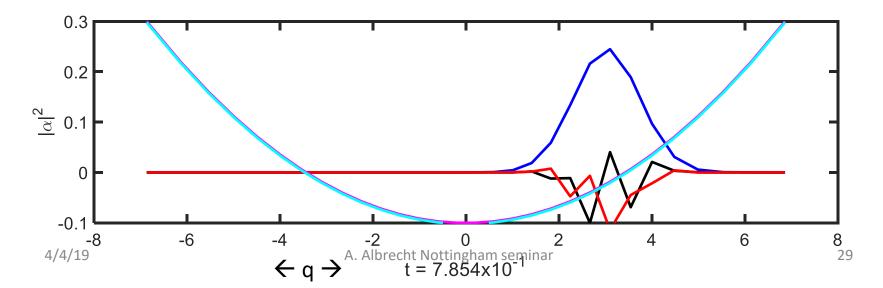
$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

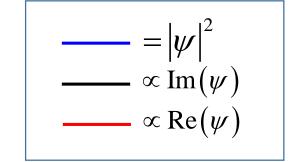
"Movie" A (Isolated SHO)

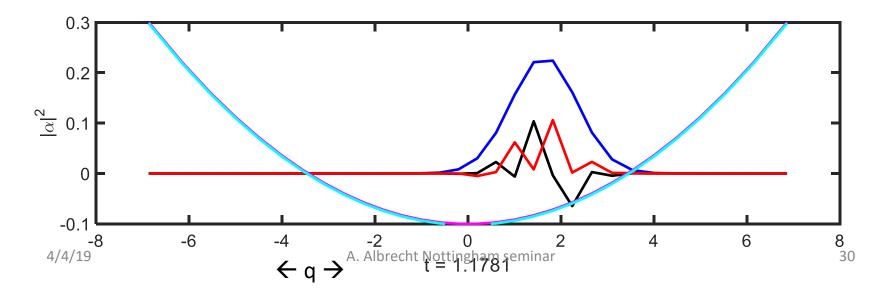


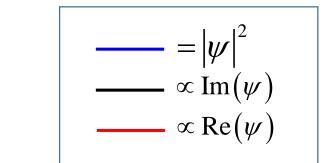


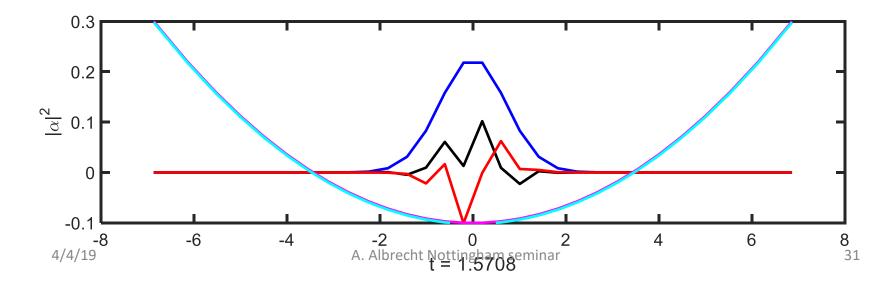


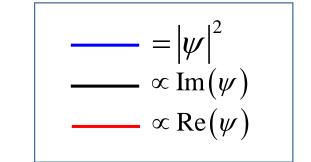


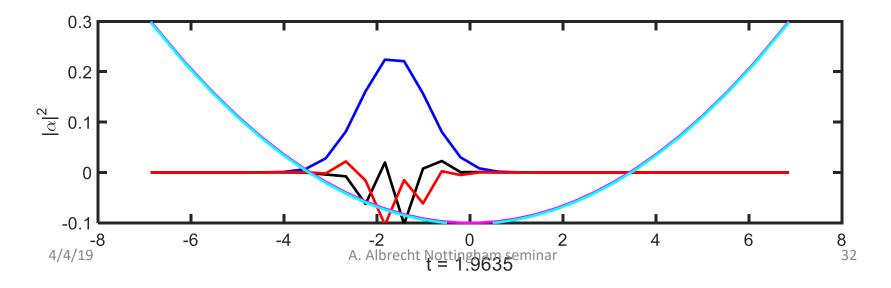


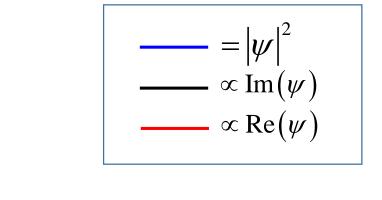


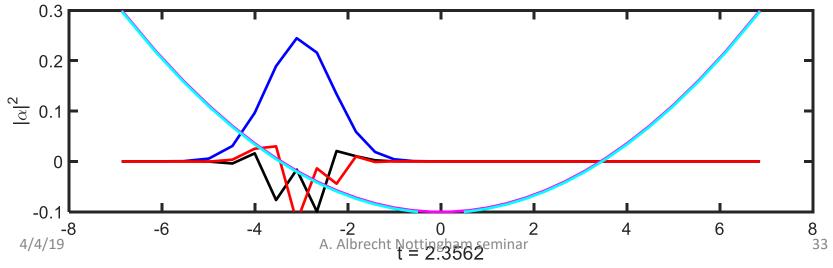


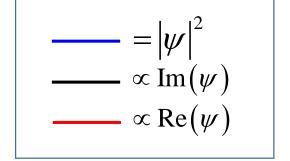


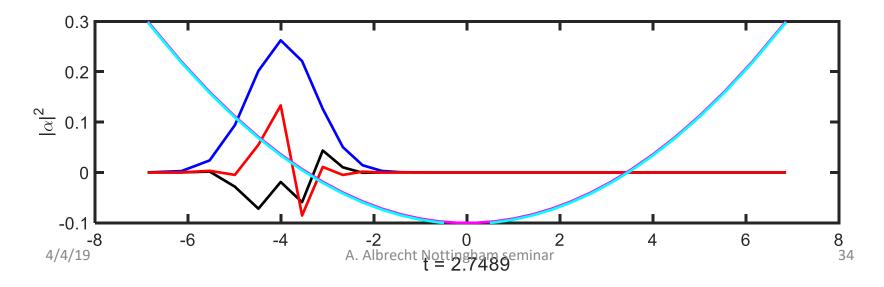


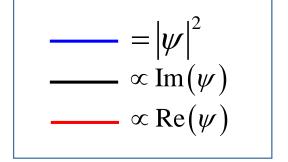


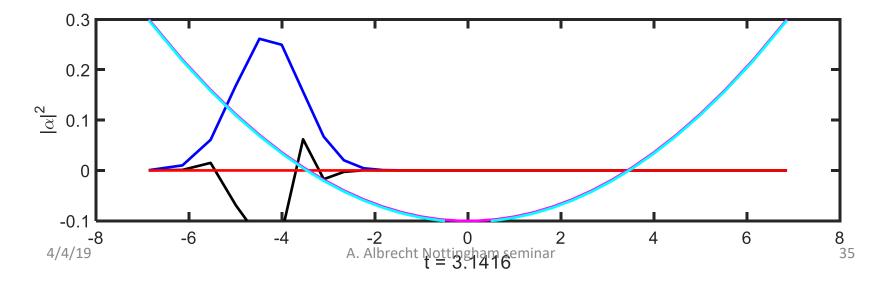


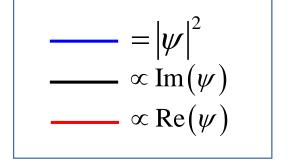


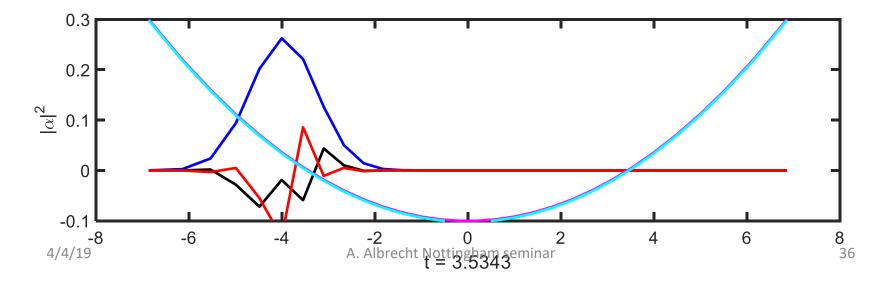


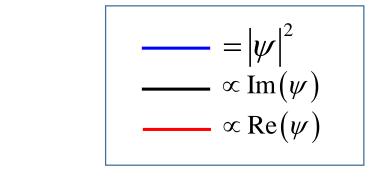


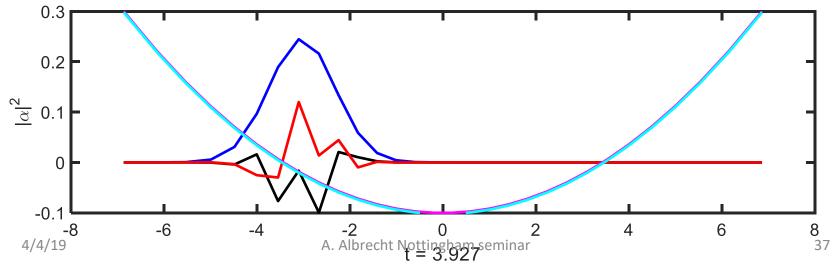


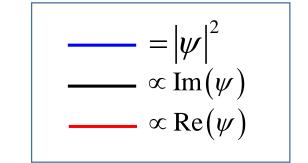


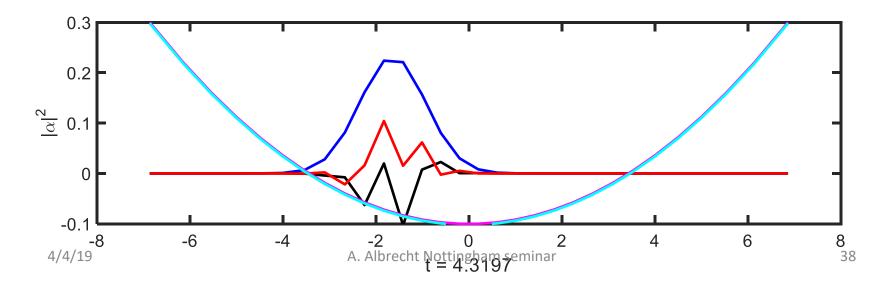


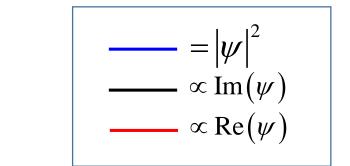


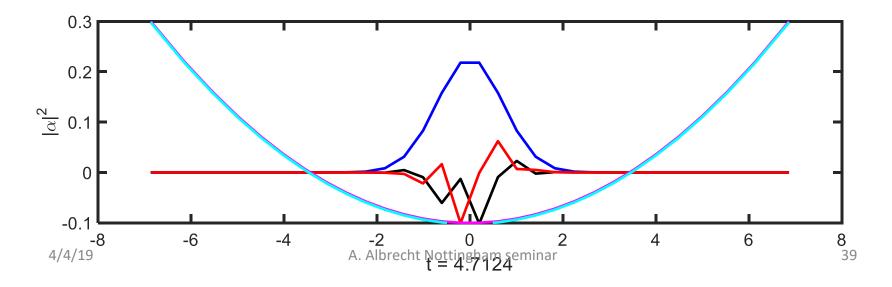


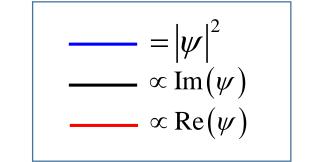


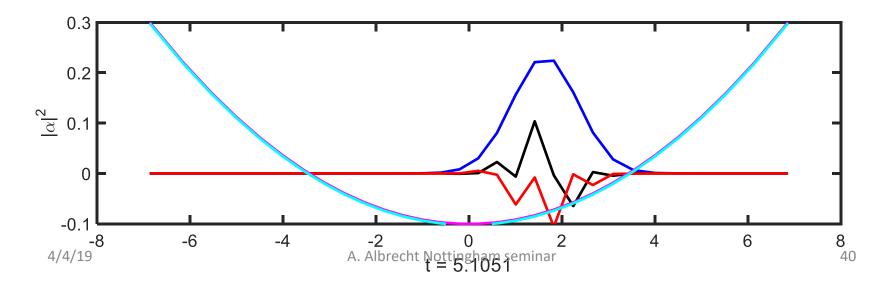


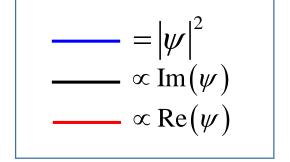


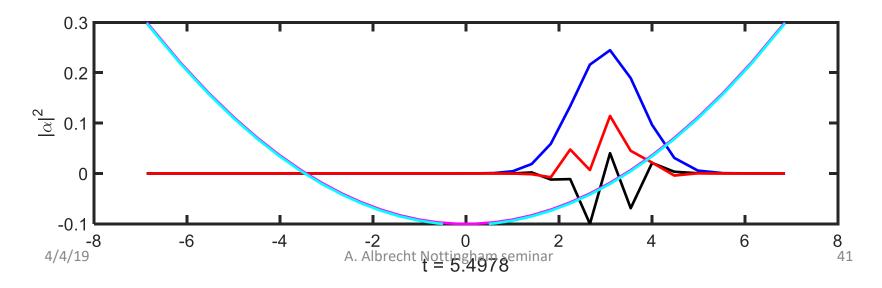


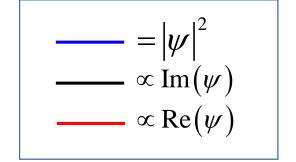


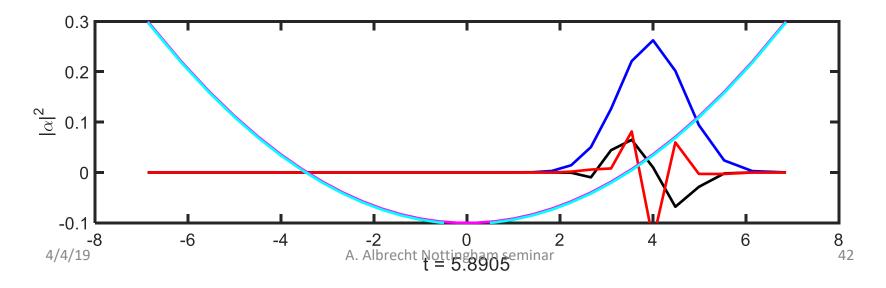


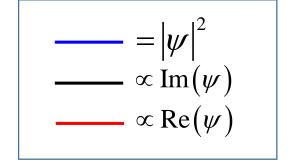


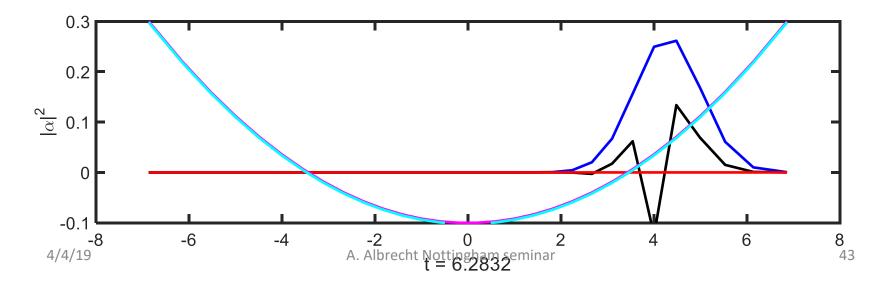






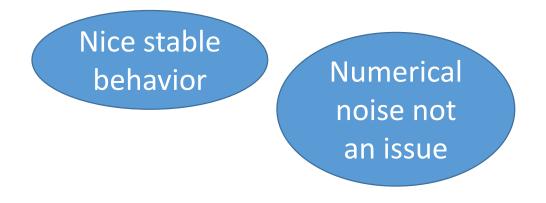






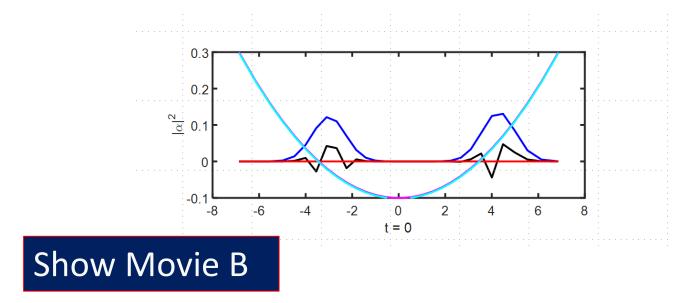
Introducing the toy model

- No interaction case $(H_I^E = \mathbf{0})$
- Model SHO with d=30 Hilbert space

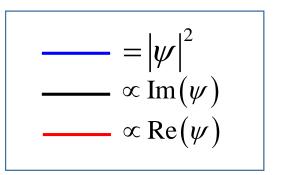


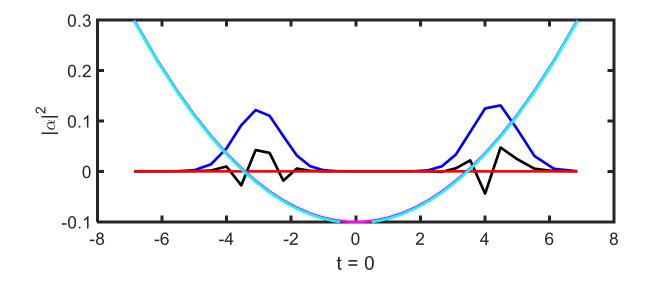
Introducing the <u>Schrödinger cat</u>

- No interaction case $(H_I^E = \mathbf{0})$
- Model SHO with d=30 Hilbert space

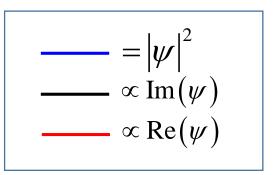


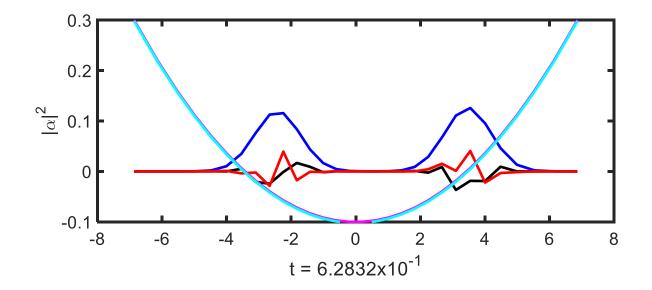




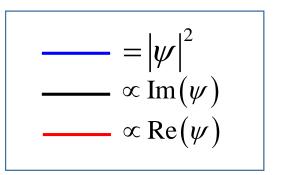


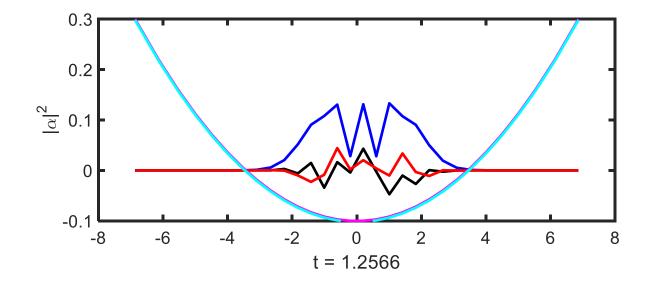




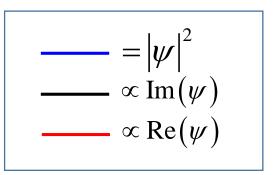


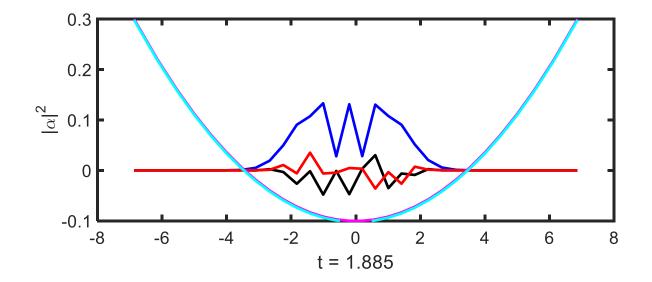




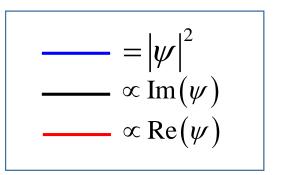


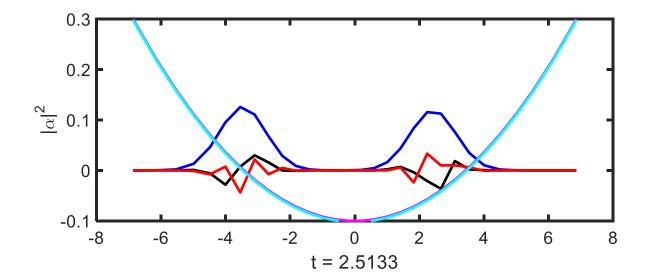




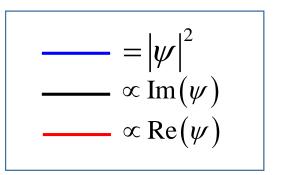


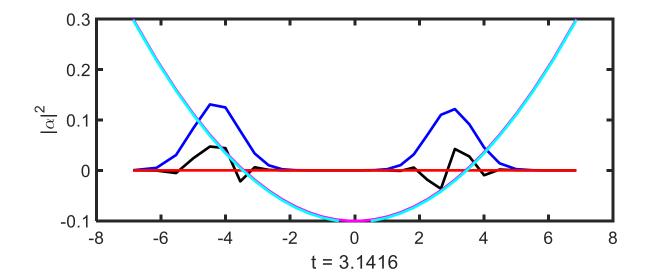




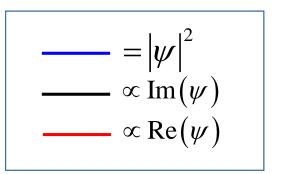


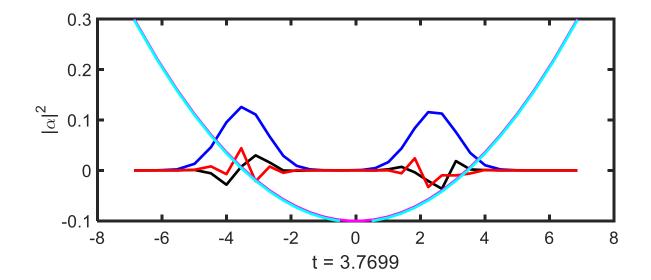




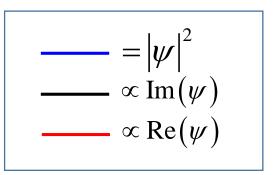


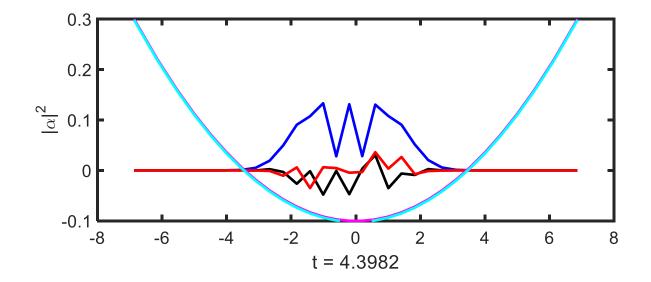




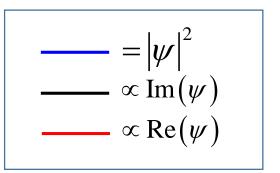


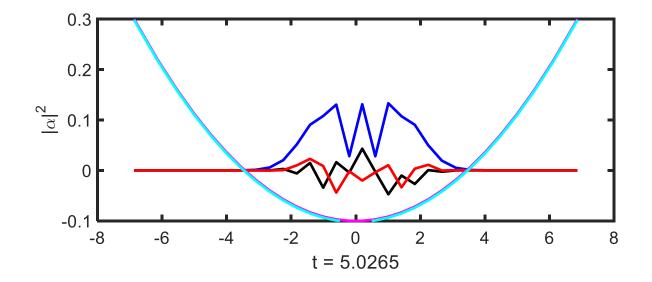




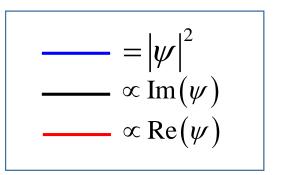


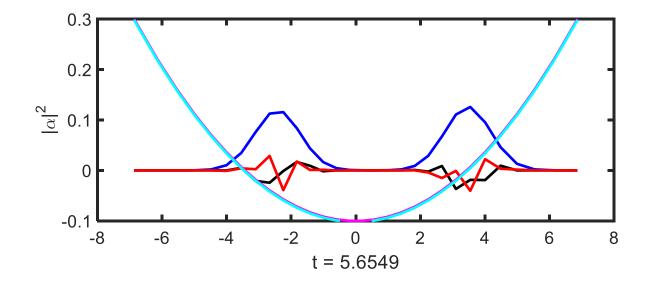




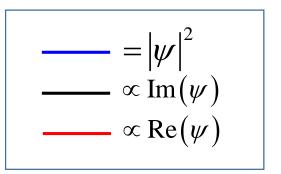


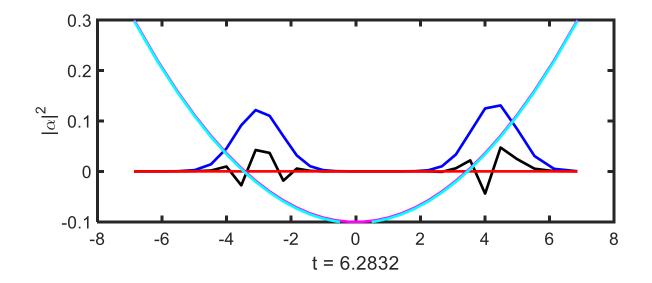












• Interactions turned on $(H_I^E \neq \mathbf{0})$

$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

This evolution will take an initial product state into a mixed state:

$$|\psi\rangle_{W} = |\psi\rangle_{S} |\psi\rangle_{E} \rightarrow \sum_{i,j} \alpha_{ij} |i\rangle_{S} |j\rangle_{E}$$

• Interactions turned on $(H_I^E \neq \mathbf{0})$

$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

This evolution will take an initial product state into a mixed state:

$$|\psi\rangle_{W} = |\psi\rangle_{S} |\psi\rangle_{E} \Longrightarrow \sum_{i,j} \alpha_{ij} |i\rangle_{S} |j\rangle_{E}$$

• Interactions turned on $(H_I^E \neq \mathbf{0})$

$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

This evolution will take an initial product state into a

ixed state:
Pure
(product)

$$|\psi\rangle_{W} = |\psi\rangle_{S} |\psi\rangle_{E} \longrightarrow \sum_{i,j} \alpha_{ij} |i\rangle_{S} |j\rangle_{E}$$

m

• Interactions turned on $(H_I^E \neq \mathbf{0})$

$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

This evolution will take an initial product state into a

mixed state:
Pure
(product)

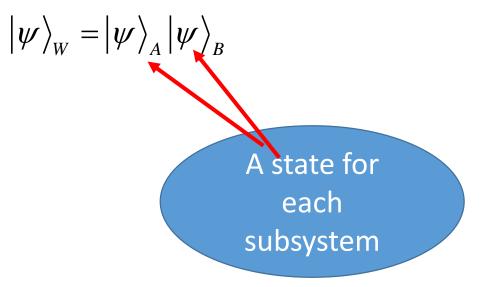
$$|\psi\rangle_{W} = |\psi\rangle_{S} |\psi\rangle_{E} \longrightarrow \sum_{i,j} \alpha_{ij} |i\rangle_{S} |j\rangle_{E}$$

$$\rho_{S} \equiv Tr_{E} \left(\left| \psi \right\rangle_{WW} \left\langle \psi \right| \right) = \left| \psi \right\rangle_{SS} \left\langle \psi \right| \Longrightarrow \text{more general } \rho_{S}$$

Some comments on entangled states

$$W = A \otimes B$$

Inclined to think:



$$W = A \otimes B$$

But the general case is:

$$\left|\psi\right\rangle_{W}=\sum_{i,j}\alpha_{ij}\left|i\right\rangle_{A}\left|j\right\rangle_{B}$$

$$W = A \otimes B$$

But the general case is:

$$\left|\psi\right\rangle_{W} = \sum_{i,j} \alpha_{ij} \left|i\right\rangle_{A} \left|j\right\rangle_{B}$$

Which gives:

$$\rho_{A} \equiv Tr_{B} \left(|\psi\rangle_{W \ W} \left\langle \psi \right| \right)$$

$$\rho_{B} \equiv Tr_{A} \left(|\psi\rangle_{W \ W} \left\langle \psi \right| \right)$$
A density
matrix for
each
subsystem

$$\langle O_A \rangle \equiv tr(\rho_A O_A) = \sum_i p_i \langle p_i | \hat{O} | p_i \rangle$$

Eigenvalues and eigenstates of ρ_A

$$\langle O_A \rangle \equiv tr(\rho_A O_A) = \sum_i p_i \langle p_i | \hat{O} | p_i \rangle$$

Eigenvalues and eigenstates of ρ_A

("Schmidt states")

$$\langle O_A \rangle \equiv tr(\rho_A O_A) = \sum_i p_i \langle p_i | \hat{O} | p_i \rangle$$

Eigenvalues and eigenstates of ρ_A

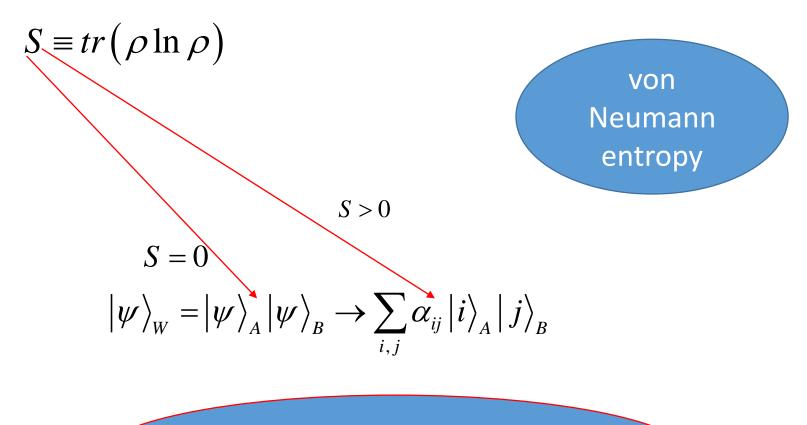
A "classical mixture" of the eigenstates (no cross terms → no "quantum coherence")

$$\langle O_A \rangle \equiv tr(\rho_A O_A) = \sum_i p_i \langle p_i | \hat{O} | p_i \rangle$$

Eigenvalues and eigenstates of ρ_A

A "classical mixture" of the eigenstates (no cross terms → no "quantum coherence")

Onset of entanglement → decoherence



Decoherence → increasing S → arrow of time In general, the density matrix eigenstates might vary significantly from one moment to the next, producing a "classical mixture of random stuff"

- In general, the density matrix eigenstates might vary significantly from one moment to the next, producing a "classical mixture of random stuff"
- Special case: The nature of the interactions between A and B lead to a reliable preference for specific "pointer" eigenstates. This is <u>Einselection</u>

- In general, the density matrix eigenstates might vary significantly from one moment to the next, producing a "classical mixture of random stuff"
- Special case: The nature of the interactions between A and B lead to a reliable preference for specific "pointer" eigenstates. This is <u>Einselection</u>

Discuss pendulum interacting with the air, leading to localized wave packed pointer states.

• Interactions turned on $(H_I^E \neq \mathbf{0})$

$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

• Interactions turned on $(H_I^E \neq \mathbf{0})$

An illustration of Einselection

$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

• Interactions turned on $(H_I^E \neq \mathbf{0})$

$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

This evolution will take an initial product state into a mixed state:

$$|\psi\rangle_{W} = |\psi\rangle_{S} |\psi\rangle_{E} \rightarrow \sum_{i,j} \alpha_{ij} |i\rangle_{S} |j\rangle_{E}$$

$$\rho_{S} \equiv Tr_{E}(|\psi\rangle_{WW}\langle\psi|) = |\psi\rangle_{SS}\langle\psi| \rightarrow \text{more general } \rho_{S}$$

• Interactions turned on $(H_I^E \neq \mathbf{0})$

$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

This evolution will take an initial product stat mixed state:

Show eigenstates and eigenvalues

$$|\psi\rangle_{W} = |\psi\rangle_{S} |\psi\rangle_{E} \rightarrow \sum_{i,j} \alpha_{ij} |i\rangle_{S} |j\rangle_{E}$$

$$\rho_{S} \equiv Tr_{E} \left(|\psi\rangle_{WW} \langle \psi| \right) = |\psi\rangle_{SS} \langle \psi| \rightarrow \text{more general}$$

 $ho_{\rm S}$

Interactions turned on $(H_I^E \neq \mathbf{0})$

$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

This evolution will take an initial product stat mixed state:

First 2 Show eigenstates and eigenvalues

$$|\psi\rangle_{W} = |\psi\rangle_{S} |\psi\rangle_{E} \rightarrow \sum_{i,j} \alpha_{ij} |i\rangle_{S} |j\rangle_{E}$$

$$\rho_{S} \equiv Tr_{E} \left(\left| \psi \right\rangle_{WW} \left\langle \psi \right| \right) = \left| \psi \right\rangle_{SS} \left\langle \psi \right| \rightarrow \text{more general } \rho_{S}$$

• Interactions turned on $(H_I^E \neq 0)$

Show Movie C

$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

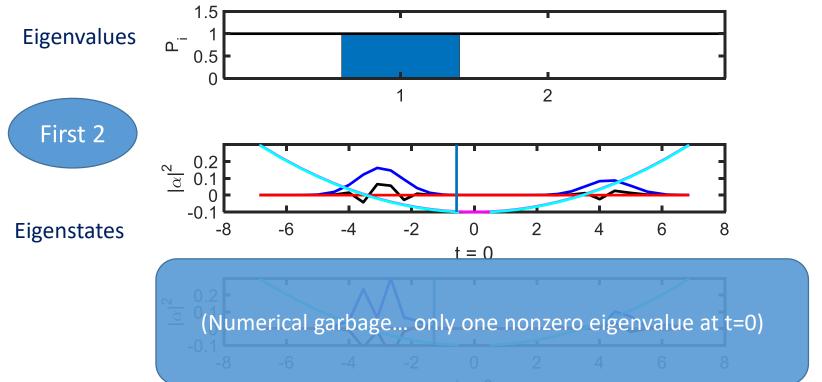
This evolution will take an initial product stat mixed state:

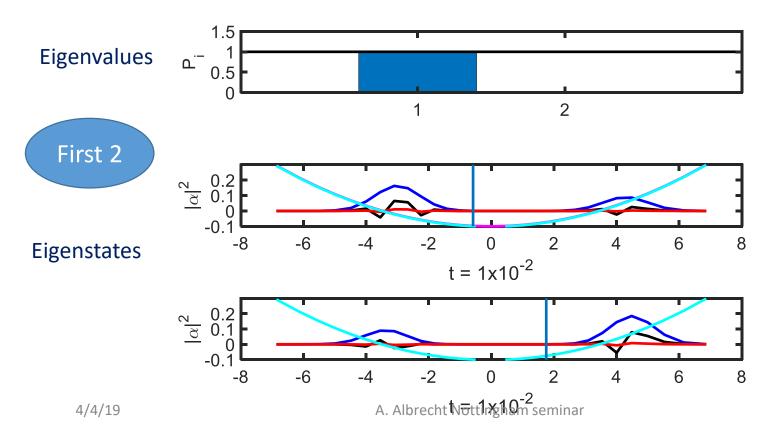
Show First 2 eigenstates and eigenvalues

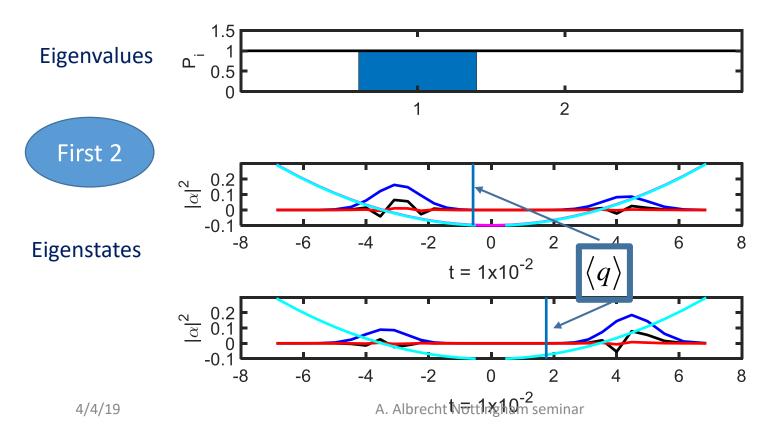
$$|\psi\rangle_{W} = |\psi\rangle_{S} |\psi\rangle_{E} \rightarrow \sum_{i,j} \alpha_{ij} |i\rangle_{S} |j\rangle_{E}$$

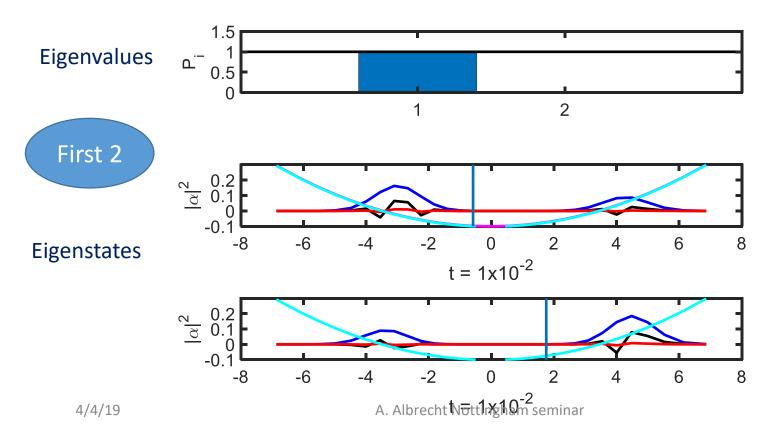
$$\rho_{S} \equiv Tr_{E}(|\psi\rangle_{WW}\langle\psi|) = |\psi\rangle_{SS}\langle\psi| \rightarrow \text{more general } \rho_{S}$$

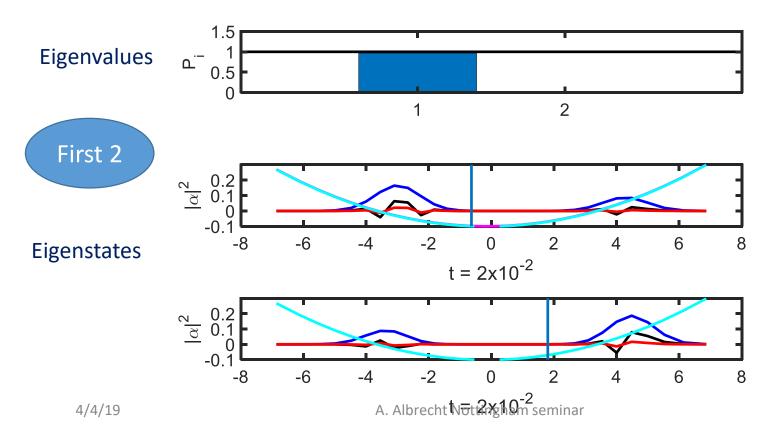


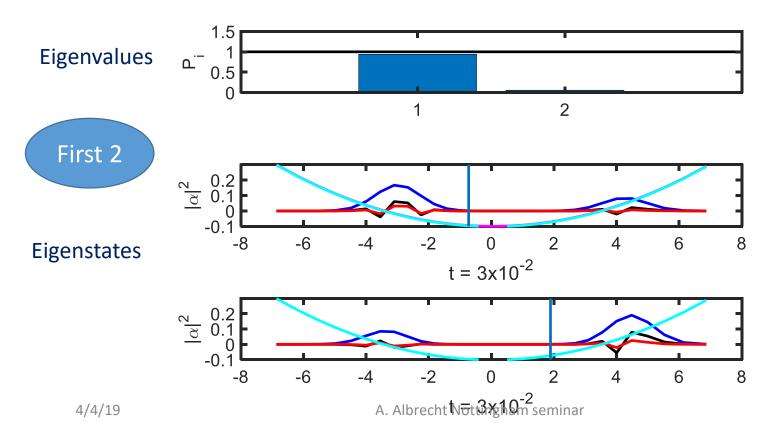


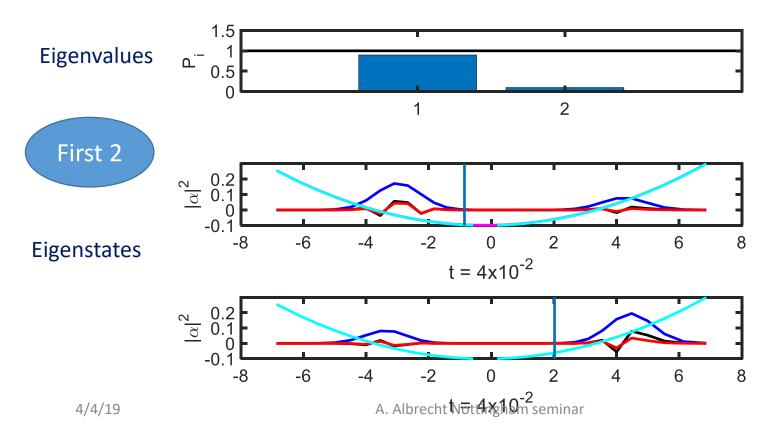


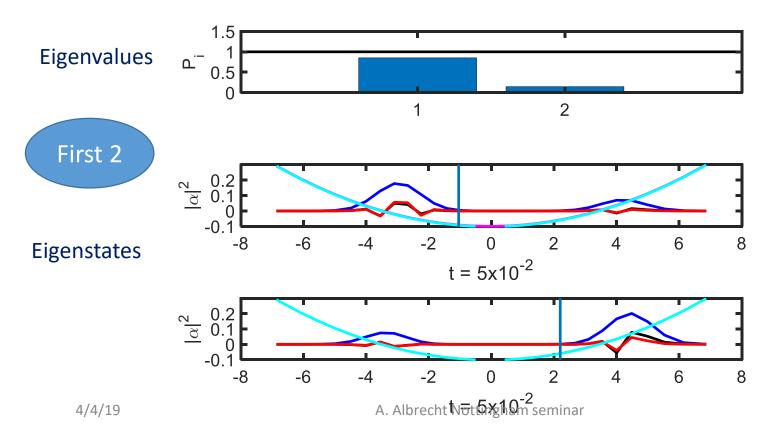


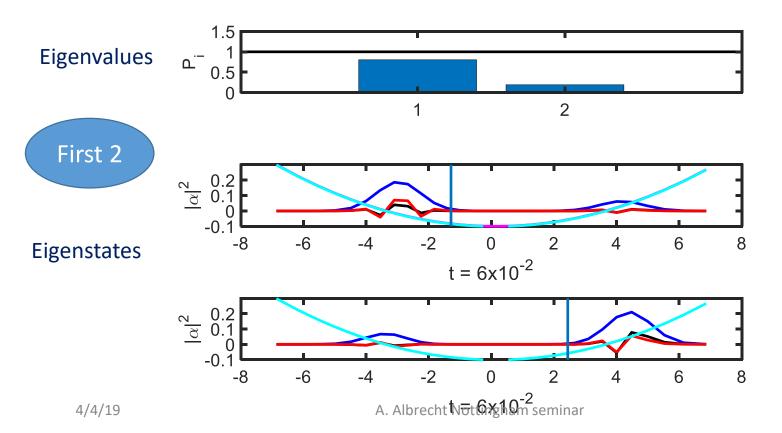


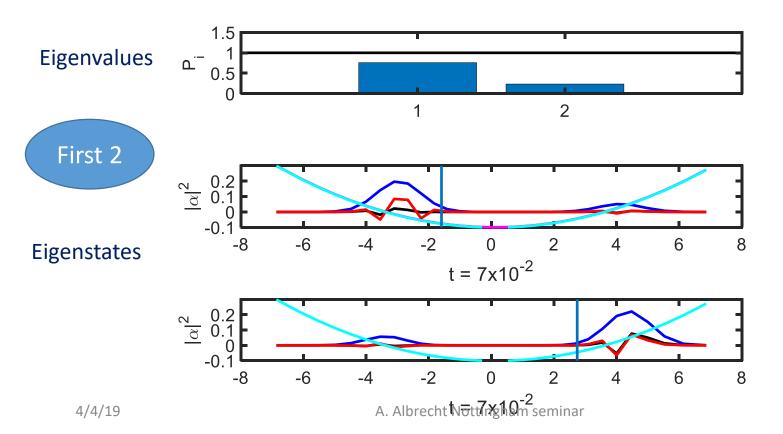


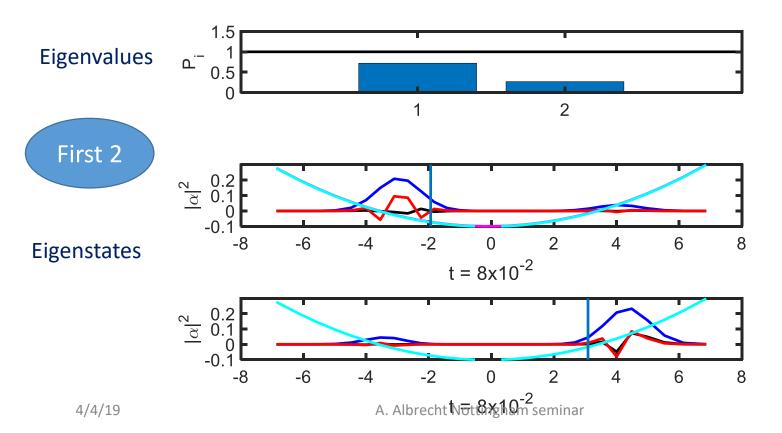


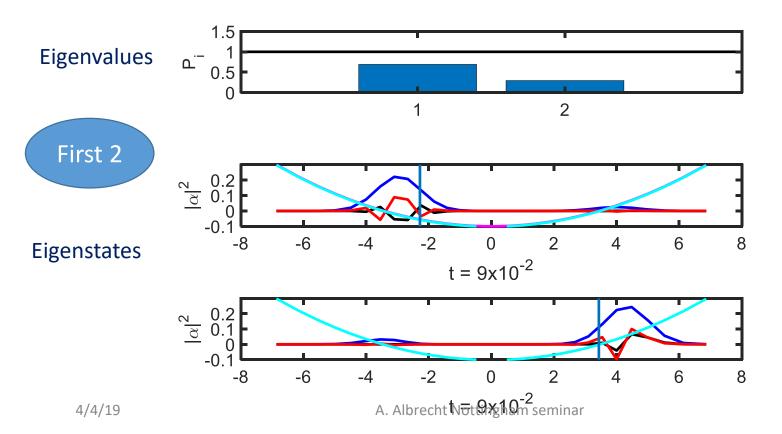


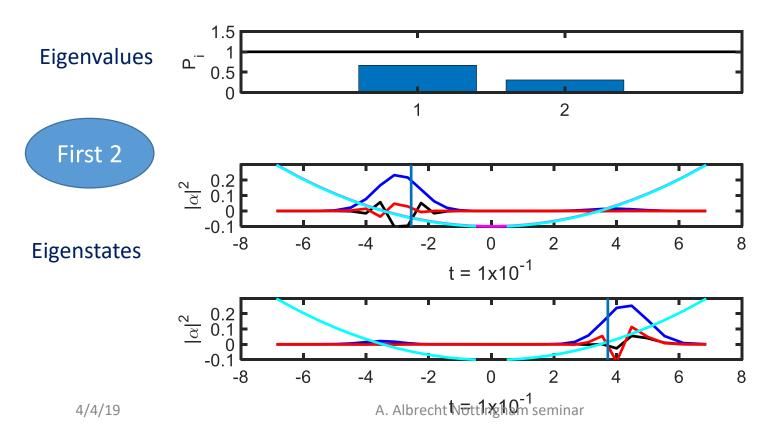


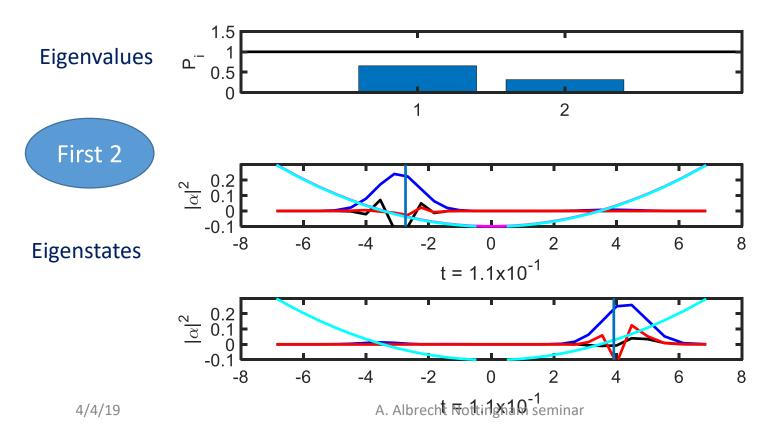


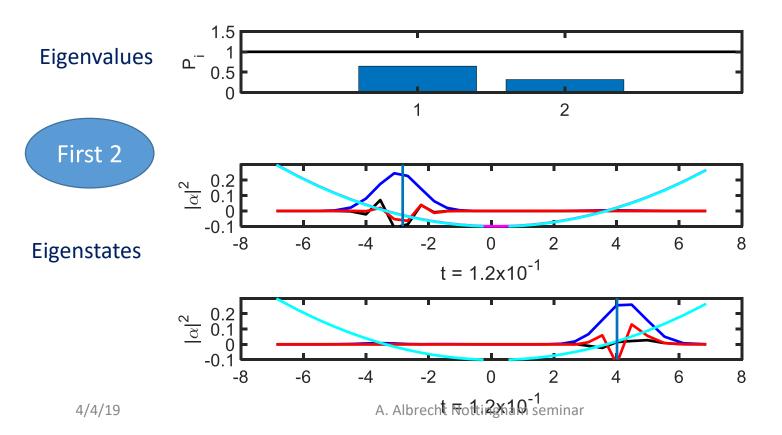


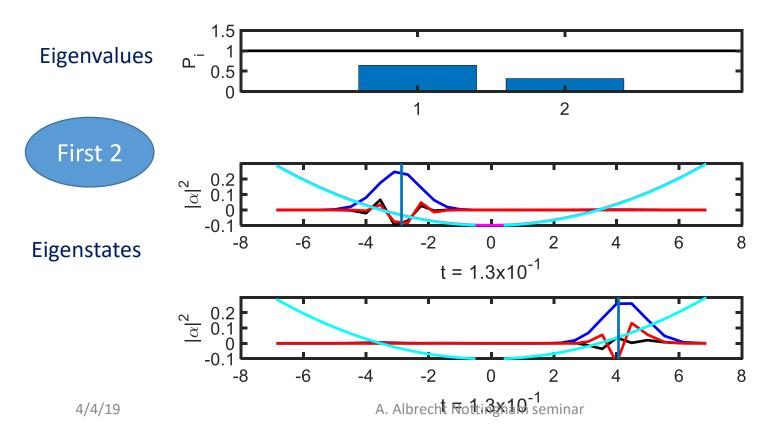


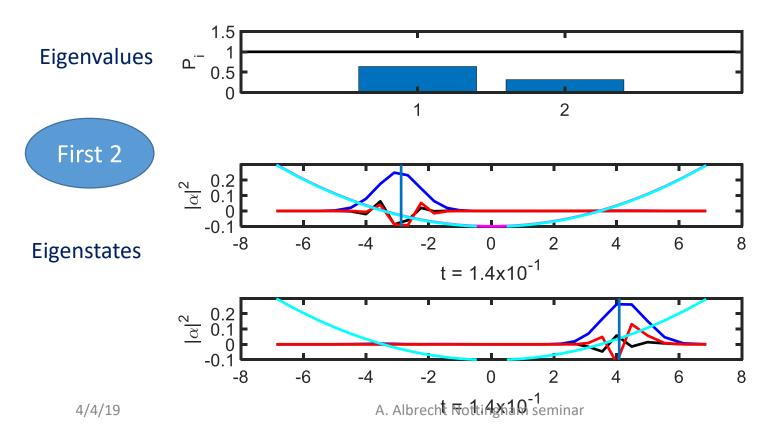




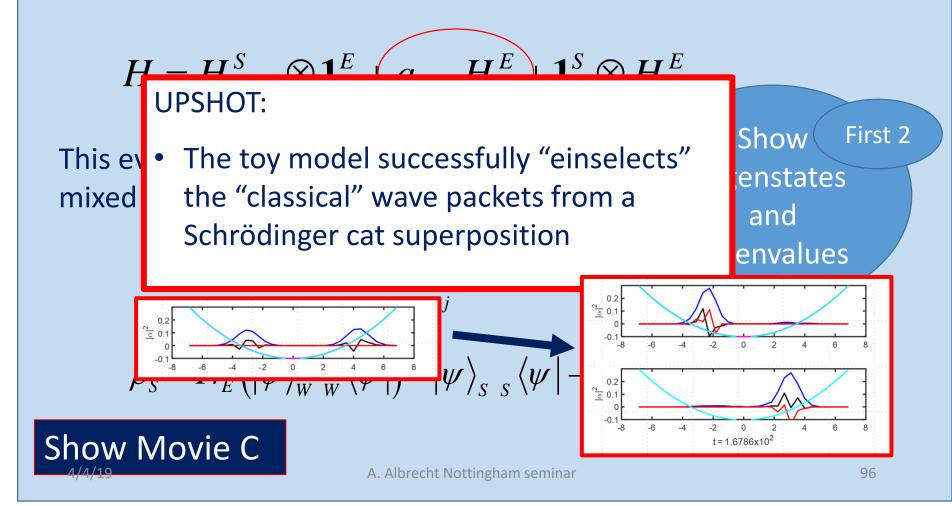








• Interactions turned on $(H_I^E \neq \mathbf{0})$



Outline

- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)

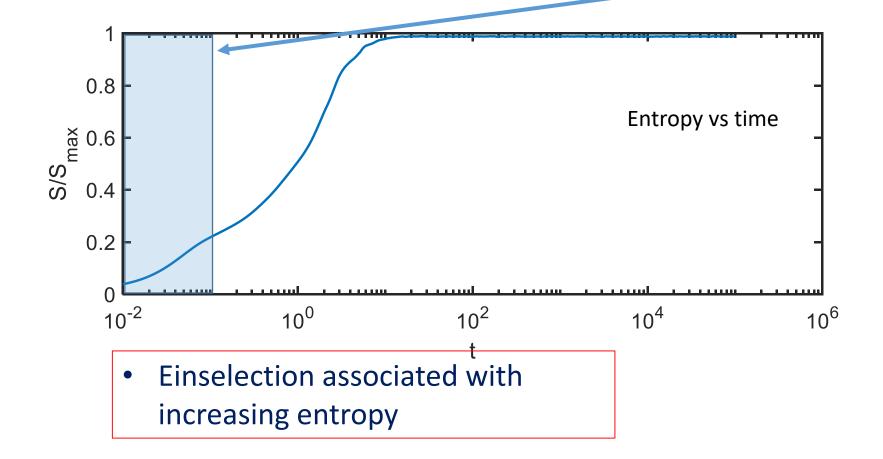
5. Conclusions

Outline

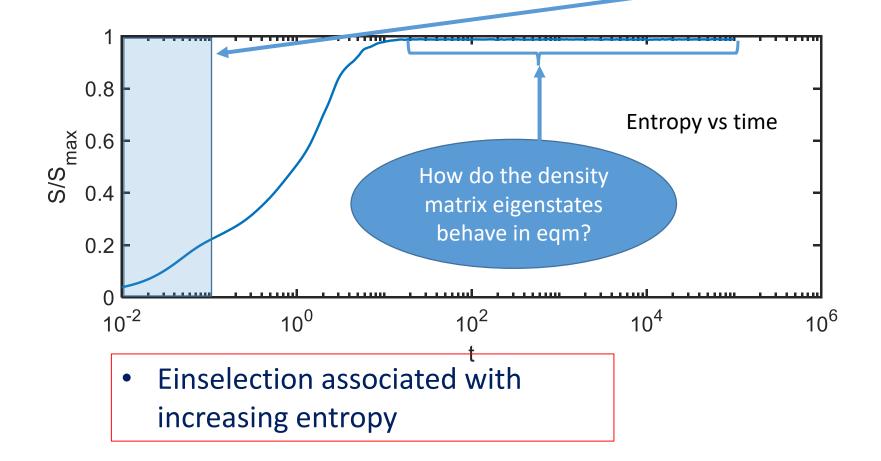
- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)

5. Conclusions

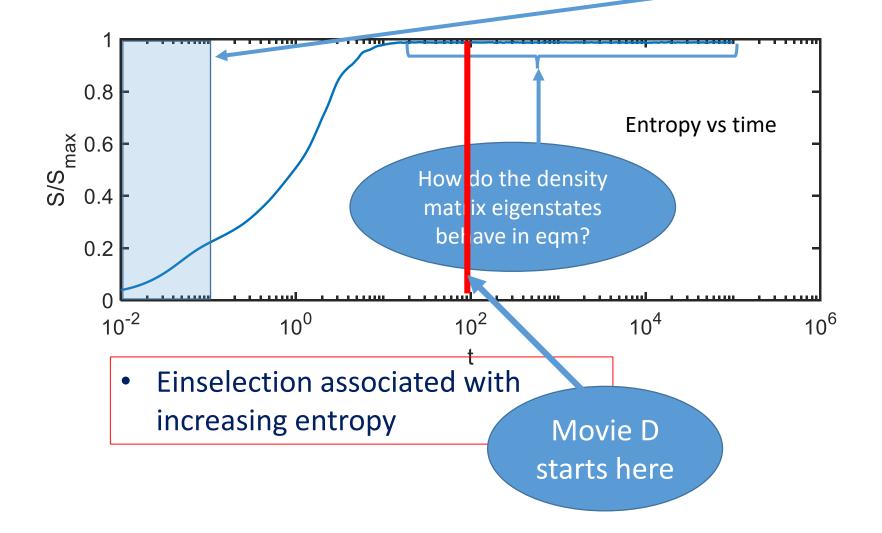
The collapsing Schrödinger cat (Movie C) was in this time window



The collapsing Schrödinger cat (Movie C) was in this time window

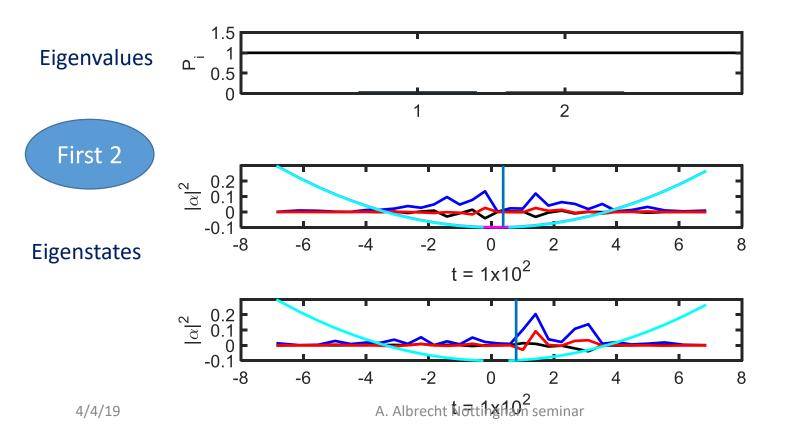


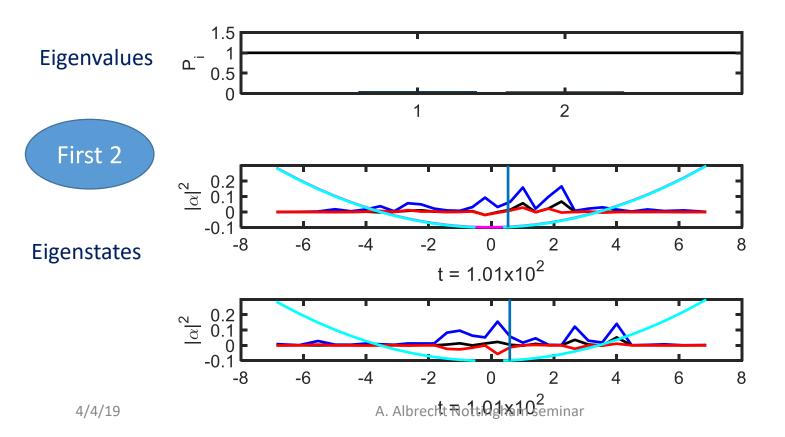
The collapsing Schrödinger cat (Movie C) was in this time window

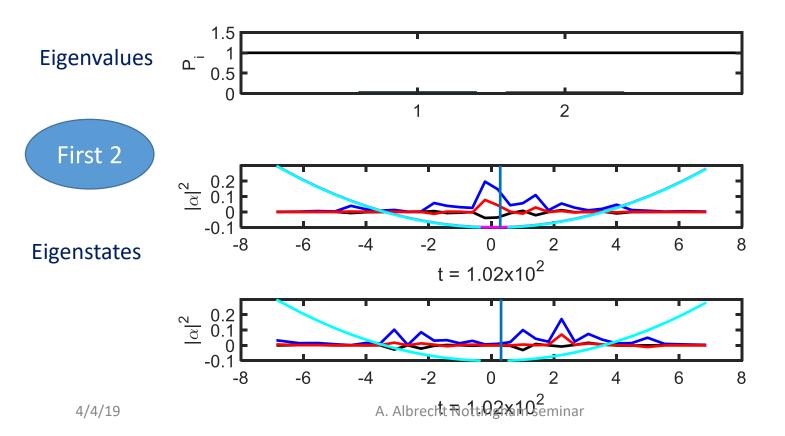


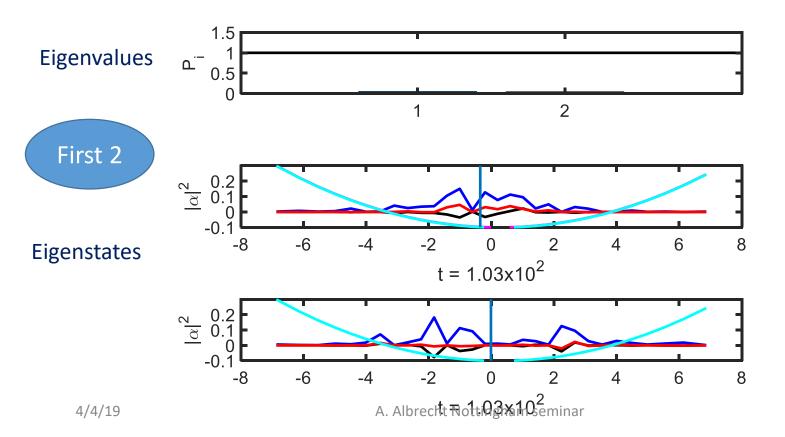


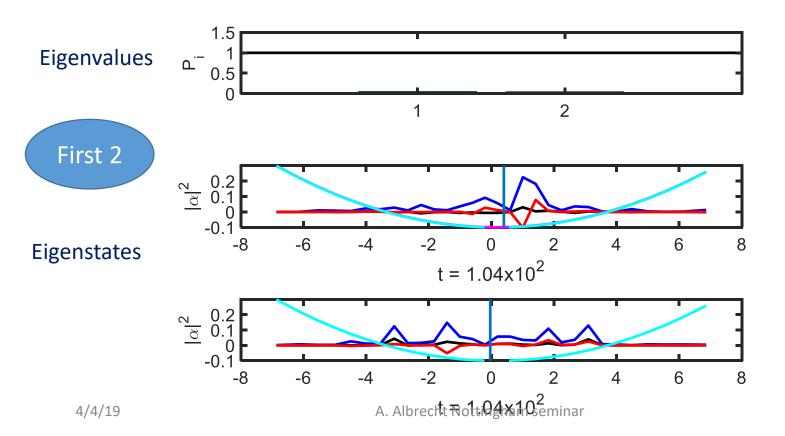
Discuss "detailed balance" (wallowing) of everett worlds at board

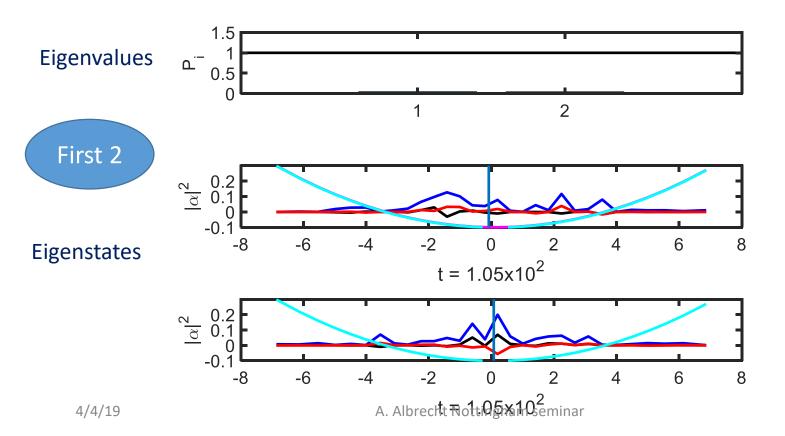


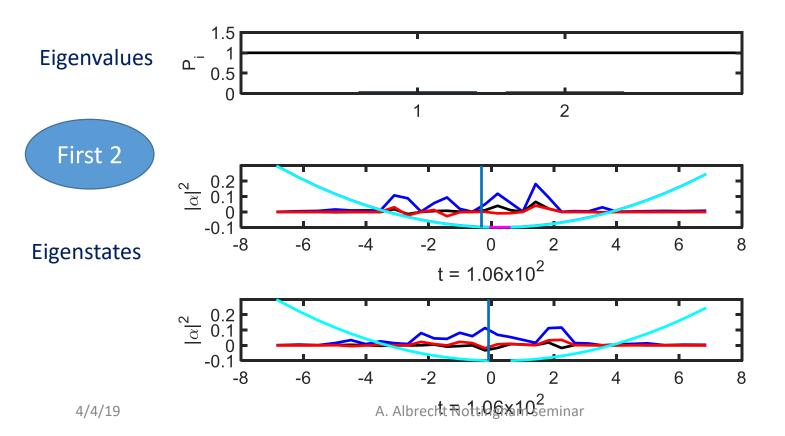


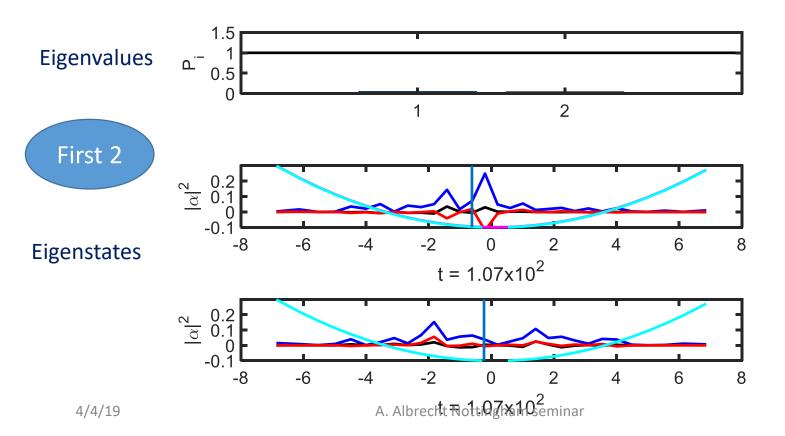


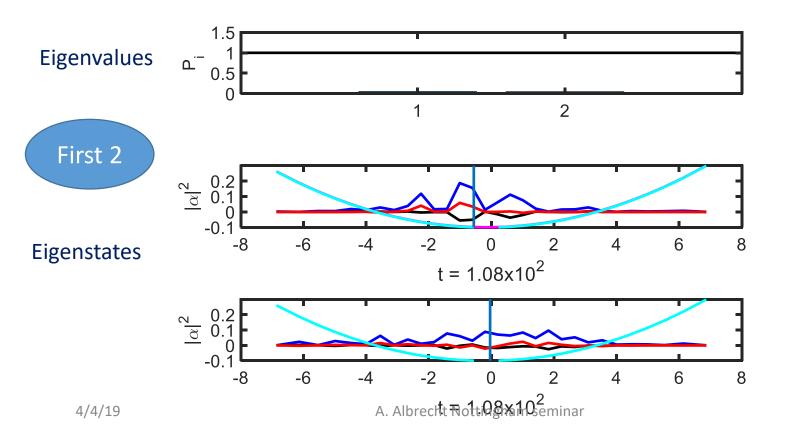


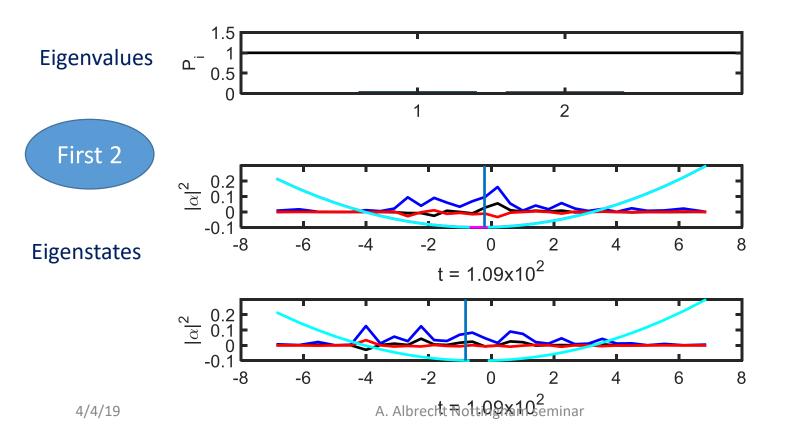


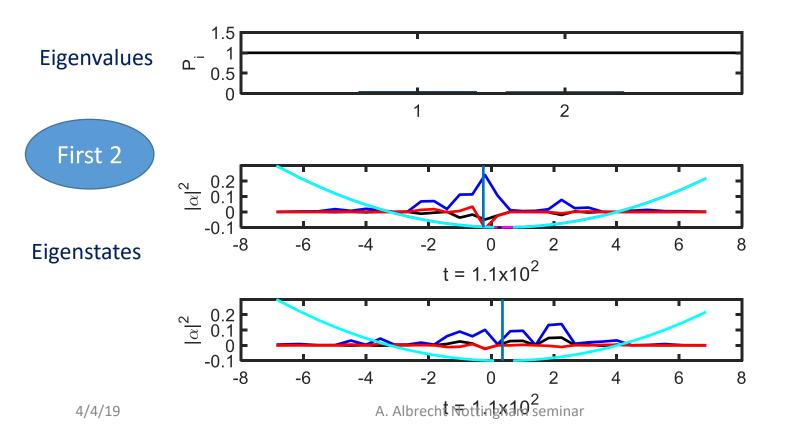


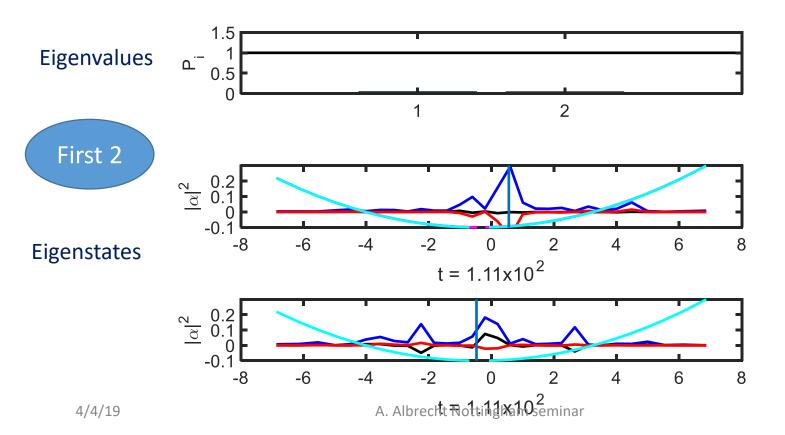


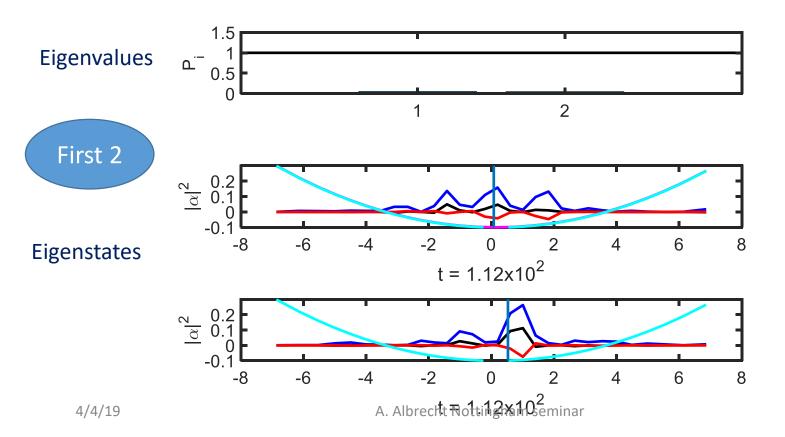


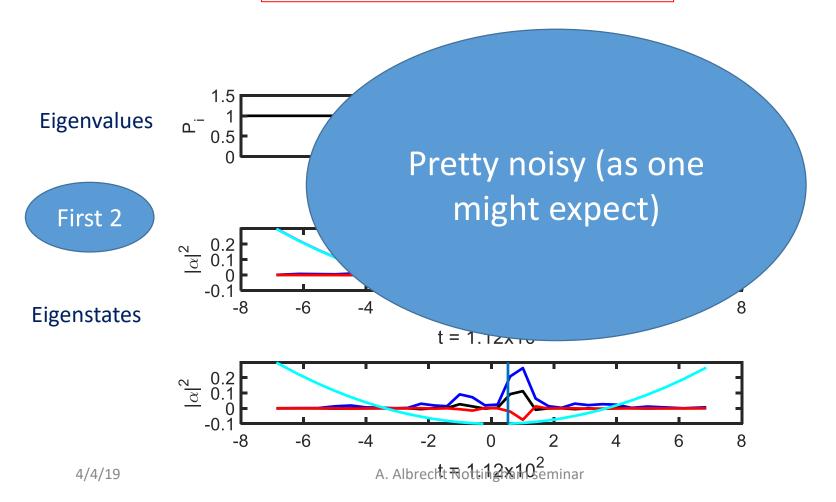


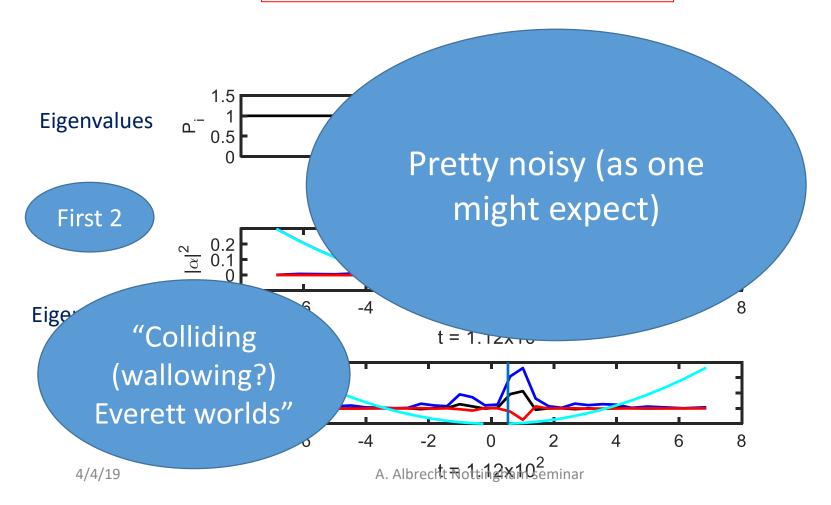












However, when the coupling strength is reduced, things get more interesting

-6

-4

-0.1

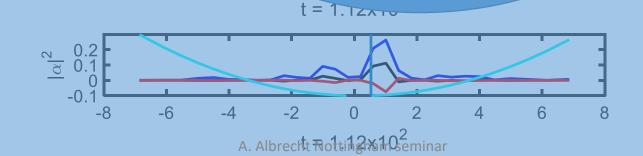
-8

y noisy (as one night expect)

Eigenstates

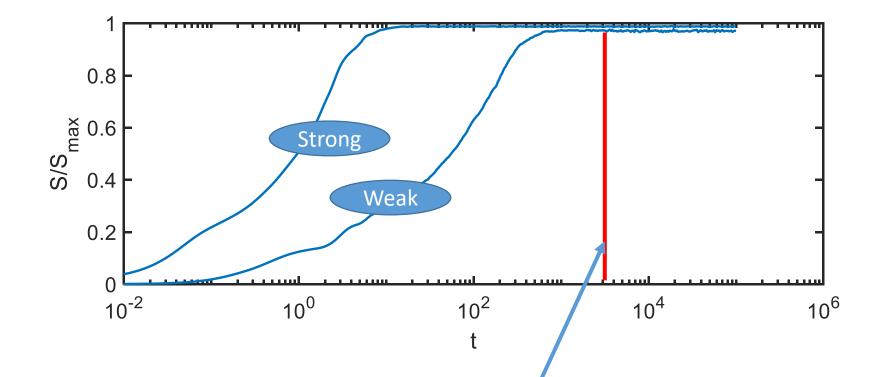
4/4/19

Eigen



118

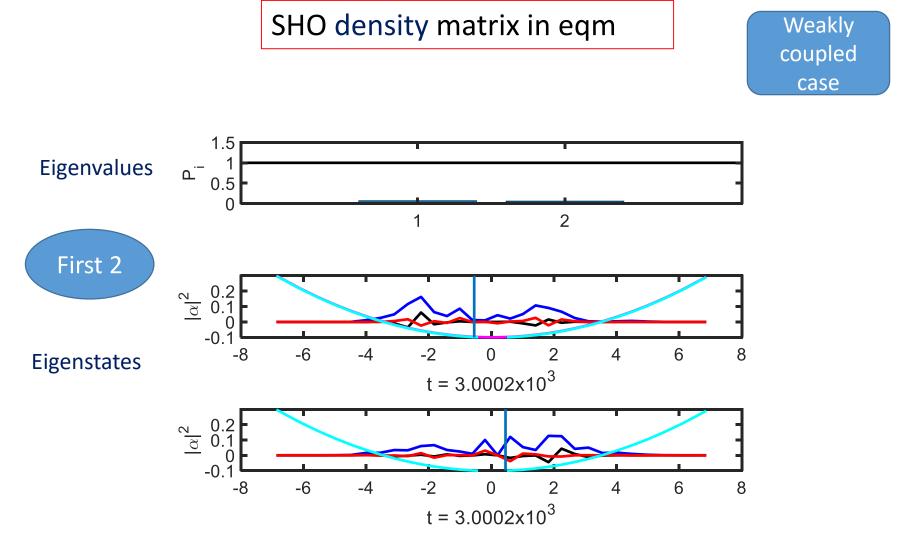
8

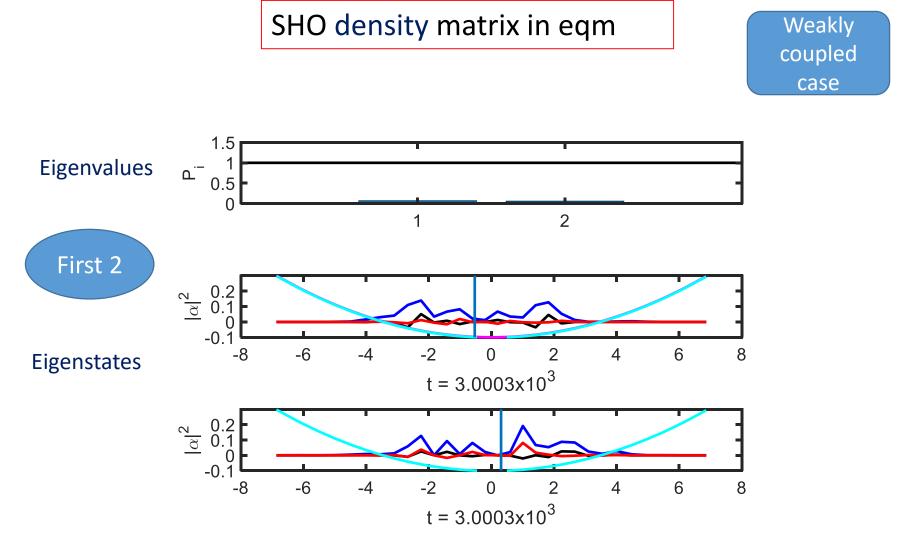


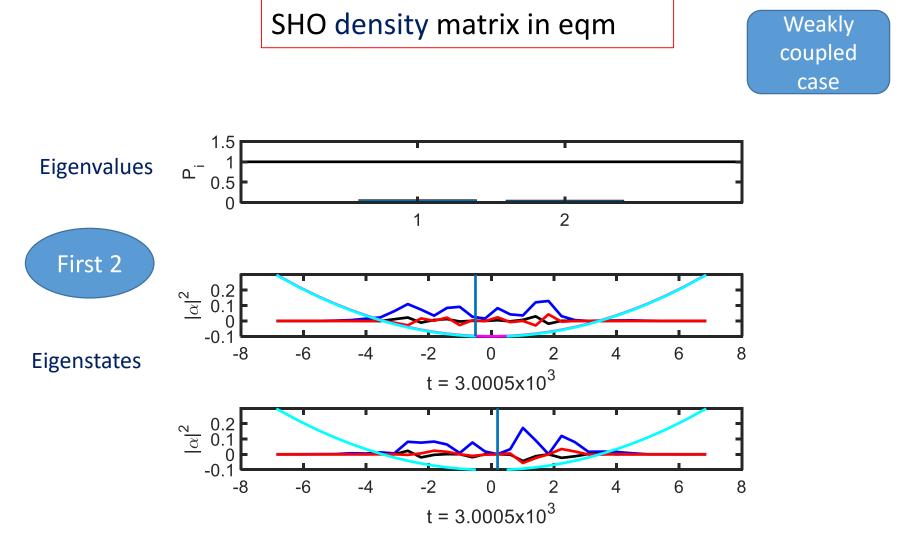
Next movie shows weak coupling case, starting here

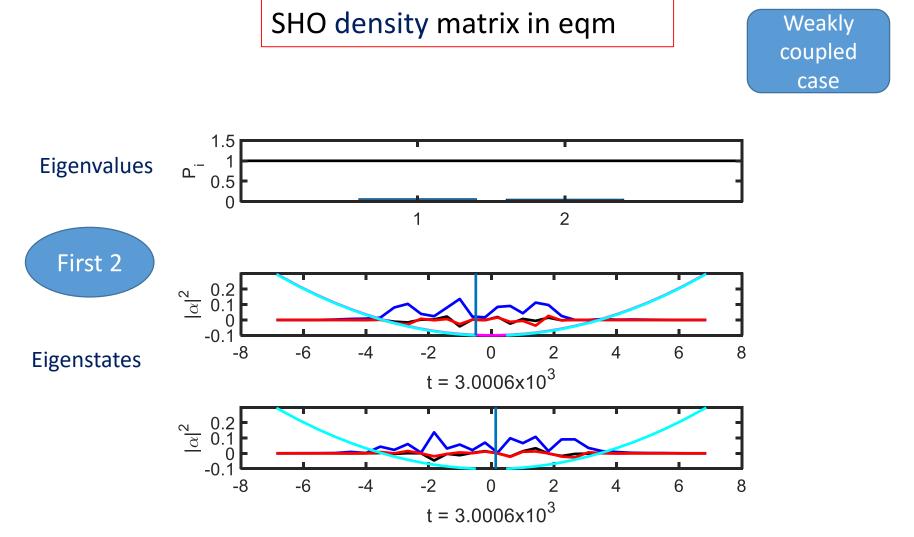


SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 2 6 -4 0 4 8 -8 Eigenstates $t = 3x10^{3}$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 0 2 6 8 -8 4 $t = 3 \times 10^3$

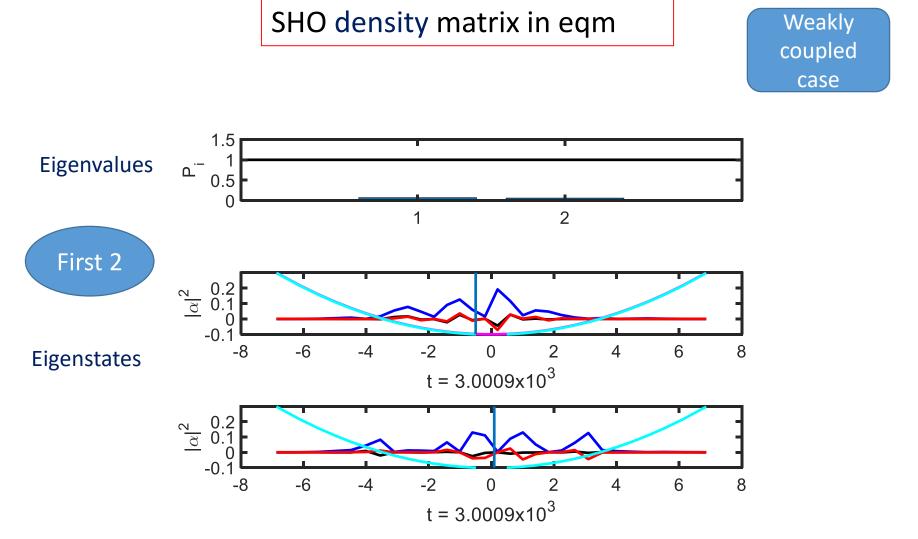


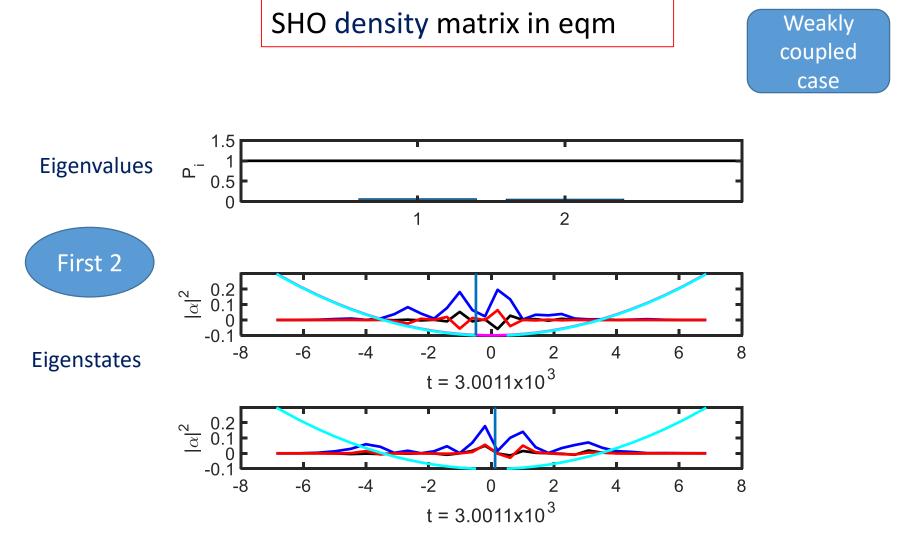


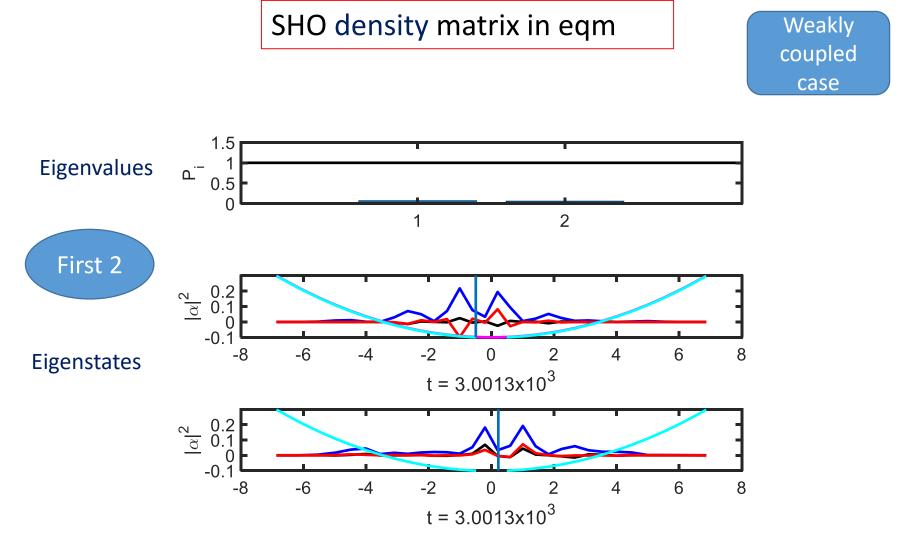


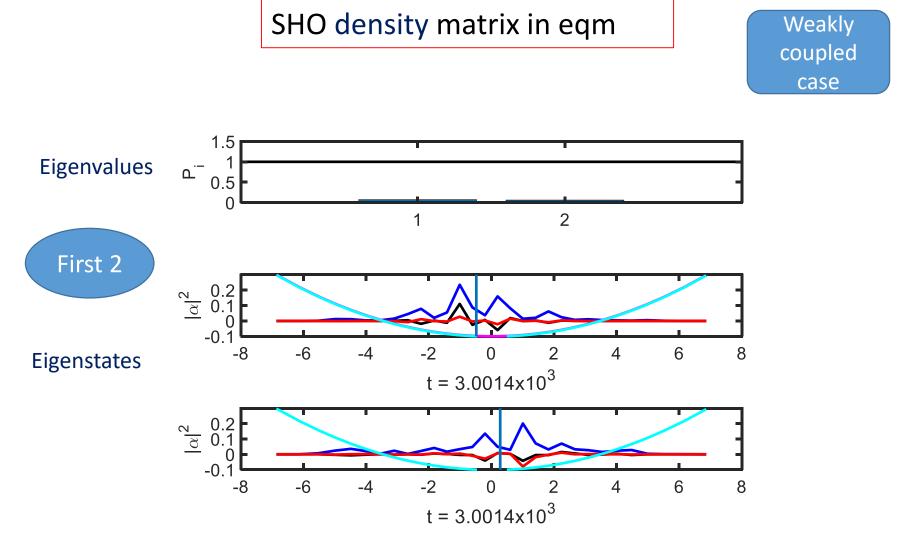


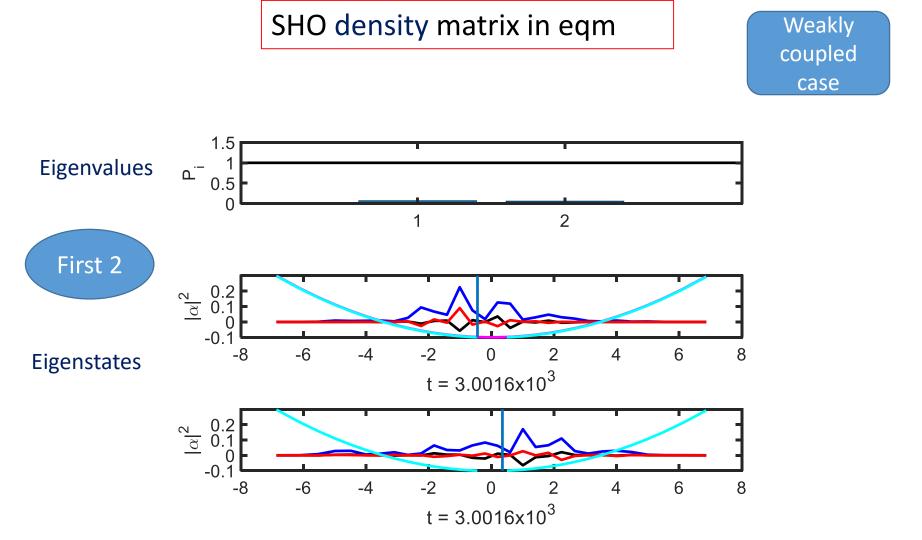
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -6 -2 6 -4 2 4 8 -8 0 Eigenstates $t = 3.0008 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 0 2 6 8 -8 -4 4 $t = 3.0008 \times 10^3$

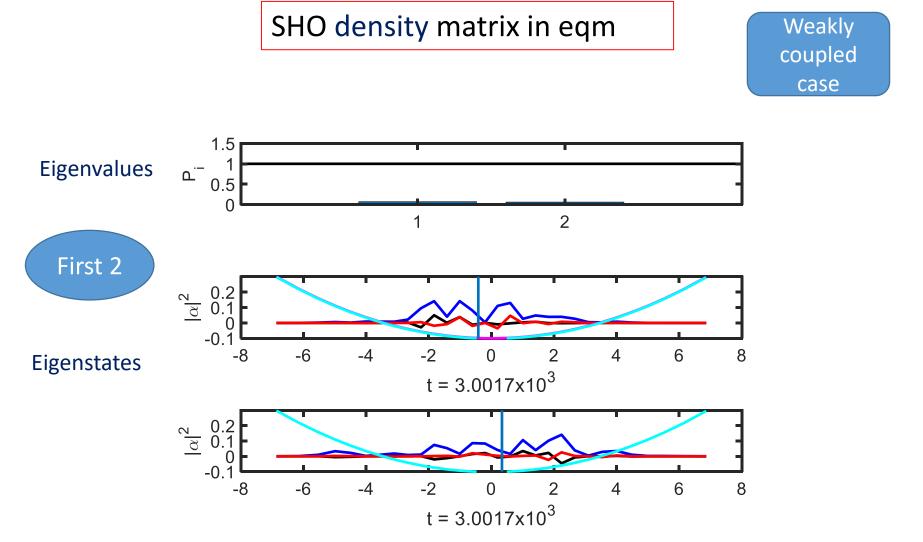


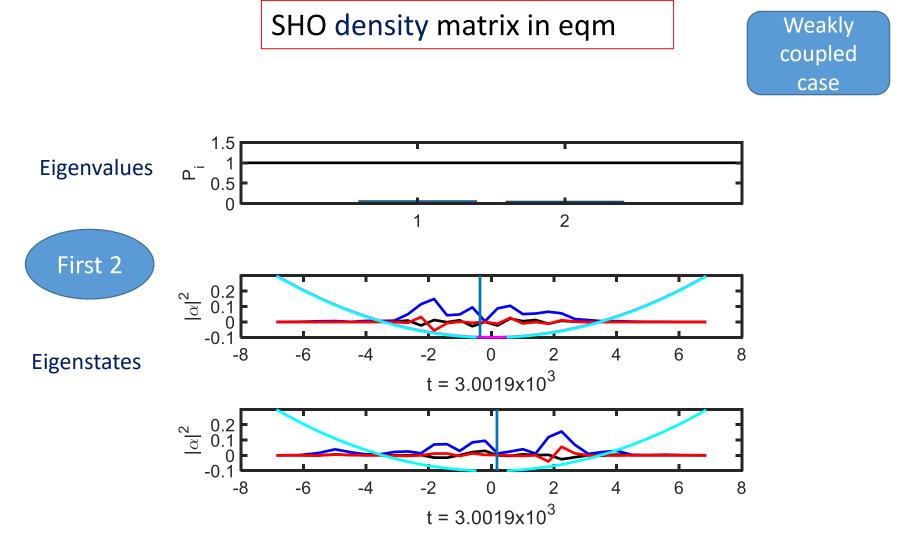


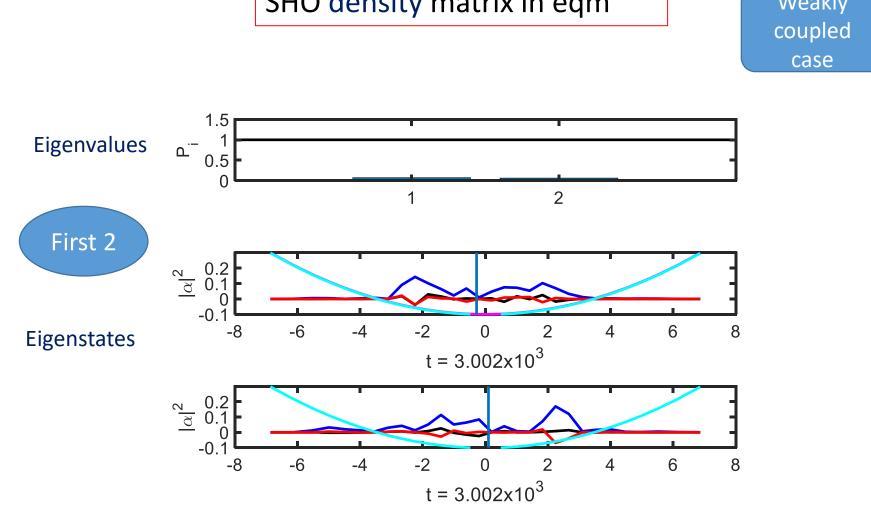




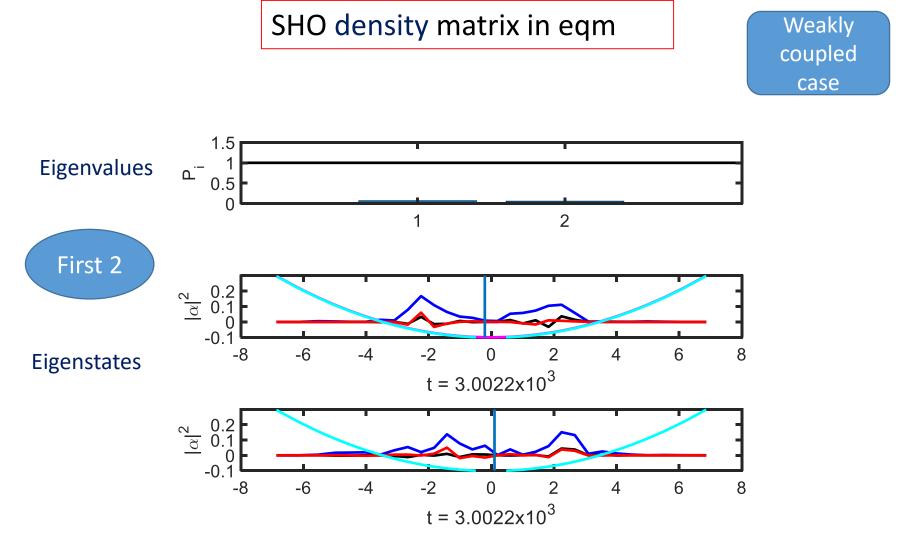


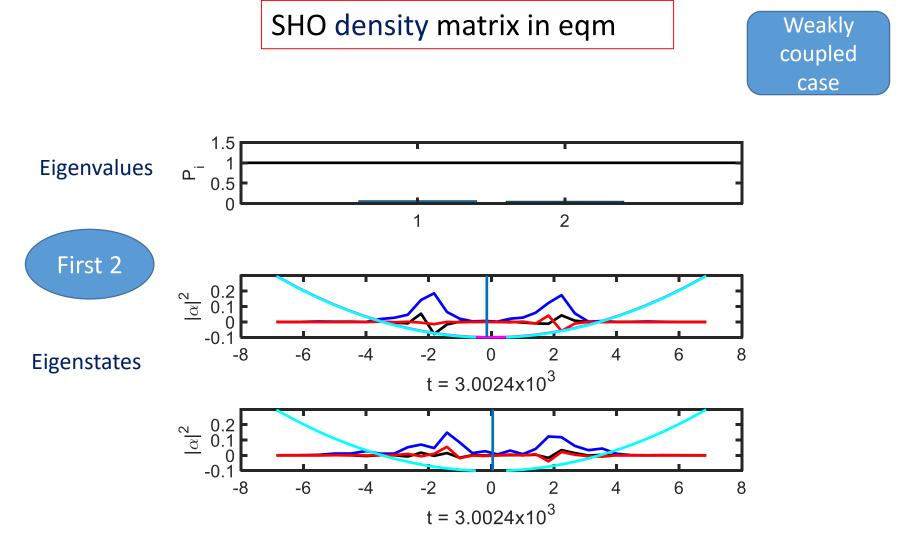


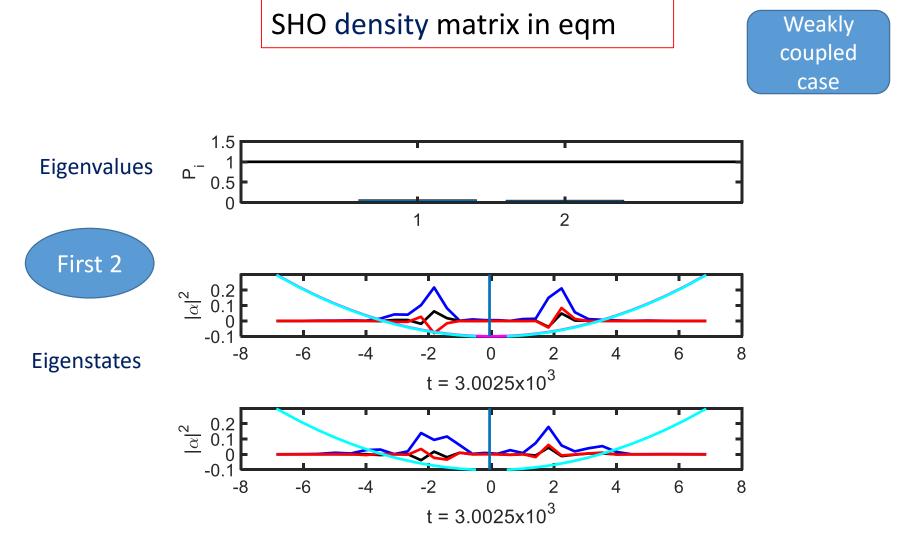


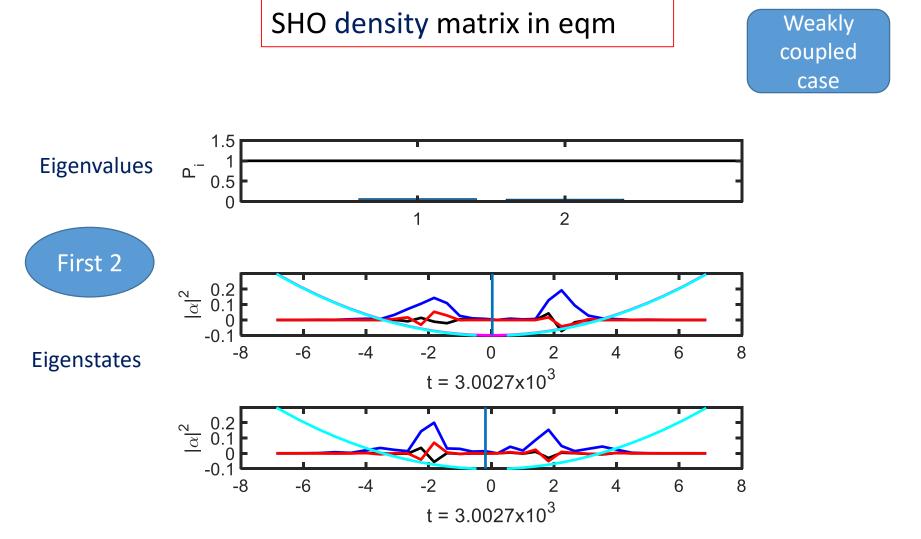


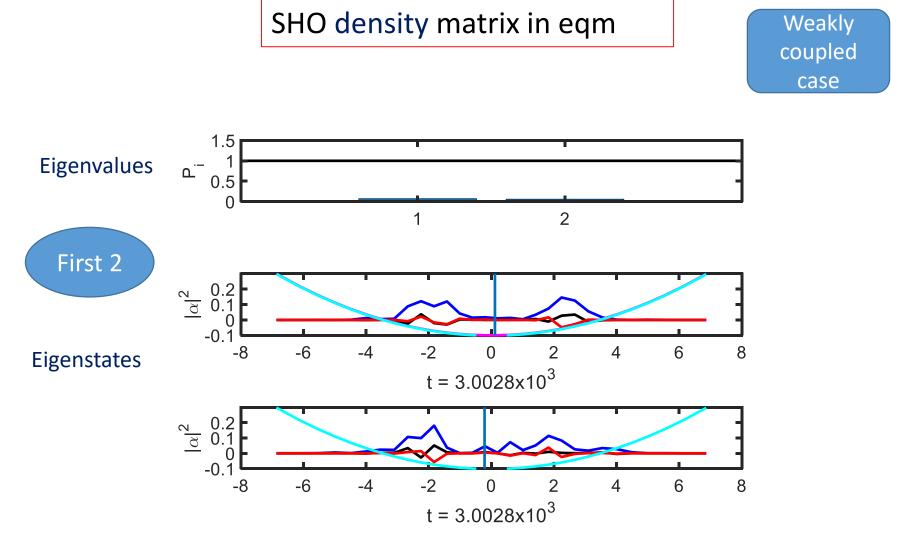
Weakly

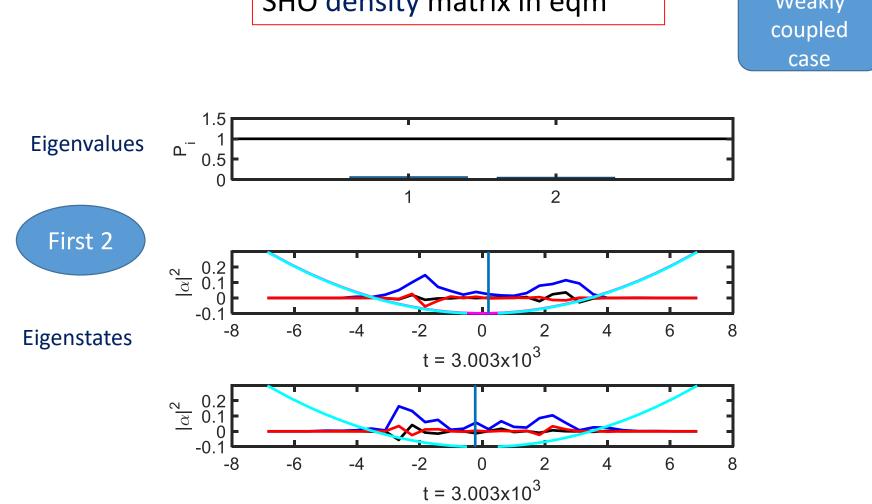








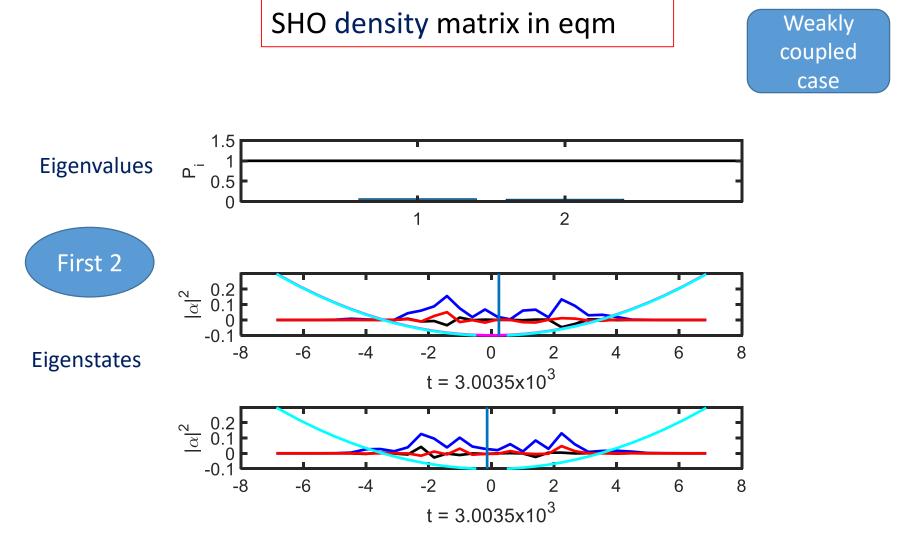


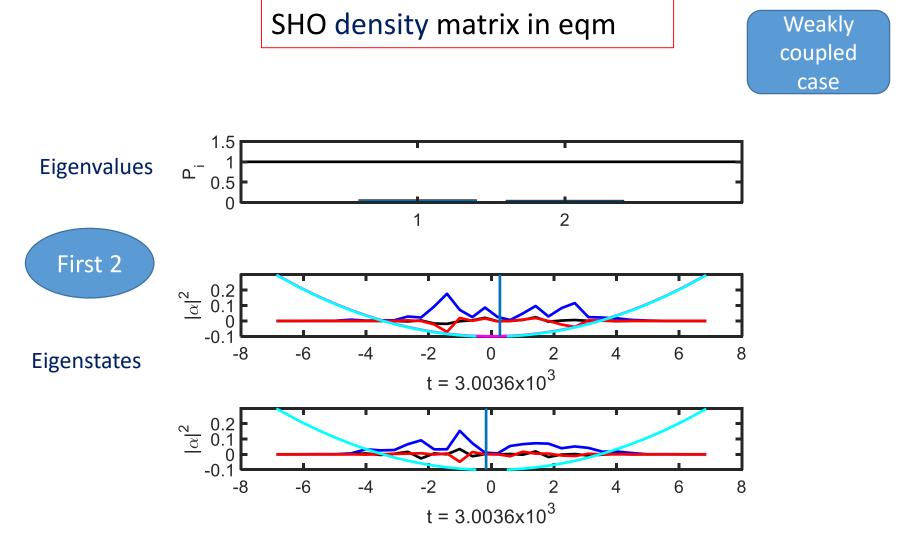


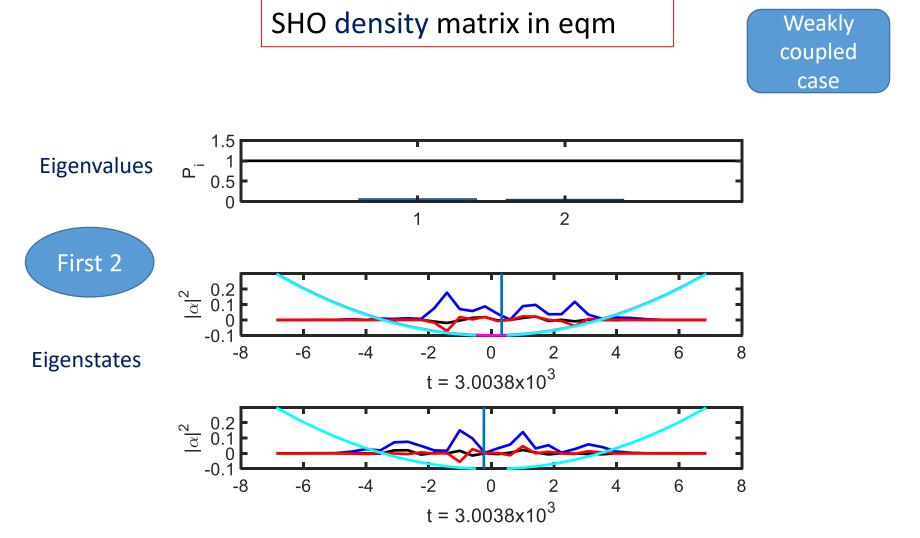


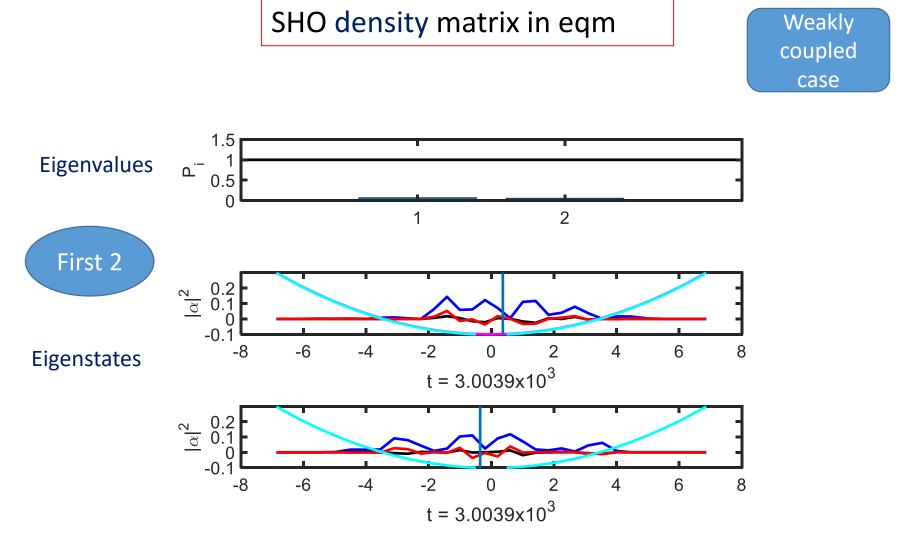
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -6 -2 6 -4 0 2 4 8 -8 Eigenstates $t = 3.0031 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 2 6 8 -8 0 4 $t = 3.0031 \times 10^3$

SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 6 -4 2 4 8 -8 0 Eigenstates $t = 3.0033 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 2 6 8 -8 0 4 $t = 3.0033 \times 10^3$









SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 6 -4 2 4 8 -8 0 Eigenstates $t = 3.0041 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 0 2 6 8 -8 -4 4 $t = 3.0041 \times 10^3$

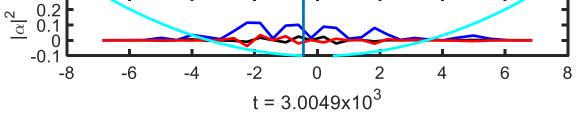
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 6 -4 2 4 8 -8 0 Eigenstates $t = 3.0042 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 0 2 6 8 -8 4 $t = 3.0042 \times 10^3$

SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 6 -4 2 4 8 -8 0 Eigenstates $t = 3.0044 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 2 6 8 -8 -4 0 4 $t = 3.0044 \times 10^3$

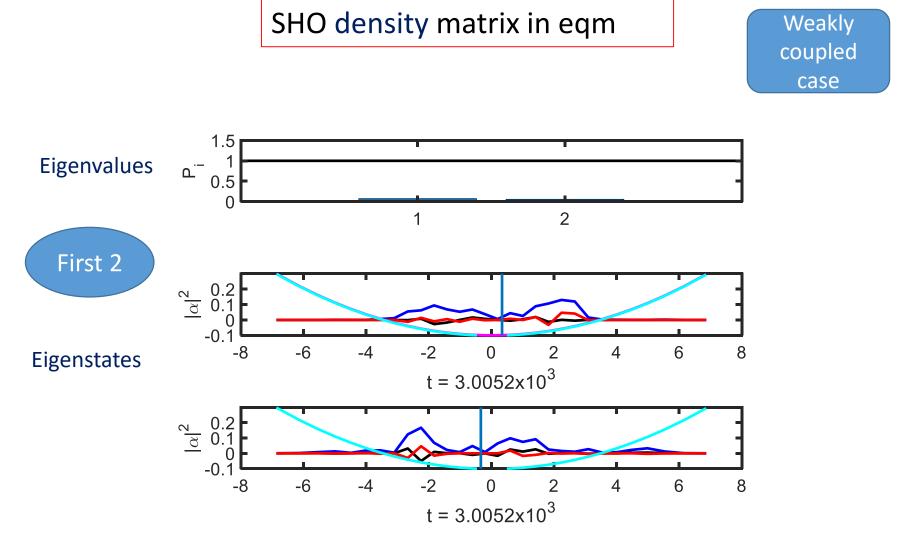
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 6 -4 2 4 8 -8 0 Eigenstates $t = 3.0046 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 2 6 8 -8 0 4 $t = 3.0046 \times 10^3$

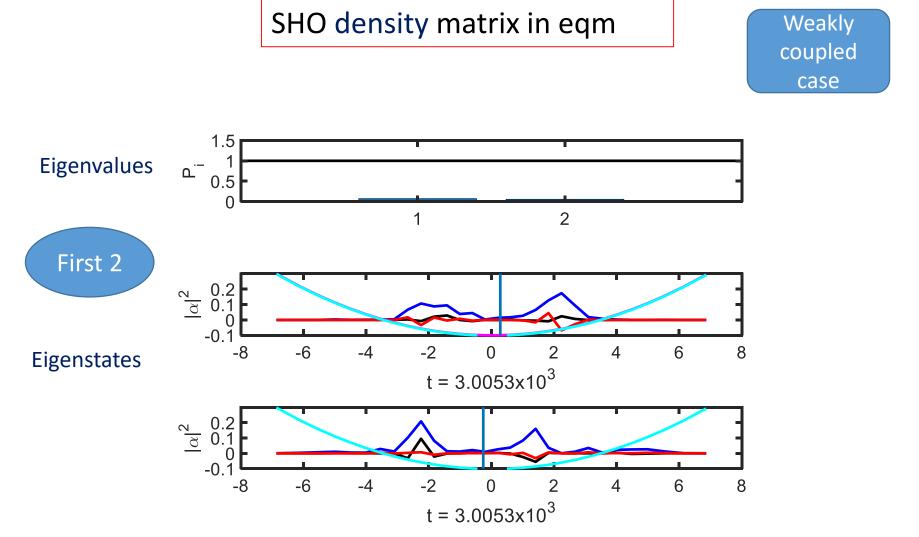
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 6 -4 2 4 8 -8 0 Eigenstates $t = 3.0047 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 2 6 8 -8 0 4 $t = 3.0047 \times 10^3$

SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 6 -4 0 2 4 8 -8 Eigenstates $t = 3.0049 \times 10^3$

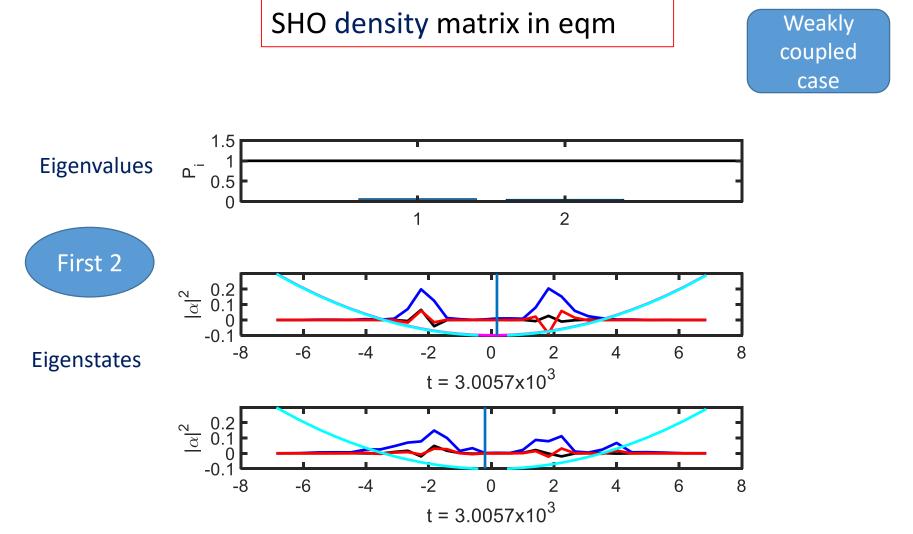


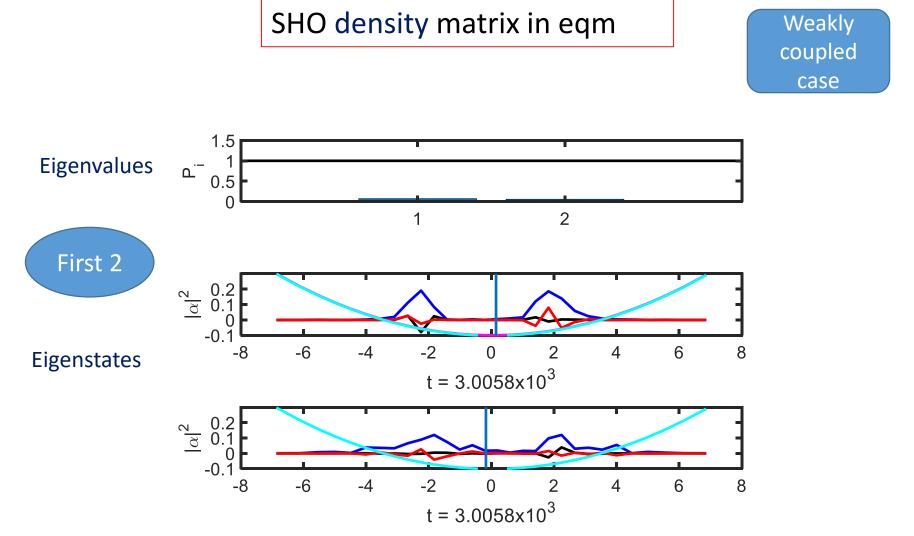
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 6 -4 0 2 4 8 -8 Eigenstates $t = 3.005 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 2 6 8 -8 0 4 $t = 3.005 \times 10^3$

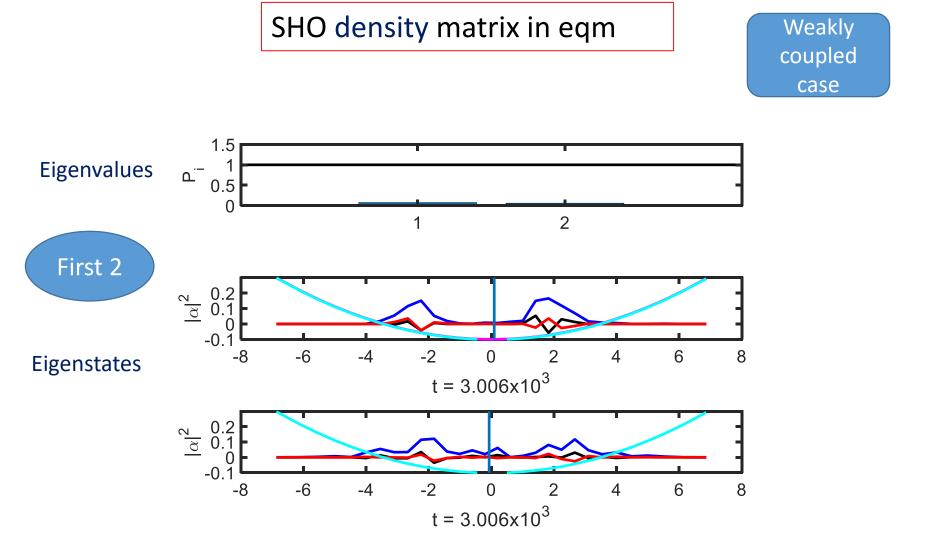




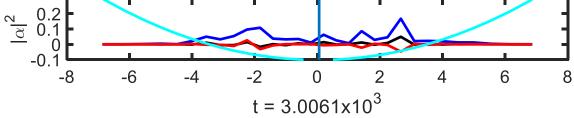
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 0 6 -4 2 4 8 -8 Eigenstates $t = 3.0055 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 0 2 6 8 -8 4 $t = 3.0055 \times 10^3$







SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 6 0 2 4 8 -8 -4 Eigenstates $t = 3.0061 \times 10^3$ 0.2



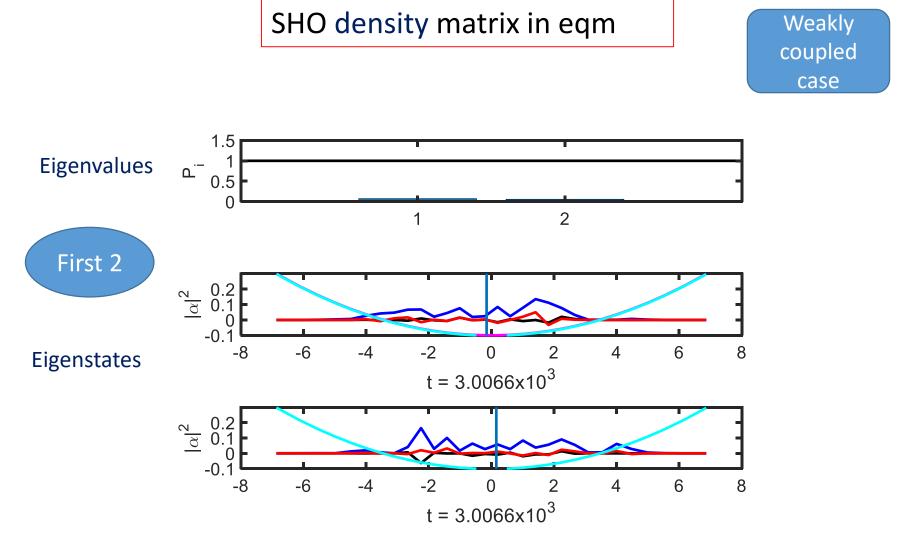
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 6 -6 -4 2 4 8 -8 0 Eigenstates $t = 3.0063 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 0 2 6 8 -8 4 $t = 3.0063 \times 10^3$

case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -6 -2 6 -4 2 4 8 -8 0 Eigenstates $t = 3.0064 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 0 2 6 8 -8 4

 $t = 3.0064 \times 10^3$

SHO density matrix in eqm

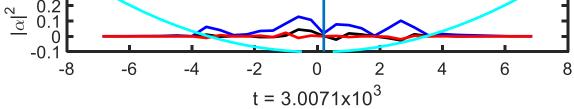
Weakly coupled

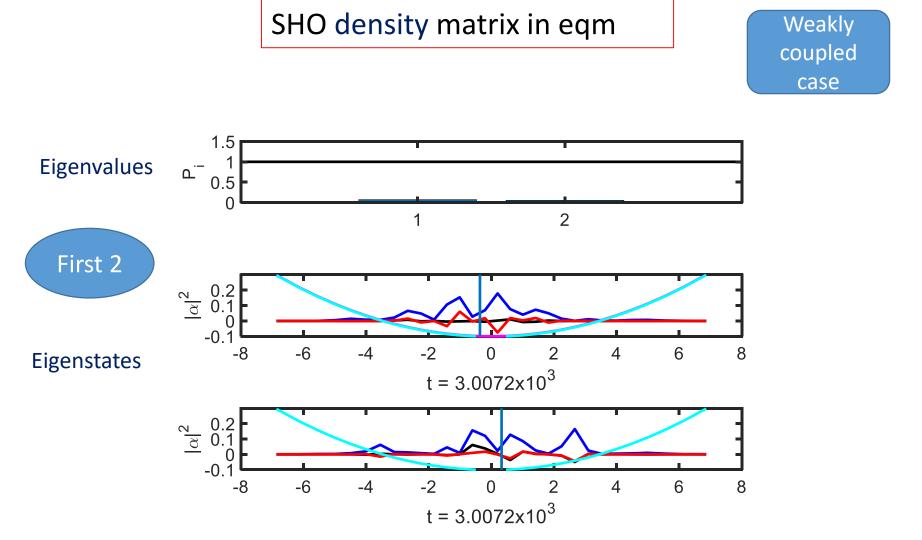


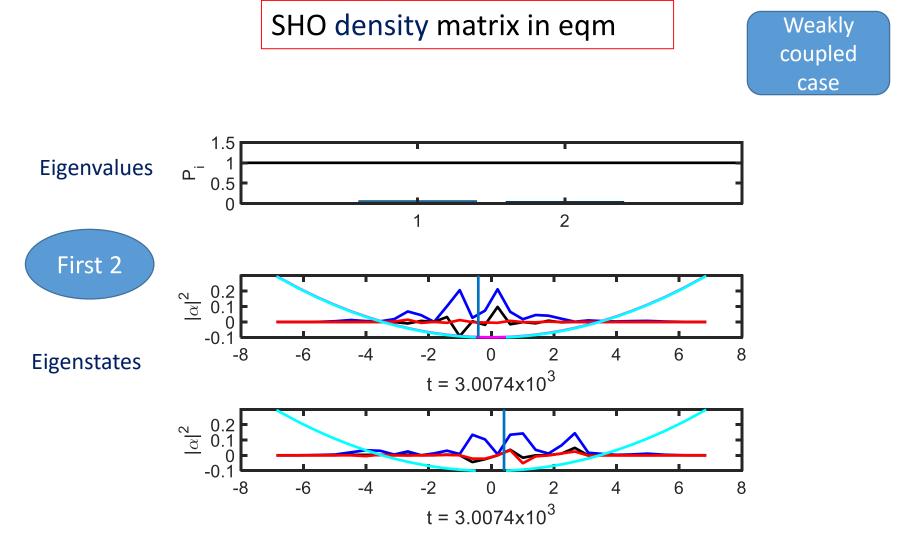
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -6 -2 6 -4 2 4 8 -8 0 Eigenstates $t = 3.0068 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 0 2 6 8 -8 4 $t = 3.0068 \times 10^3$

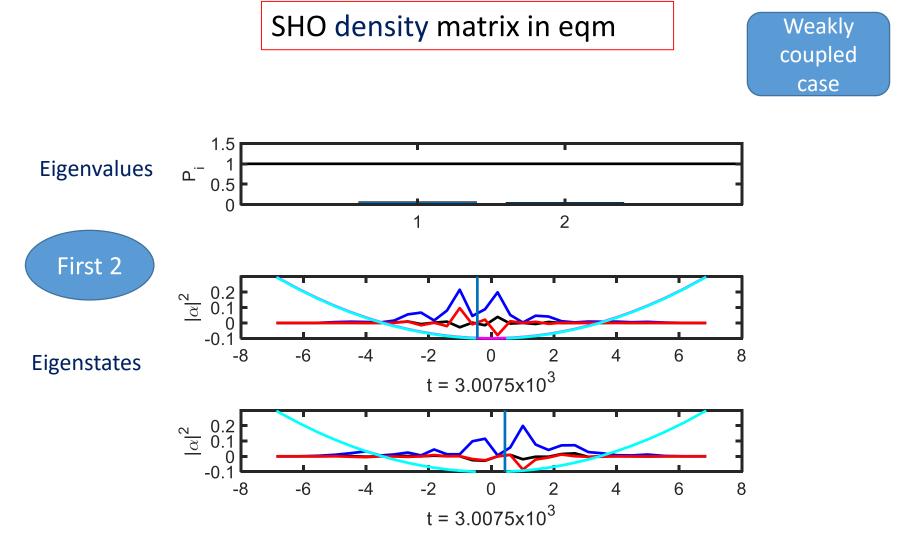
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -6 -2 6 2 4 8 -8 -4 0 Eigenstates $t = 3.0069 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 0 2 6 8 -8 -4 4 $t = 3.0069 \times 10^3$

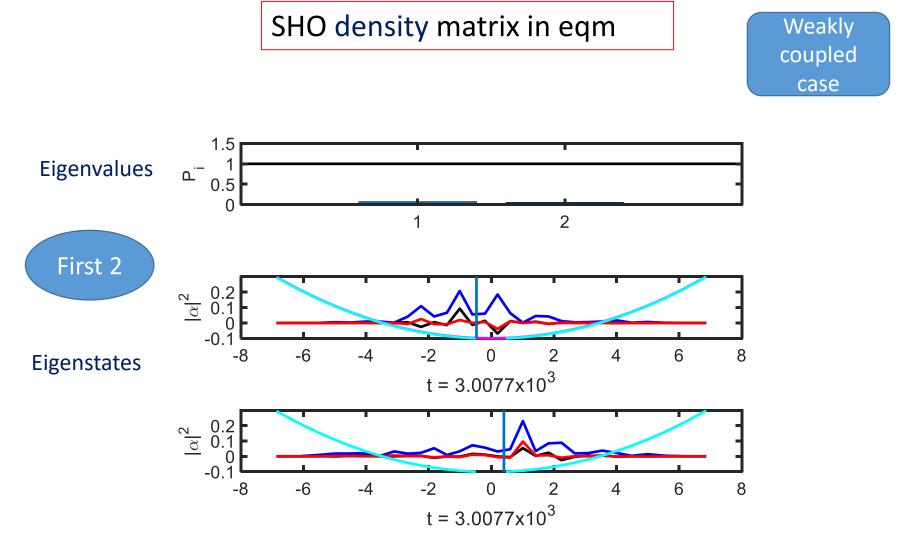
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -6 -2 6 -4 2 4 8 -8 0 Eigenstates $t = 3.0071 \times 10^3$ 0.2

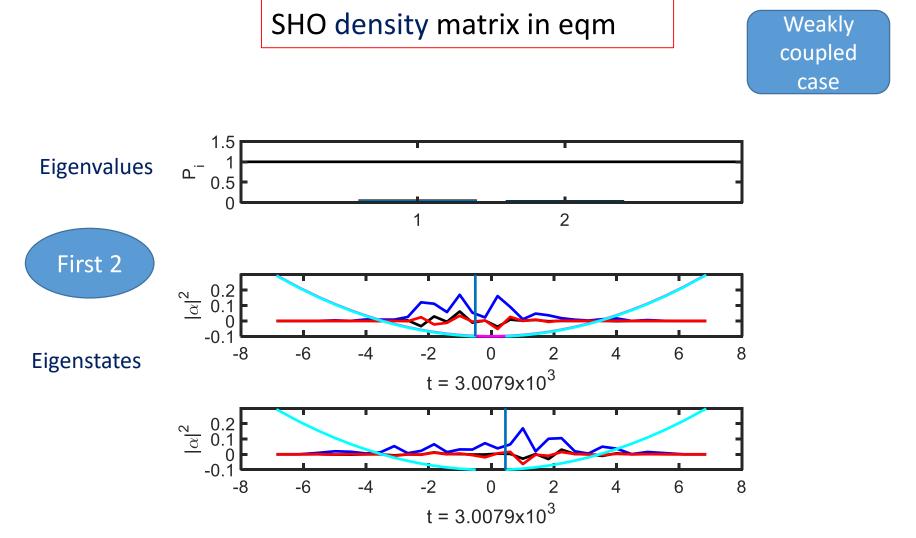




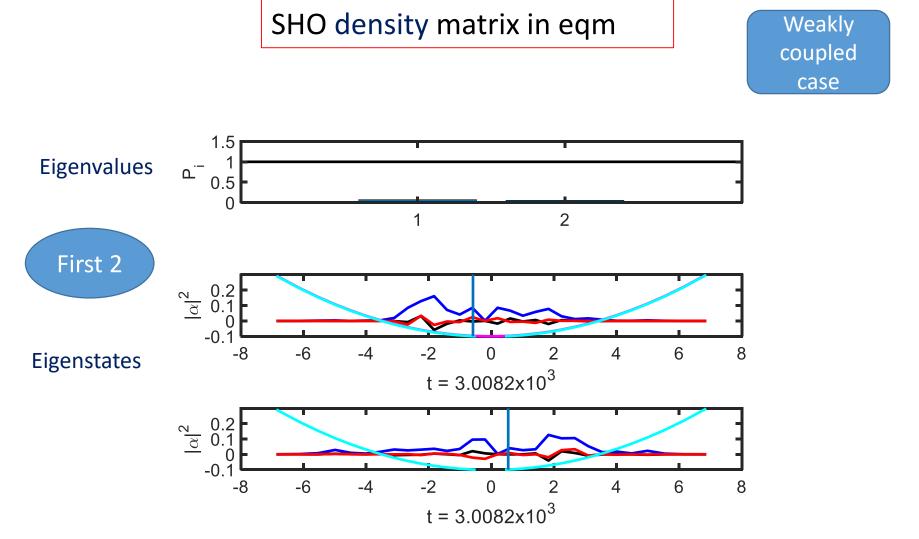


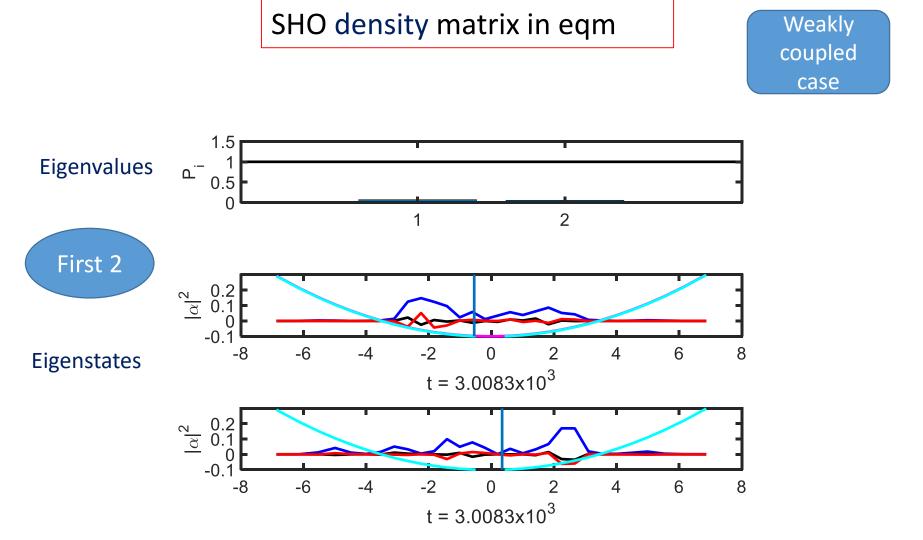


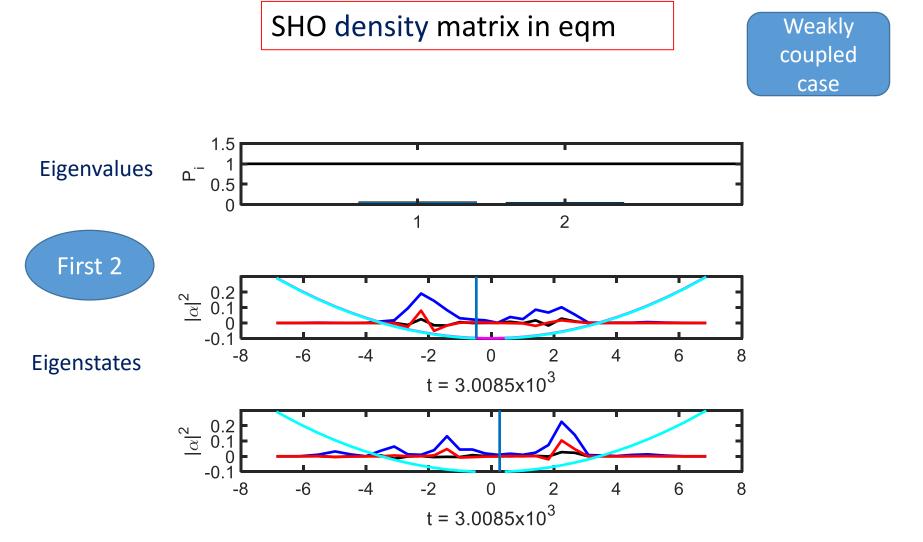


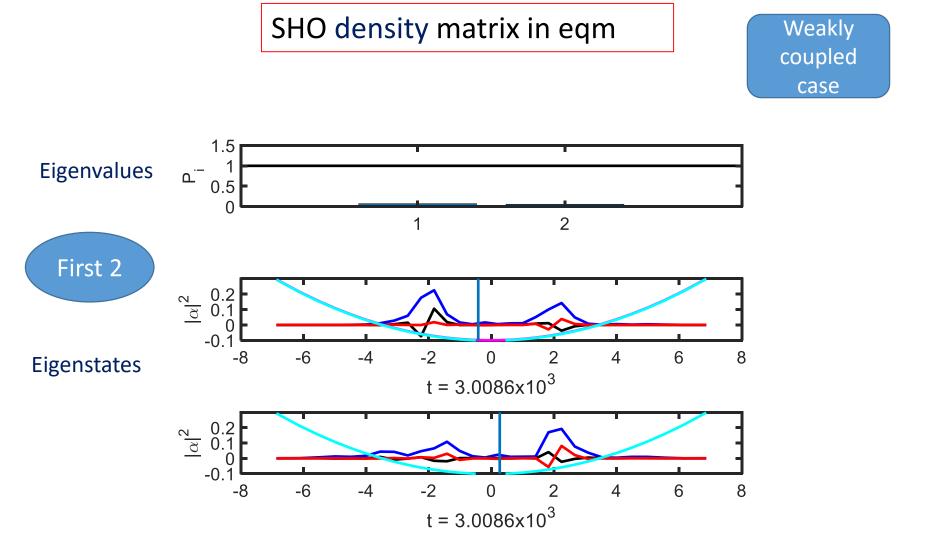


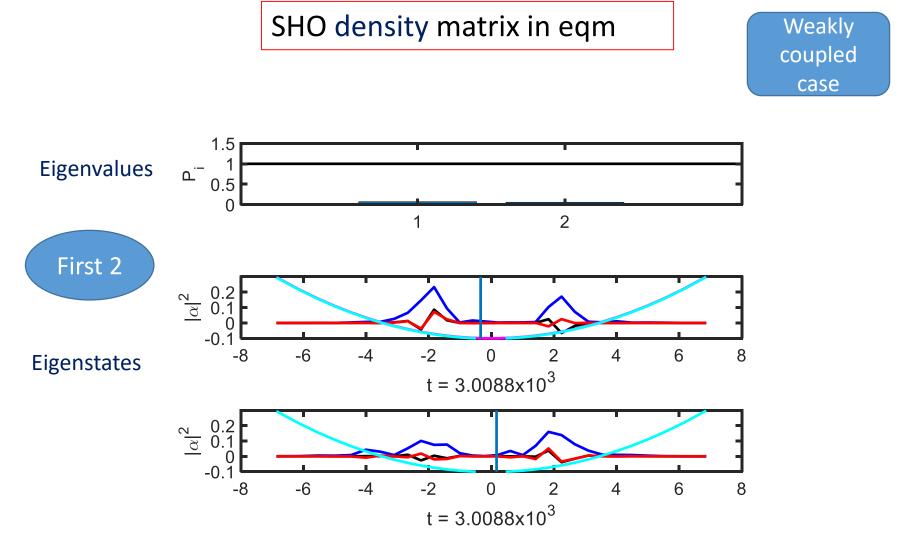
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 6 -6 -4 2 4 8 -8 0 Eigenstates $t = 3.008 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 0 2 6 8 -8 -4 4 $t = 3.008 \times 10^3$



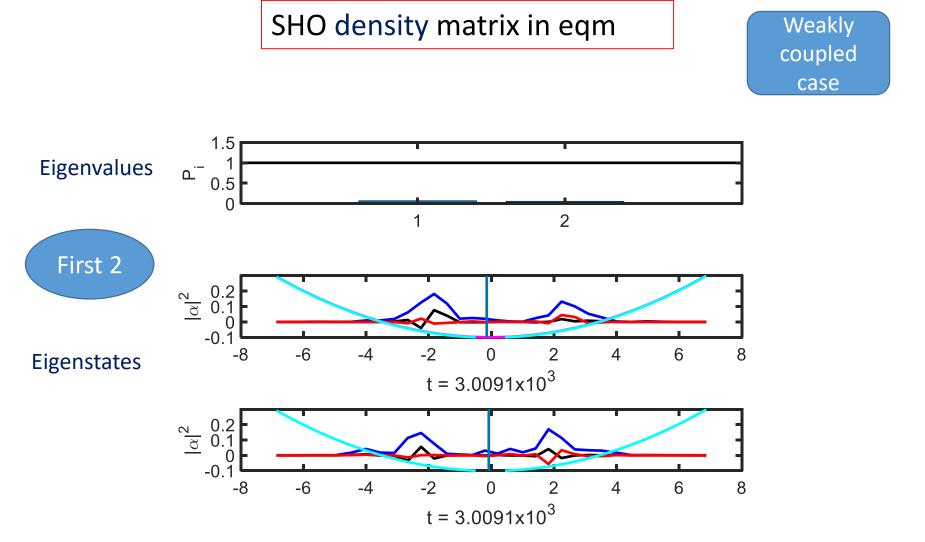


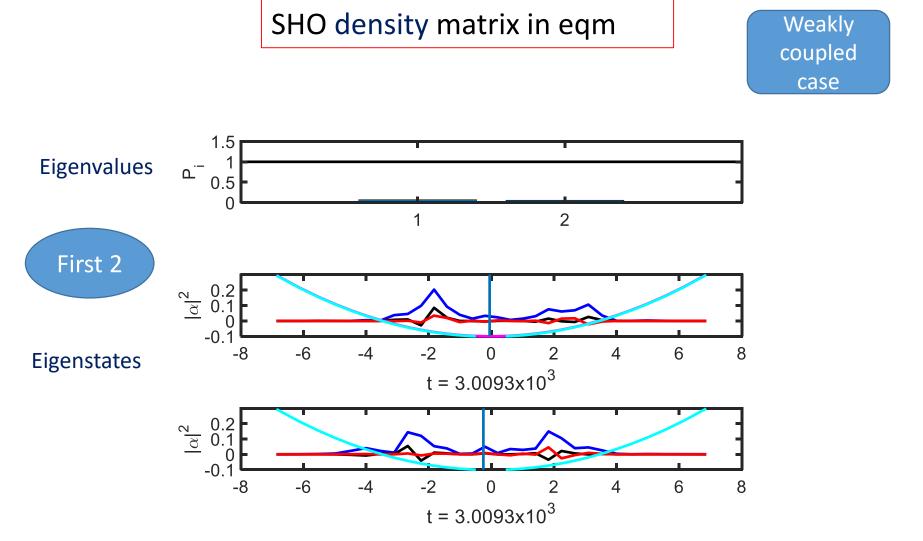






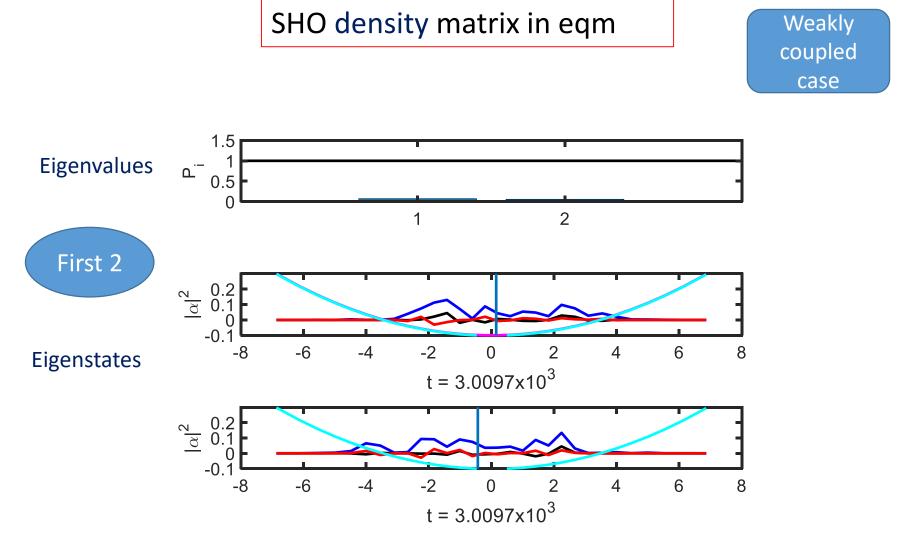
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 2 6 -4 0 4 8 -8 Eigenstates $t = 3.009 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 0 2 6 8 -8 4 $t = 3.009 \times 10^3$

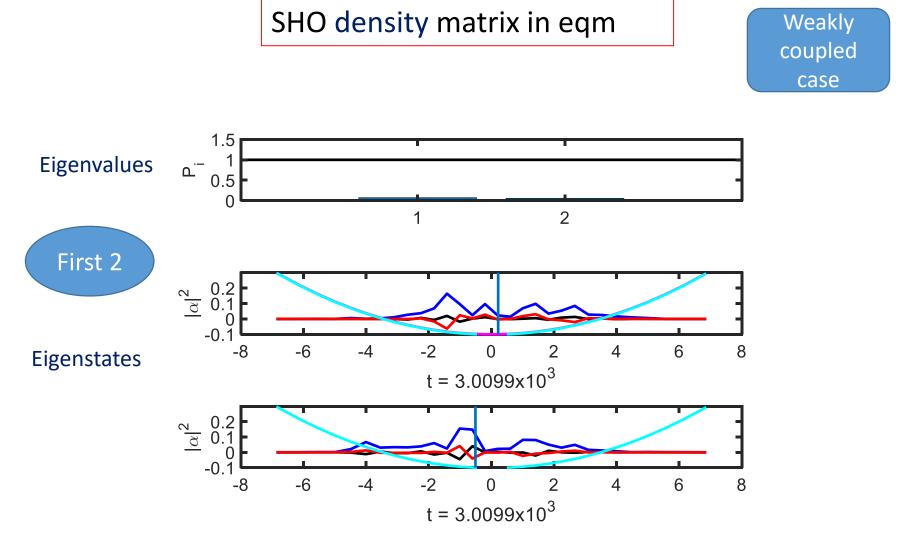


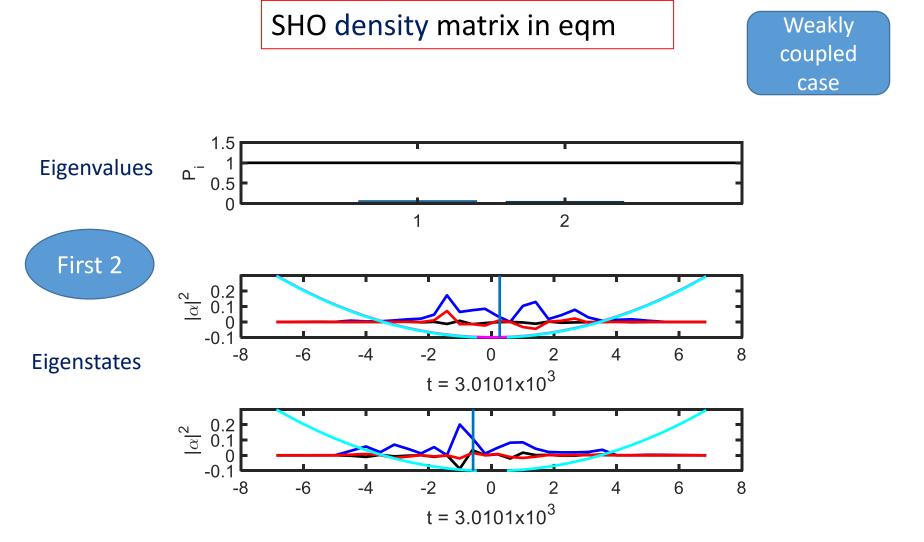


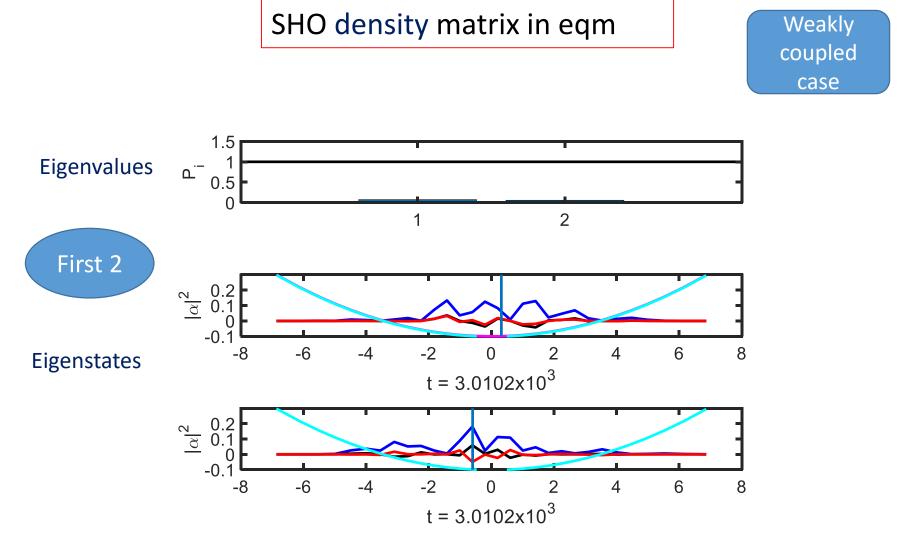
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 6 -6 -4 2 4 8 -8 0 Eigenstates $t = 3.0094 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 0 2 6 8 -8 -4 4 $t = 3.0094 \times 10^3$

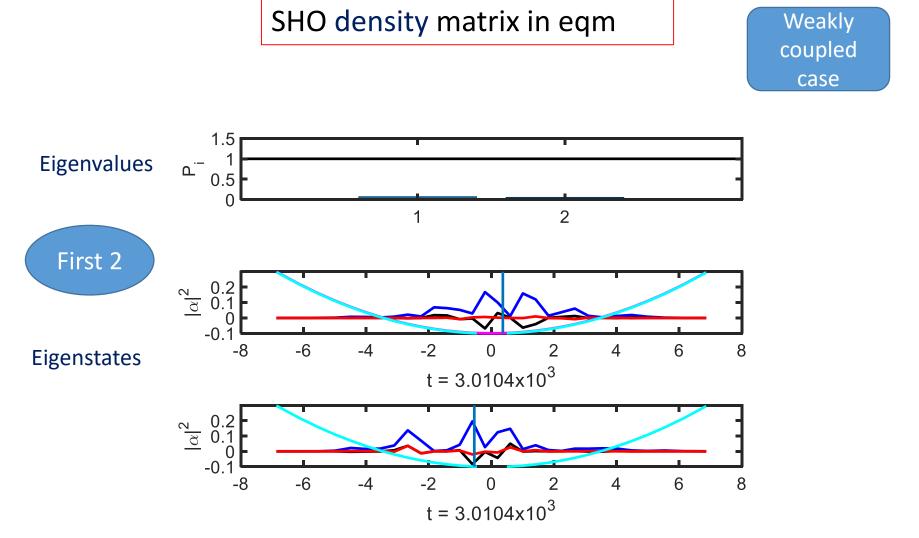
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 6 -4 2 4 8 -8 0 Eigenstates $t = 3.0096 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 2 6 8 -8 0 4 $t = 3.0096 \times 10^3$

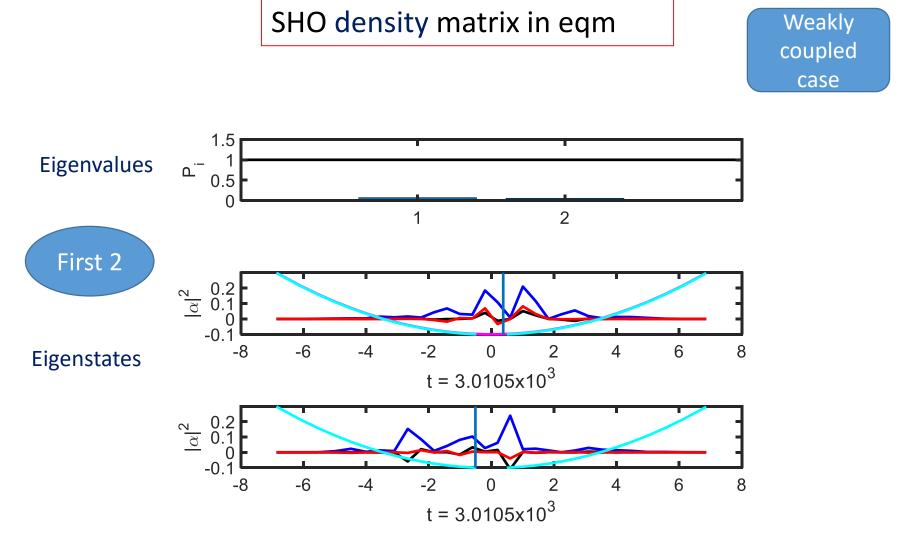


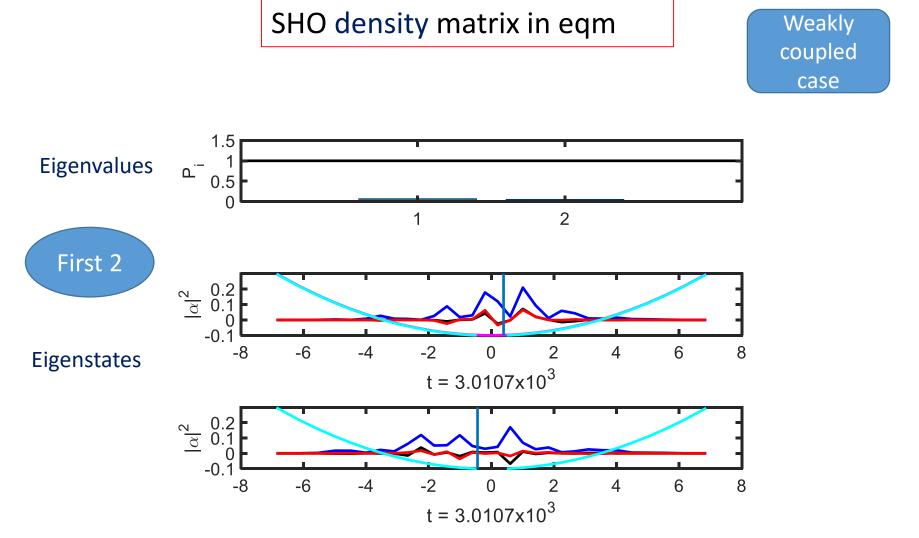


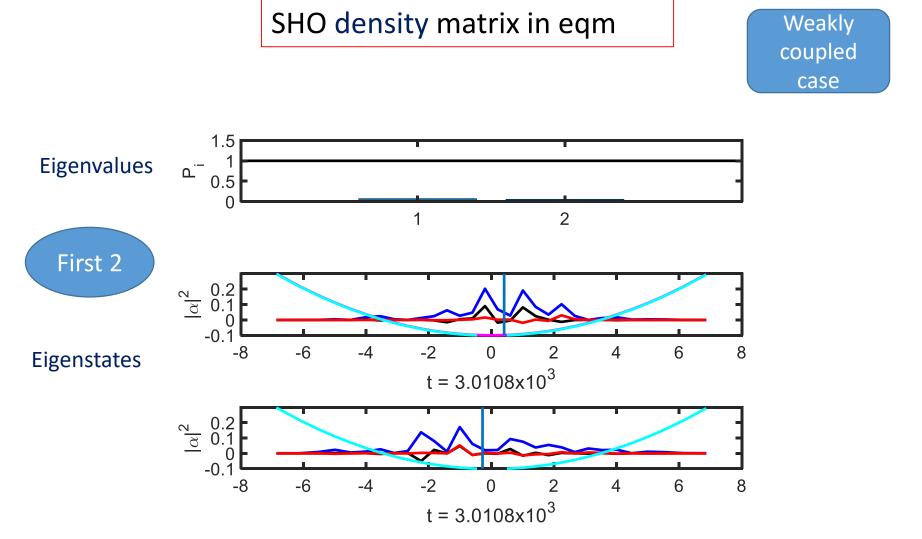












SHO density matrix in eqm coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 6 -4 2 4 8 -8 0 Eigenstates $t = 3.011 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 2 6 8 -8 0 4

 $t = 3.011 \times 10^3$

Weakly

SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1

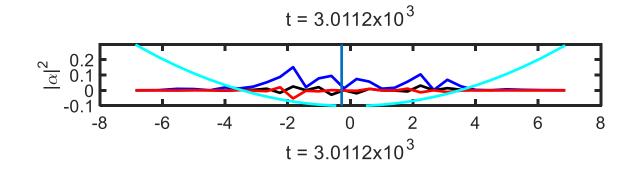
-2

-6

-8

-4

Eigenstates



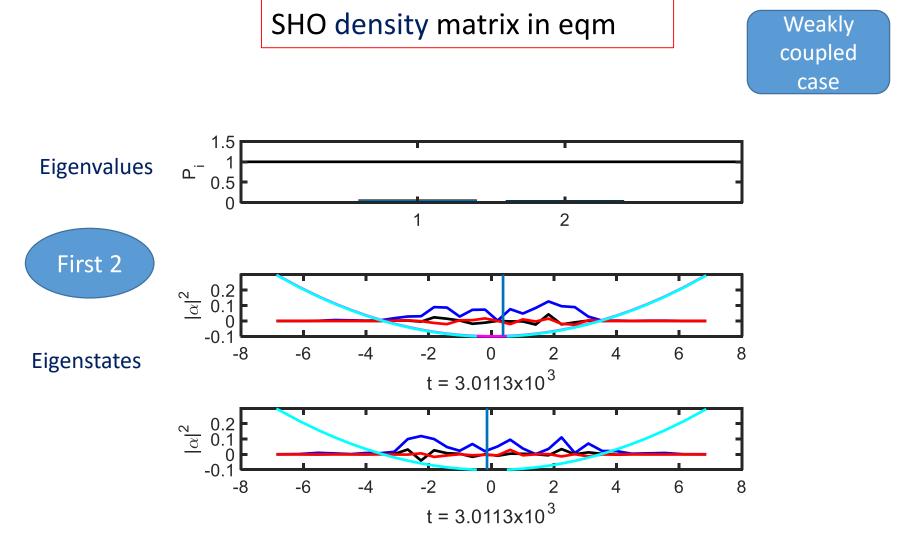
0

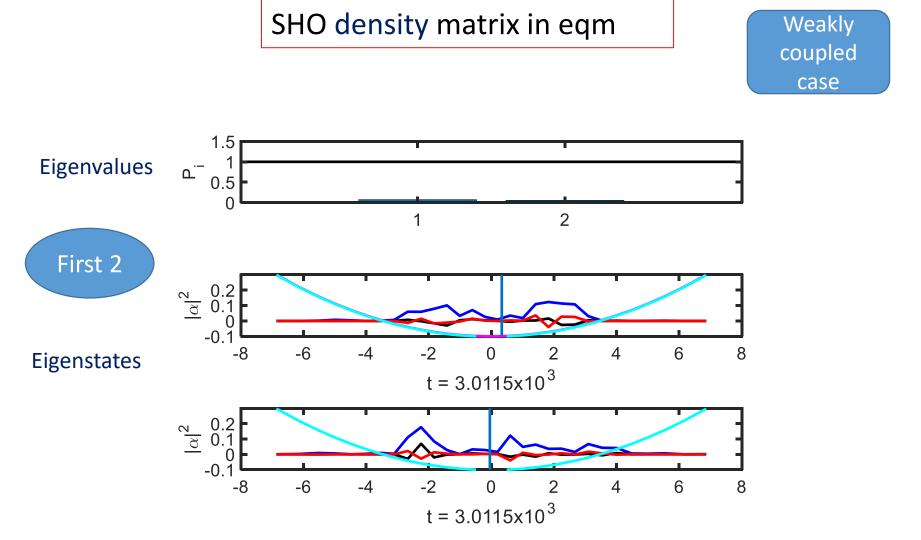
2

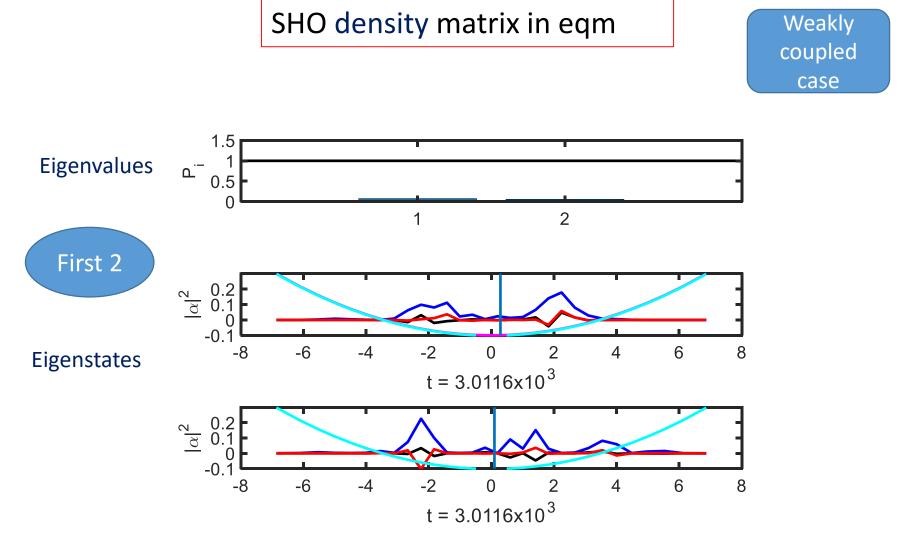
6

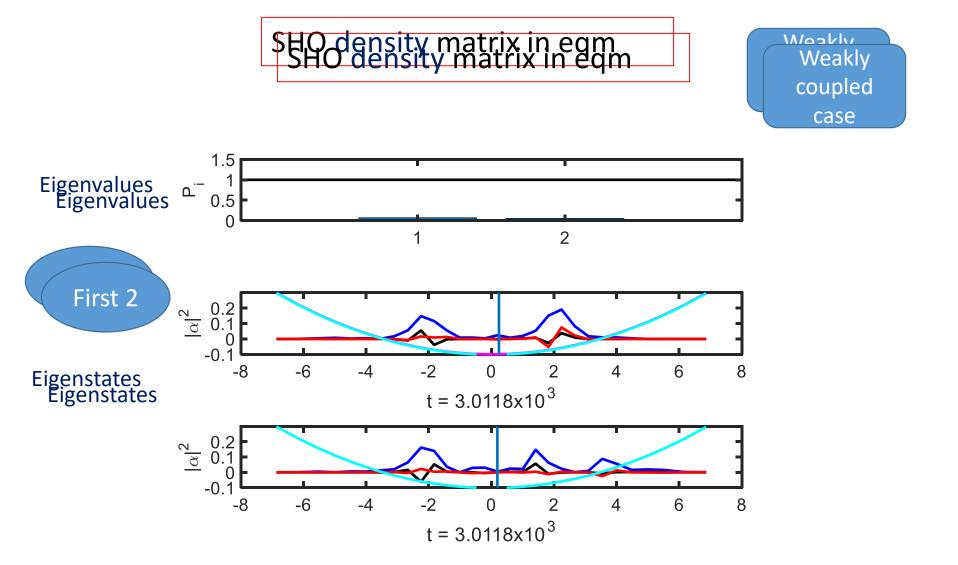
8

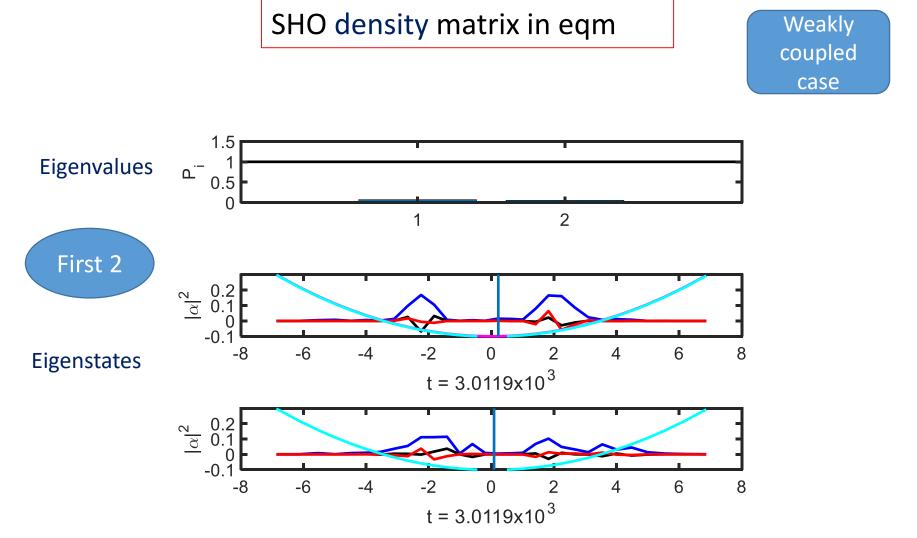
4











SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 0 6 -4 2 4 8 -8 Eigenstates $t = 3.0121 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 0 2 6 8 -8 -4 4 $t = 3.0121 \times 10^3$

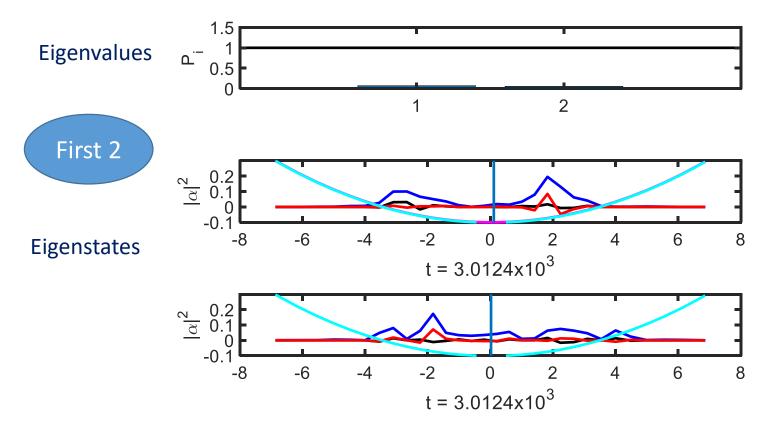
SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 6 -4 0 2 4 8 -8 Eigenstates $t = 3.0123 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 2 6 8 -8 0 4 $t = 3.0123 \times 10^3$

SHO density matrix in eqm

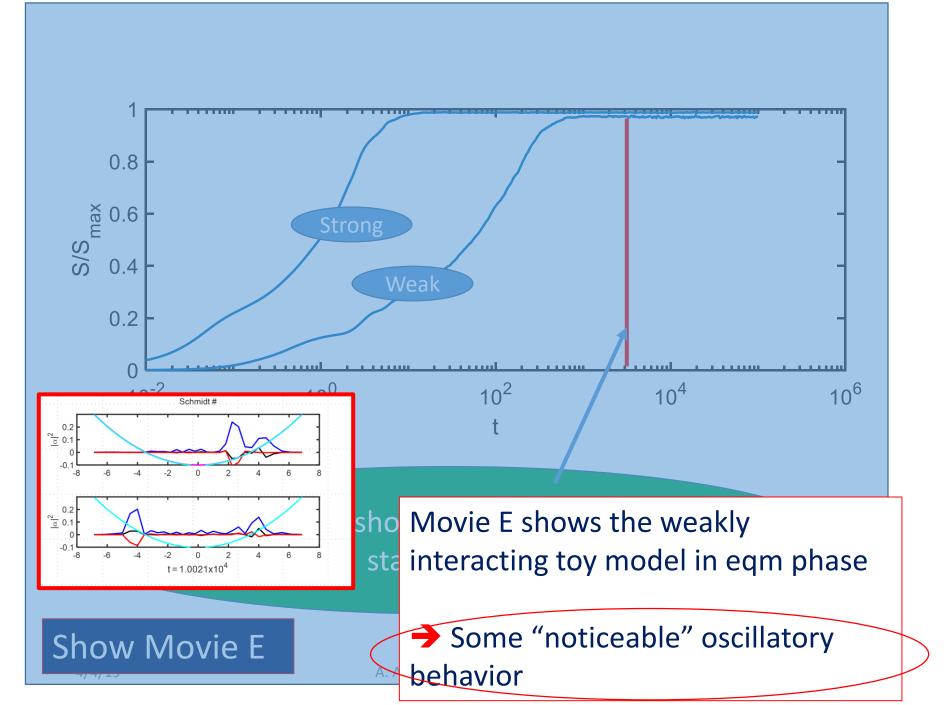
Weakly

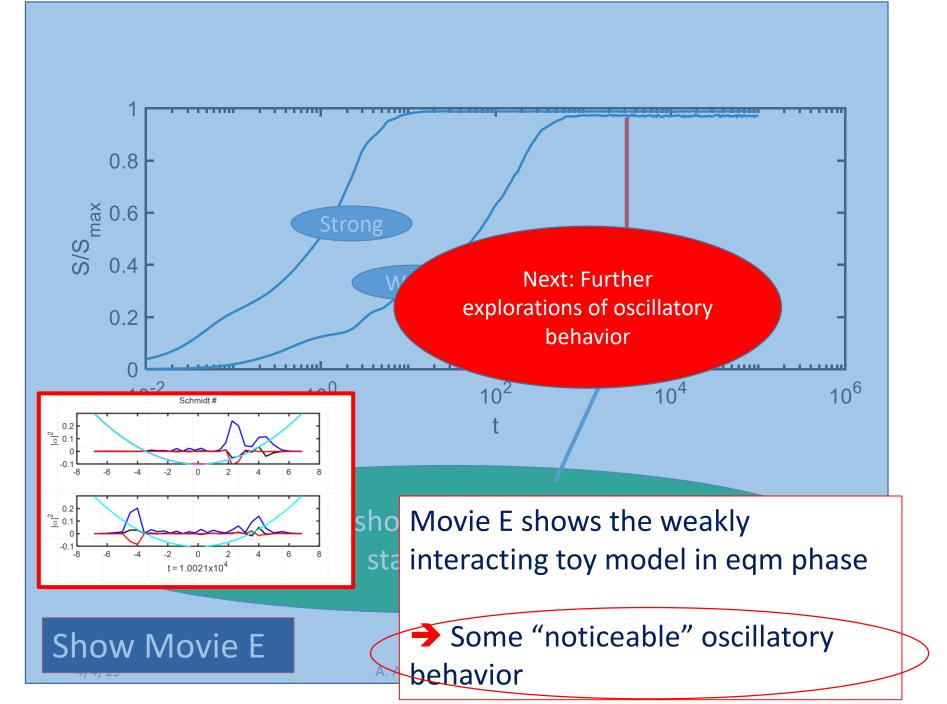
coupled

case

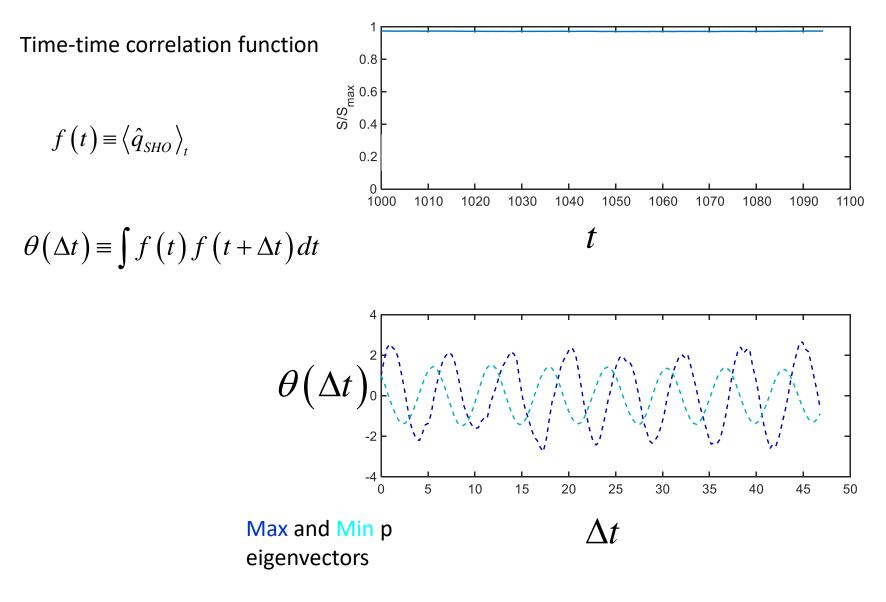


SHO density matrix in eqm Weakly coupled case 1.5 Eigenvalues 1 ۳.– 0.5 0 2 1 First 2 0.2 0.1 $|\alpha|^2$ 0 -0.1 -2 -6 6 -4 0 2 4 8 -8 Eigenstates $t = 3.0126 \times 10^3$ 0.2 $|\alpha|^2$ С -0.1 -2 -6 -4 0 2 6 8 -8 4 $t = 3.0126 \times 10^3$

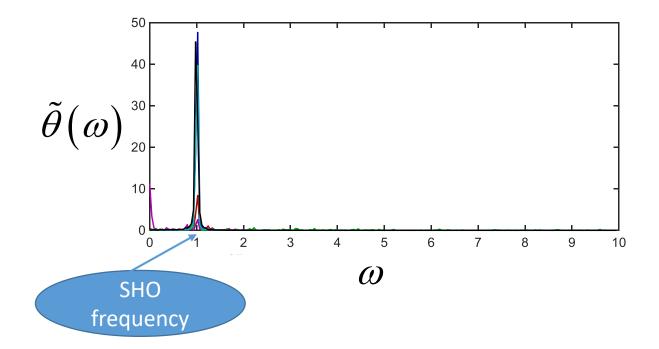


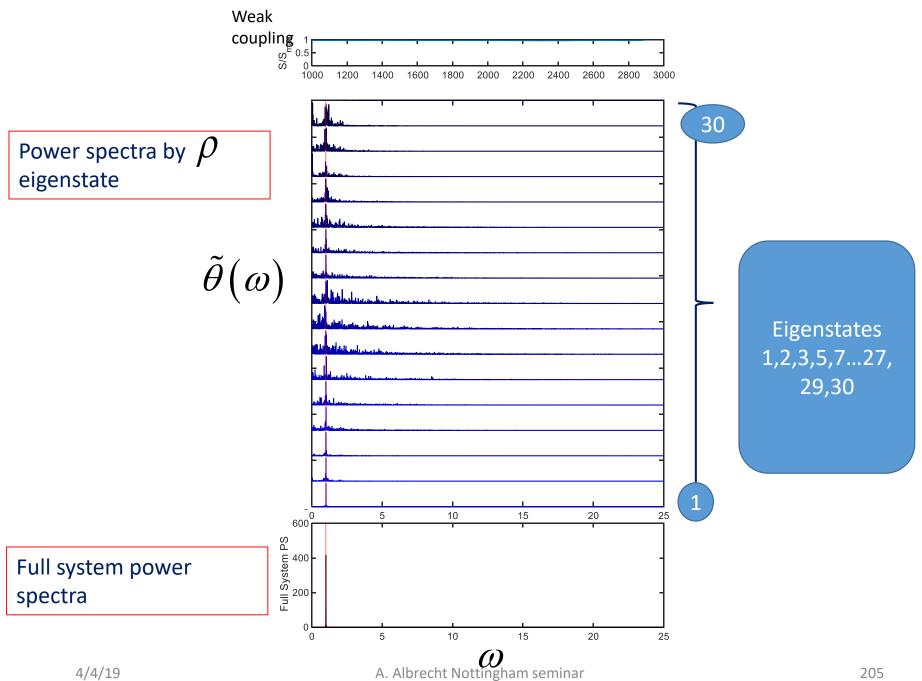


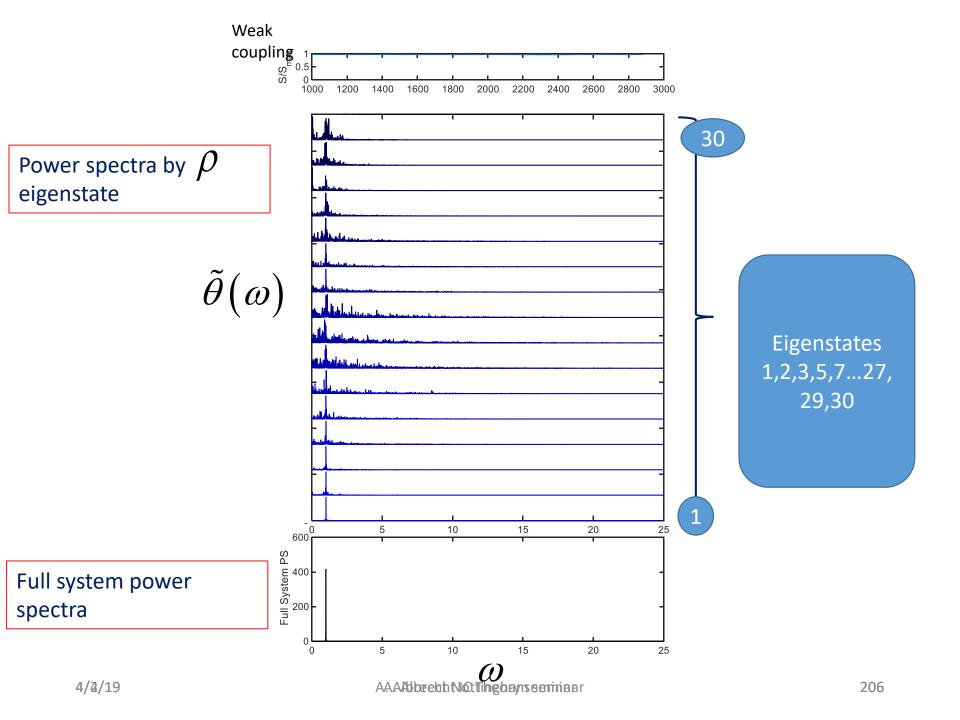
Further analysis:



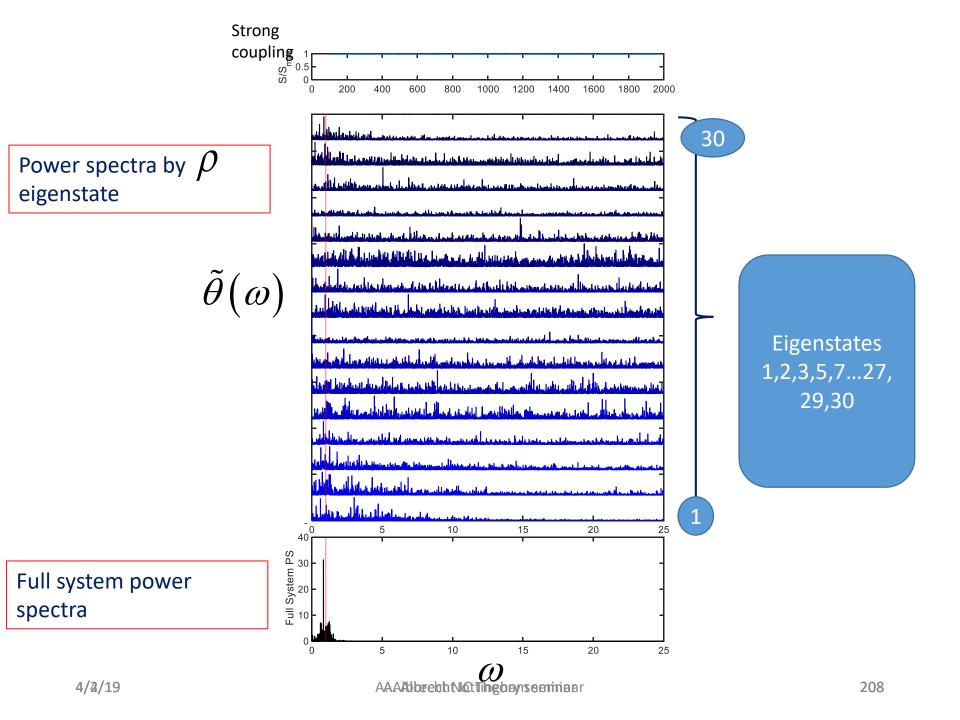
Power spectrum

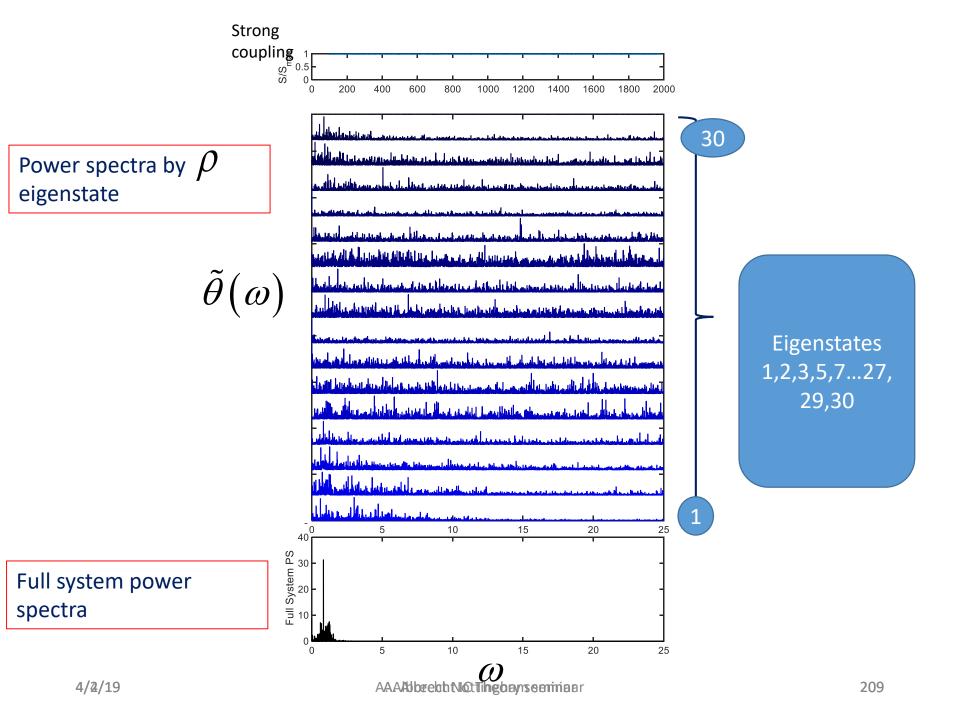


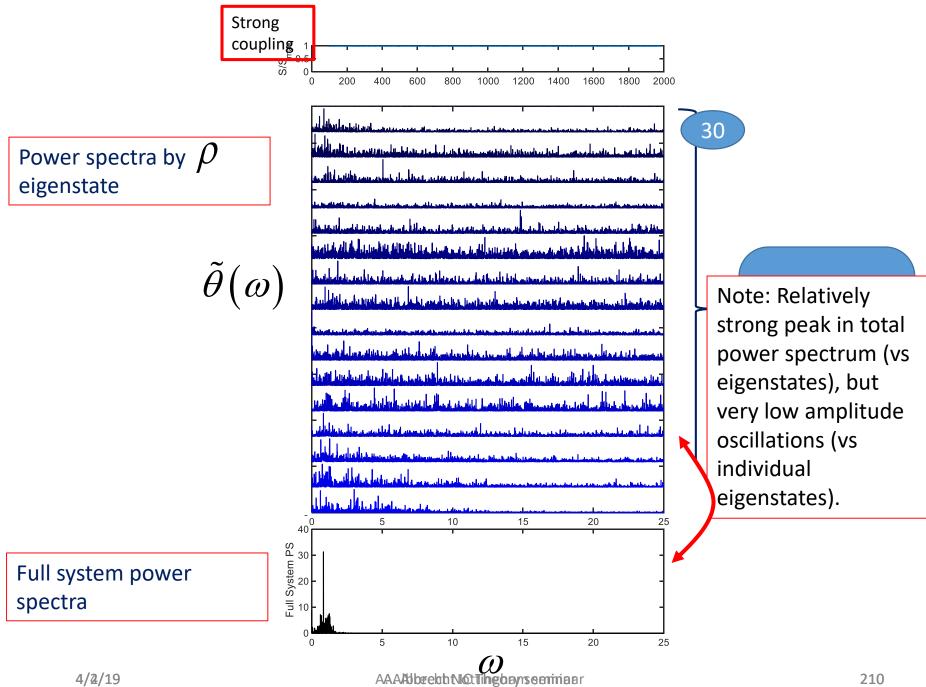




Now look at strong coupling, where Eqm seemed more noisy







Upshot:

- Strong oscillatory signal in <q> for ρ eigenstates for weakly coupled case (despite messy overall wavefunction shapes)
- → No such signal for strongly coupled case
- → Both cases show strong oscillatory signal for $\langle q\rho \rangle$ but amplitude is small.

Upshot:

Strong oscillatory signal in <q> for ρ eigenstates for weakly coupled case (despite messy overall wavefunction shapes)

→ No such signal for strongly coupled case

→ Both cases show strong oscillatory signal for $\langle q\rho \rangle$ but amplitude is small.

NEXT: A consistent histories approach

- → Generally, in the path integral formulation of QM interference among paths plays an important role
- → CH formalism identifies paths where interferences effects are NOT important. These are the paths to which probabilities can be assigned, and which are classical in that sense.
- → We have found that the messiness of the eqm physics of our toy model shows up as histories that degrade after a couple of SHO periods
- CH gives interesting test of "coherent state as most robust state" result from master equation work.

- → Generally, in the path integral formulation of QM interference among paths plays an important role
- → CH formalism identifies paths where interferences effects are NOT important. These are the paths to which probabilities can be assigned, and which are classical in that sense.
- → We have found that the messiness of the eqm physics of our toy model shows up as histories that degrade after a couple of SHO periods
- → CH gives interesting test of "coherent state as most robust state" result from master equation work.

- → Generally, in the path integral formulation of QM interference among paths plays an important role
- CH formalism identifies paths where interferences effects are NOT important. These are the paths to which probabilities can be assigned, and which are classical in that sense.
- → We have found that the messiness of the eqm a discrete time grid) toy model shows up as histories that degrade after a couple of SHO periods
- → CH gives interesting test of "coherent state as most robust state" result from master equation work.

(Define histories on

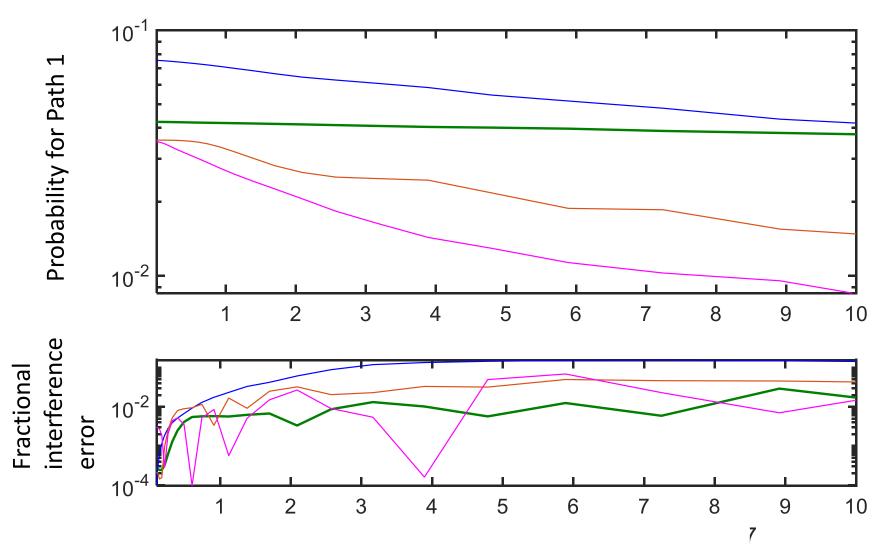
- → Generally, in the path integral formulation of QM interference among paths plays an important role
- → CH formalism identifies paths where interferences effects are NOT important. These are the paths to which probabilities can be assigned, and which are classical in that sense.
- → We have found that the messiness of the eqm physics of our toy model shows up as histories that degrade after a couple of SHO periods
- → CH gives interesting test of "coherent state as most robust state" result from master equation work.

Consistent histories (CH):

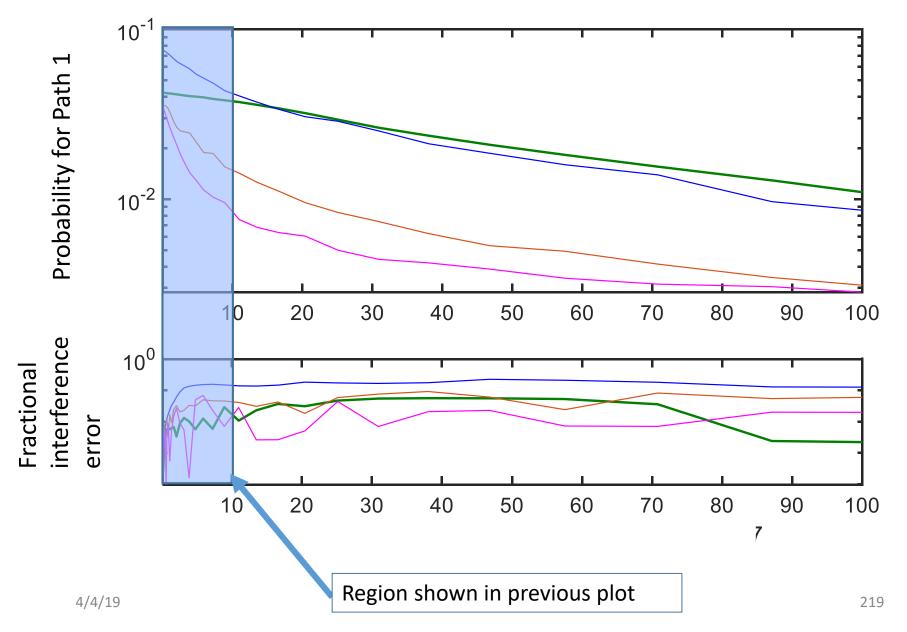
- → Generally, in the path integral formulation of QM interference among paths plays an important role
- → CH formalism identifies paths where interferences effects are NOT important. These are the paths to which probabilities can be assigned, and which are classical in that sense.
- → We have found that the messiness of the eqm physics of our toy model shows up as histories that degrade after a couple of SHO periods

→ <u>CH gives interesting test of "coherent state as most robust</u> <u>state" result from master equation work.</u>

Histories built from coherent states (green) degrade more slowly than histories built from other states.



Histories built from coherent states (green) degrade more slowly than histories built from other states. <u>But eventually the coherent state paths degrade too.</u>



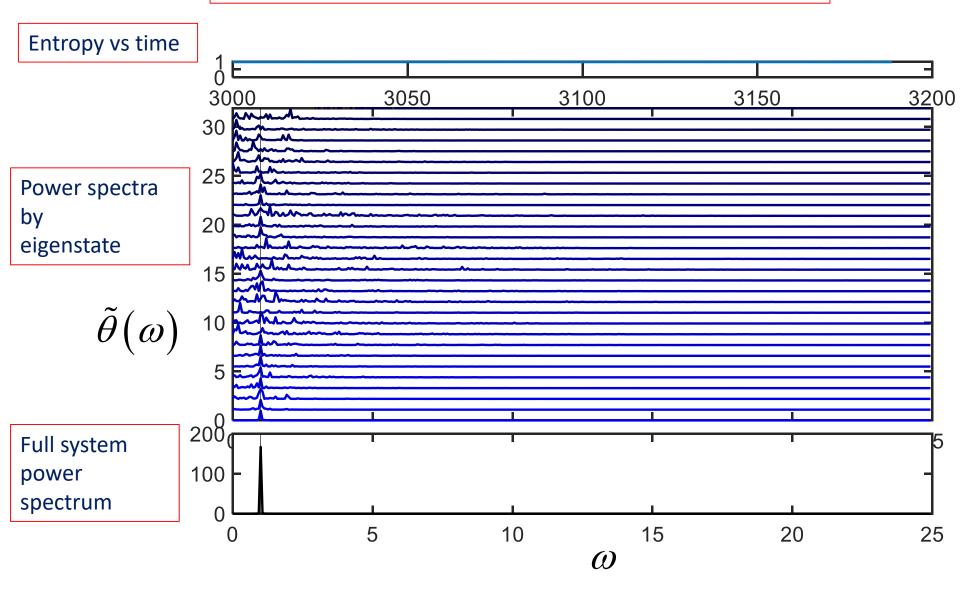
Lessons from eqm studies in the adapted CL model:

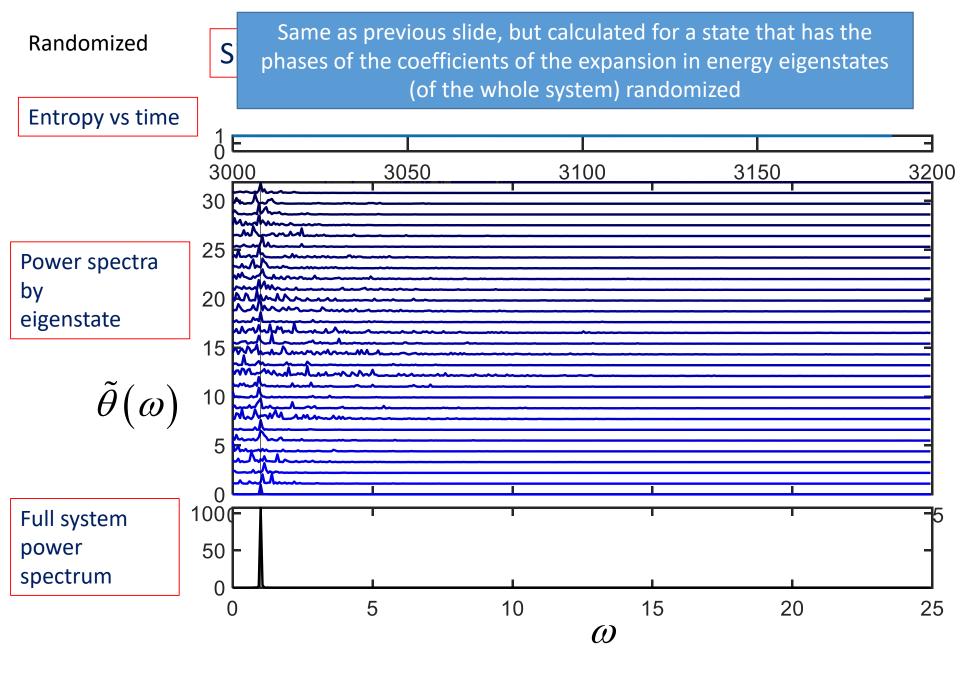
- Plenty of messiness in eqm ("wallowing or interfering Everett worlds")
- → Still, some intriguing sign of "classicality" show up in the power spectra of density matrix eigenstates.
- Consistent Histories formalism shows some classical behavior (which degrades after a couple of SHO periods)
- → Further discussion of implications at end of talk.

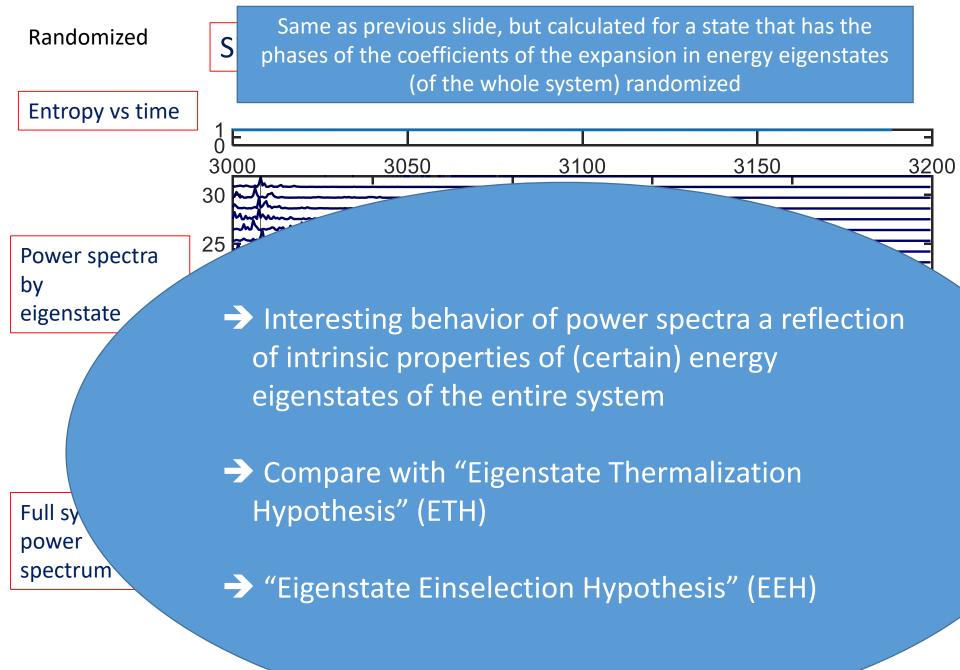
- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)

- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)

Similar power spectra to those shown above







- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)

- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)

- The adapted CL (ACL) model reproduced results about decoherence and einselection known from the standard CL model
- The ACL allows the exploration of these phenomena under conditions not accessible to the CL model (specifically eqm.)
- We have found einselection phenomena are highly degraded in eqm., but not completely destroyed.
- A suggestion: Classical phenomena may persist in equilibrium on a time scale short compared to a decoherence time... interesting for cosmology.

- The adapted CL (ACL) model reproduced results about decoherence and einselection known from the standard CL model
- The ACL allows the exploration of these phenomena under conditions not accessible to the CL model (specifically eqm.)
- We have found einselection phenomena are highly degraded in eqm., but not completely destroyed.
- A suggestion: Classical phenomena may persist in equilibrium on a time scale short compared to a decoherence time... interesting for cosmology.

- The adapted CL (ACL) model reproduced results about decoherence and einselection known from the standard CL model
- The ACL allows the exploration of these phenomena under conditions not accessible to the CL model (specifically eqm.)
- We have found einselection phenomena are highly degraded in eqm., but not completely destroyed.
- A suggestion: Classical phenomena may persist in equilibrium on a time scale short compared to a decoherence time... interesting for cosmology.

- The adapted CL (ACL) model reproduced results about decoherence and einselection known from the standard CL model
- The ACL allows the exploration of these phenomena under conditions not accessible to the CL model (specifically eqm.)
- We have found einselection phenomena are highly degraded in eqm., but not completely destroyed.
 Interesting connection with

Interesting connection with cosmological motivations

A suggestion: Classical phenomena may persist in equilibrium on a time scale short compared to a decoherence time... interesting for cosmology.

- The adapted CL (ACL) model reproduced results about decoherence and einselection known from the standard CL model
- → The ACL allows the exploration of these phenomena under conditions not accessible to the CL model (specifically eqm.)
- We have found einselection phenomena are highly degraded in eqm., but not completely destroyed.
- A suggestion: Classical phenomena may persist in equilibrium on a time scale short compared to a decoherence time... interesting for cosmology.

- The adapted CL (ACL) model reproduced results about decoherence and einselection known from the standard CL model
- The ACL allows the exploration of these phenomena under conditions not accessible to the CL model (specifically eqm.)
- We have found einselection phenomena are highly degraded in eqm., but not completely destroyed.
- A suggestion: Classical phenomena may persist in equilibrium on a time scale short compared to a decoherence time... interesting for cosmology.

"classical" features emerged in eqm when the decoherence time was extended (weak coupling) to be longer than the SHO period

- The adapted CL (ACL) model reproduced results about decoherence and einselection known from the standard CL model
- The ACL allows the exploration of these phenomena under conditions not accessible to the CL model (specifically eqm.)
- We have found einselection phenomena are highly degraded in eqm., but not completely destroyed.
- A suggestion: Classical phenomena may persist time scale short compared to a decoherence ti cosmology.

"classical" features emerged in eqm when the decoherence time was extended (weak coupling) to be longer than the SHO period Compare with "de Sitter Equilibrium" scenario where decoherence time is the age of the Universe