

262N1.0 (P1)

~~(2) II-1~~

II Dimensional Analysis

i) Dimensional analysis handout. (Copies available on the course web site)

ii) Estimating eqm in the universe

Compare interaction rates with the expansion rate.

$$\Gamma = n \sigma |v| = \text{interaction (or reaction) rate}$$

↑ ↑
number cross-section
density of
relevant particles

H = expansion rate

Case 1: Everything relativistic:

Only relevant energy scale in Γ is T
using energy units, and dimensional analysis

$$\left. \begin{array}{l} n \sim T^3, \quad \sigma \sim T^{-2} \\ |v| \approx 1 \end{array} \right\} \Gamma \sim T$$

~~(2) II-2~~
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$$H^2 = \frac{8\pi}{3} G \rho$$

(in early universe can neglect Λ and $\frac{K}{R_0 a^2}$)

$$\rho \sim T^4$$

$$\text{write } G = m_p^{-2} \quad (m_p \approx 10^{19} \text{ GeV})$$

$$H^2 \approx m_p^{-2} T^4$$

When is $\Gamma \gg H$?

$$T \gg m_p^{-1} T^2$$

$$m_p \gg T$$

so eqm achieved for $T \ll m_p \approx 10^{19} \text{ GeV}$

$\Rightarrow$ eqm possible starting in the very early universe.

~~② II 3~~

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Case 2: "Four-Fermi" form for σ

There are many other possibilities than "Case 1", where everything was relativistic. Here is ~~an~~ one example that illustrates an important ~~point~~ point.

Some interactions are mediated by a vector "Boson" with mass M_x . In the limit where all interacting particles are relativistic but $T \ll M_x$, M_x appears with a "-4" power in σ .

to get the units right

$$\sigma \propto T^n M_x^{-4} \rightarrow n = 2$$

$$\Gamma \approx \frac{T^5}{M_x^4}$$

$$\Gamma \gg H \Rightarrow \frac{T^5}{M_x^4} \gg \frac{T^2}{m_p}$$

$$T \gg m_x \left(\frac{m_x}{m_p} \right)^{1/3}$$

for example, if $M_x \approx 100 \text{ GeV}$ ~~fixed~~
 (as it is for "weak interaction")

$$T \gg 10^{-4} \text{ GeV} = 100 \text{ MeV}$$

The point of choosing this particular "case 2" was to show that for some processes eqm is achieved for temperatures above some critical ~~very~~ value. The processes drop out of eqm as the expanding universe drops below the critical value.

iii) Properties of a freely propagating thermal distribution of photons.

Eq 2.25 from ~~Reg~~ Reydin:

$$E(f) df = \frac{8\pi h}{c^3} \frac{f^3 df}{\exp(hf/kT) - 1}$$

= ~~number~~ ^{energy} density of photons with frequencies between f and $f+df$.

M=0

Particles in eqm:

Phase space distribution function

$$f(\vec{p}, t_1) = [\exp(\frac{E}{T_1}) \pm 1]^{-1}$$

$$f(\vec{p}) \cdot dV \cdot d\vec{p} = \# \text{ of particles in vol } dV \text{ with } p < p < p + dp$$

Follow For non-interacting particles:

Follow vol dV and momentum between t_1 + t_2 (assume eqm at t_1)For a photon traveling in an FRW universe
 $p = \frac{1}{\lambda} \sim \frac{1}{a}$ (see p 40 in Ryden)

→ If we ~~follow~~ allow ΔV to evolve in a ~~comoving~~ comoving way, # of particles is ~~the~~ unchanged, as long as we follow the momenta as they redshift.

$$\Delta V_1 \rightarrow \Delta V_2 = \Delta V_1 \left(\frac{a_2}{a_1} \right)^3$$

$$p_1 \rightarrow p_2 = p_1 \left(\frac{a_1}{a_2} \right)$$

$$f(\vec{p}_2, t_2) dV_2 d^3p_2 = f(\vec{p}_1, t_1) dV_1 d^3p_1$$

$$f(p_2, t_2) = f(\vec{p}_1, t_1) \frac{dV_1}{dV_2} \frac{d^3p_1}{d^3p_2} = f(\vec{p}_1, t_1)$$

$$f(p_2, t_2) = \left[\exp\left(\frac{p_1}{T_1}\right) \pm 1 \right]^{-1} = \left[\exp\left(\frac{p_2 \left(\frac{a_1}{a_2}\right)}{T_1} \right) \pm 1 \right]^{-1}$$

$$T_2 = T_1 \left(\frac{a_1}{a_2}\right)$$

$$f(p_2, t_2) = \left[\exp\left(\frac{p_2}{T_2}\right) \pm 1 \right]^{-1}$$