

# Les Houches Lectures on Cosmic Inflation

## Four Parts

- 1) Introductory material
- 2) Entropy, Tuning and Equilibrium in Cosmology
- 3) Classical and quantum probabilities in the multiverse
- 4) de Sitter equilibrium cosmology

Andreas Albrecht; UC Davis  
Les Houches Lectures; July-Aug 2013

## Part 4 Outline

1. de Sitter Equilibrium cosmology
2. Cosmic curvature from de Sitter Equilibrium cosmology

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## de Sitter Equilibrium (dSE) cosmology

- Take ideas from Holography,  $\Lambda$  to construct a finite cosmology

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## de Sitter Equilibrium (dSE) cosmology

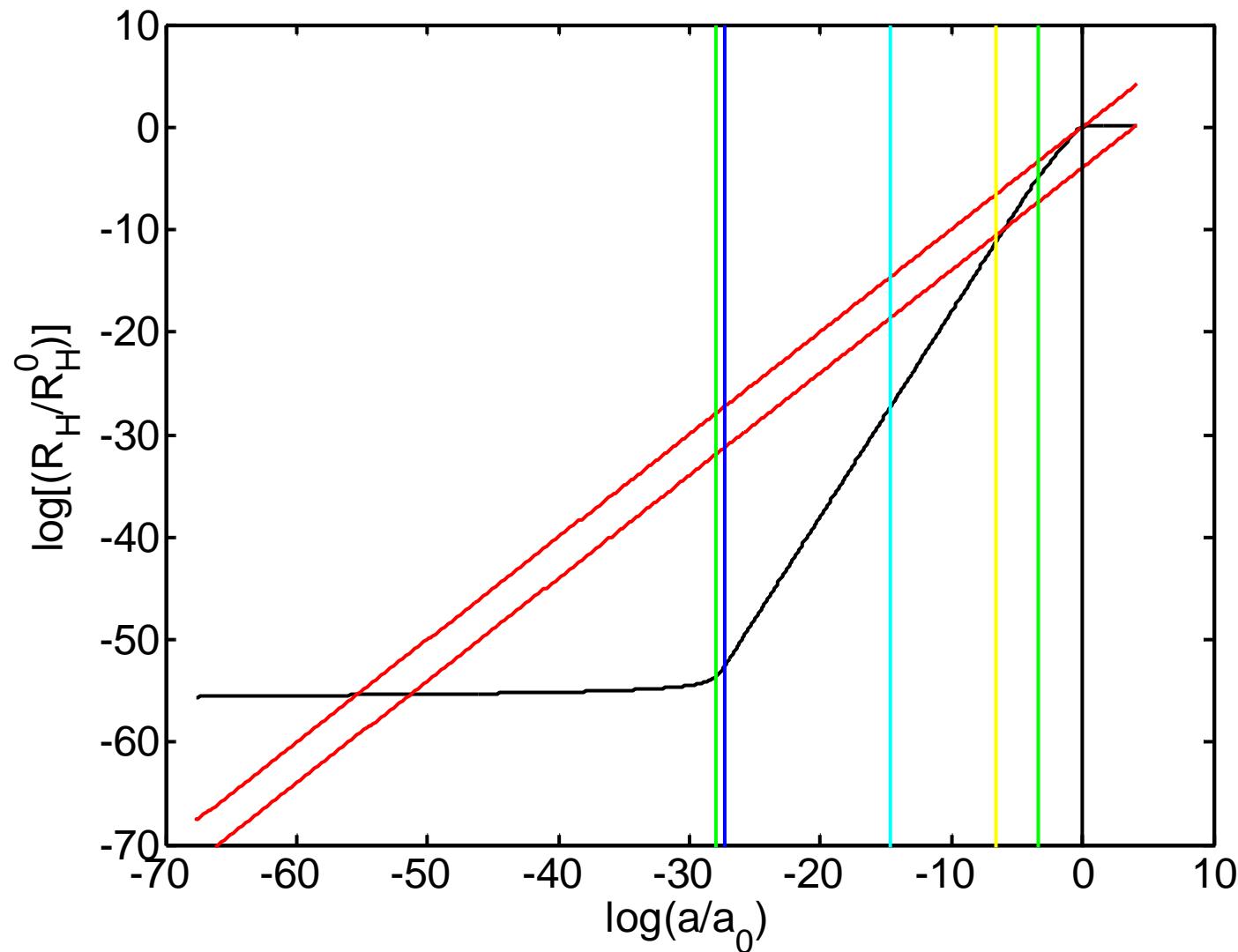
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AA: *arXiv:1104.3315*

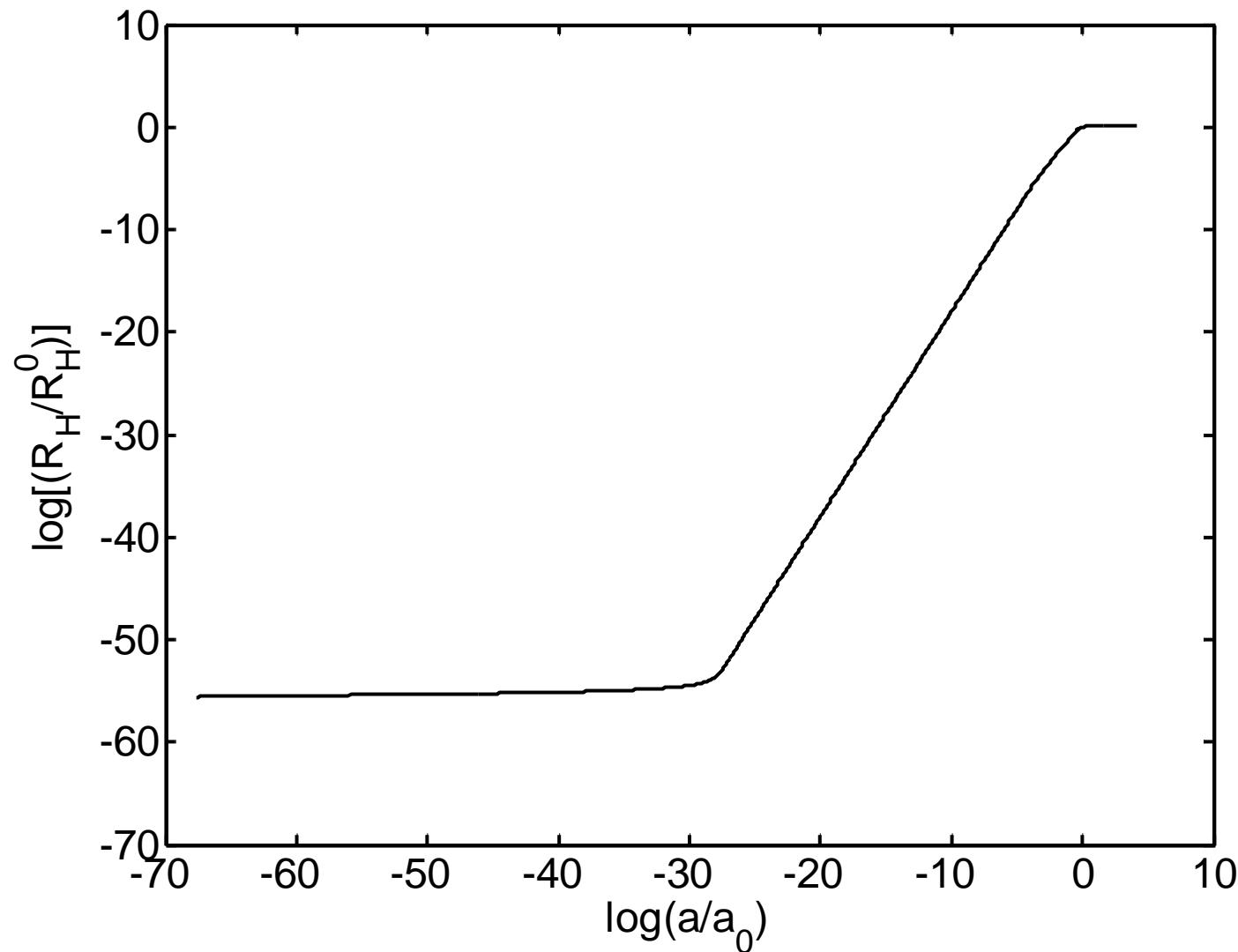
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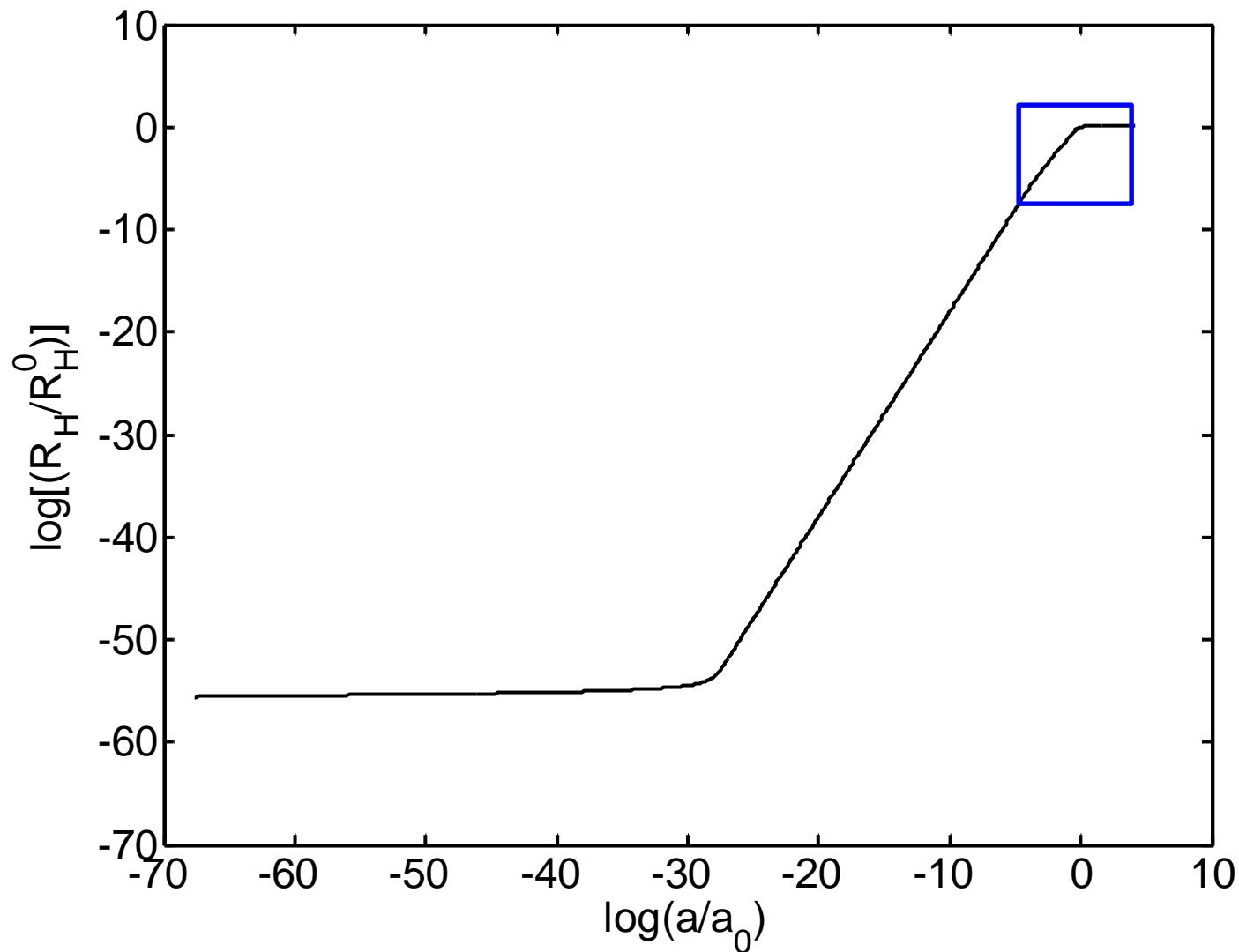
AA & Sorbo: *hep-th/0405270*

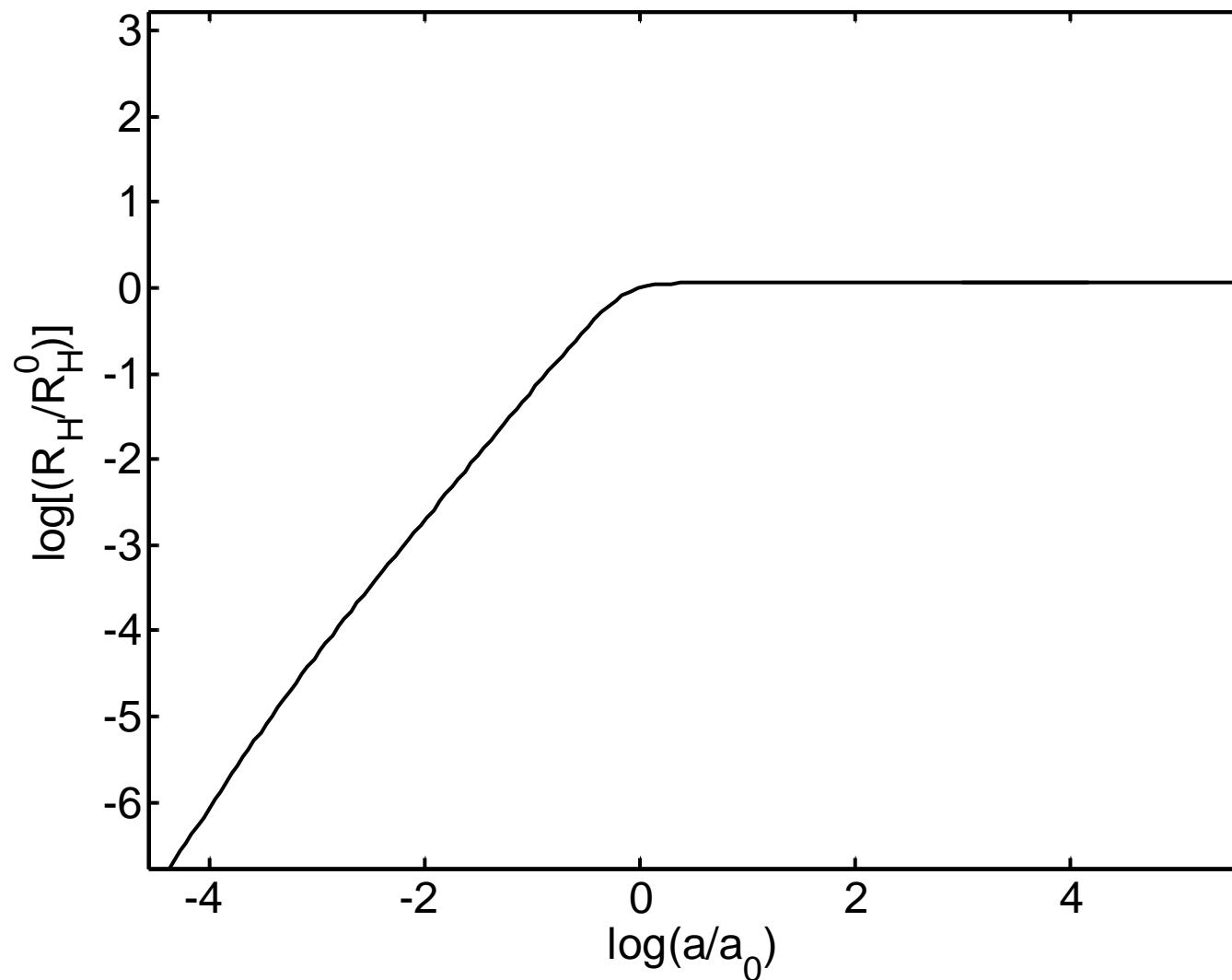
# Evolution of Cosmic Length



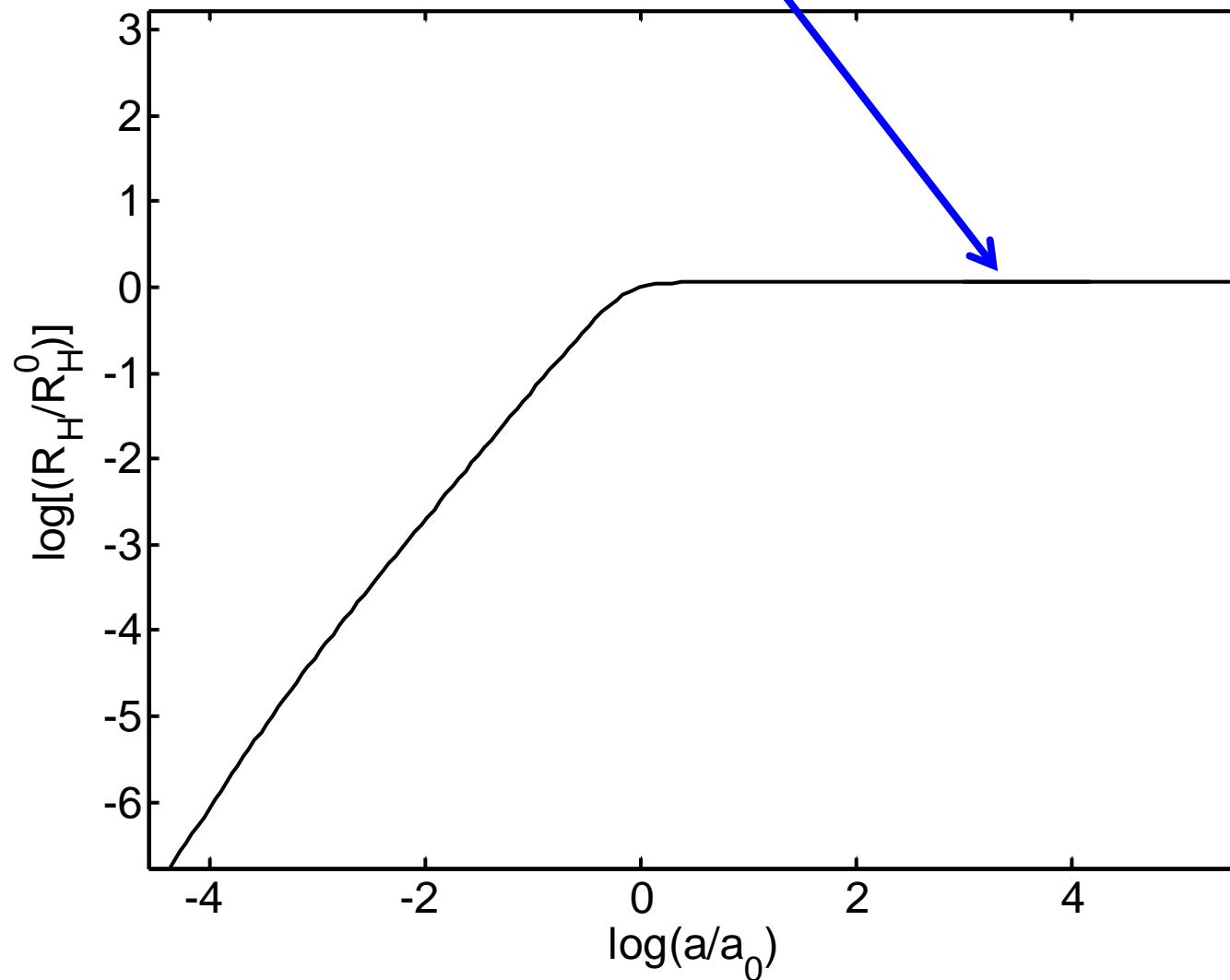
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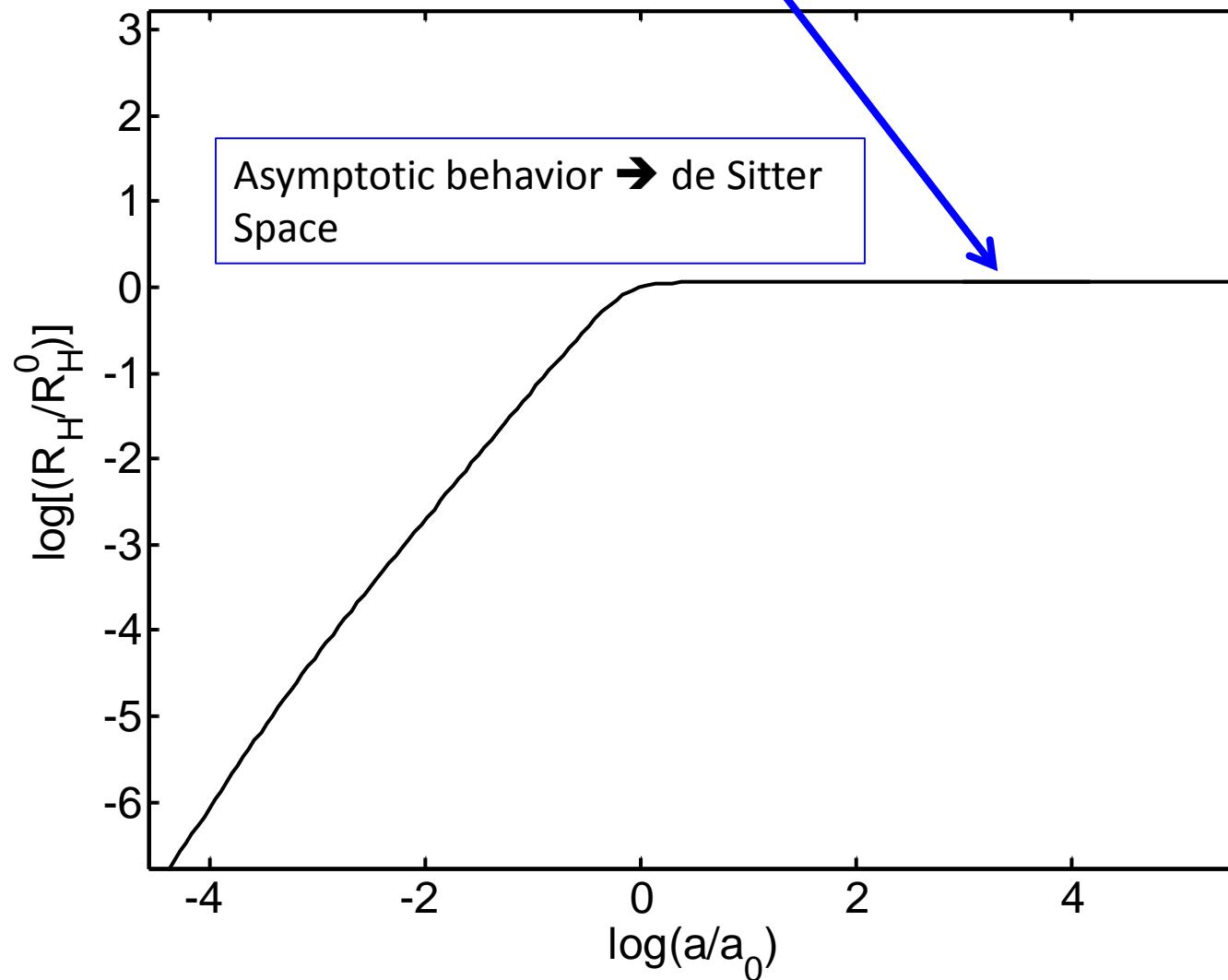




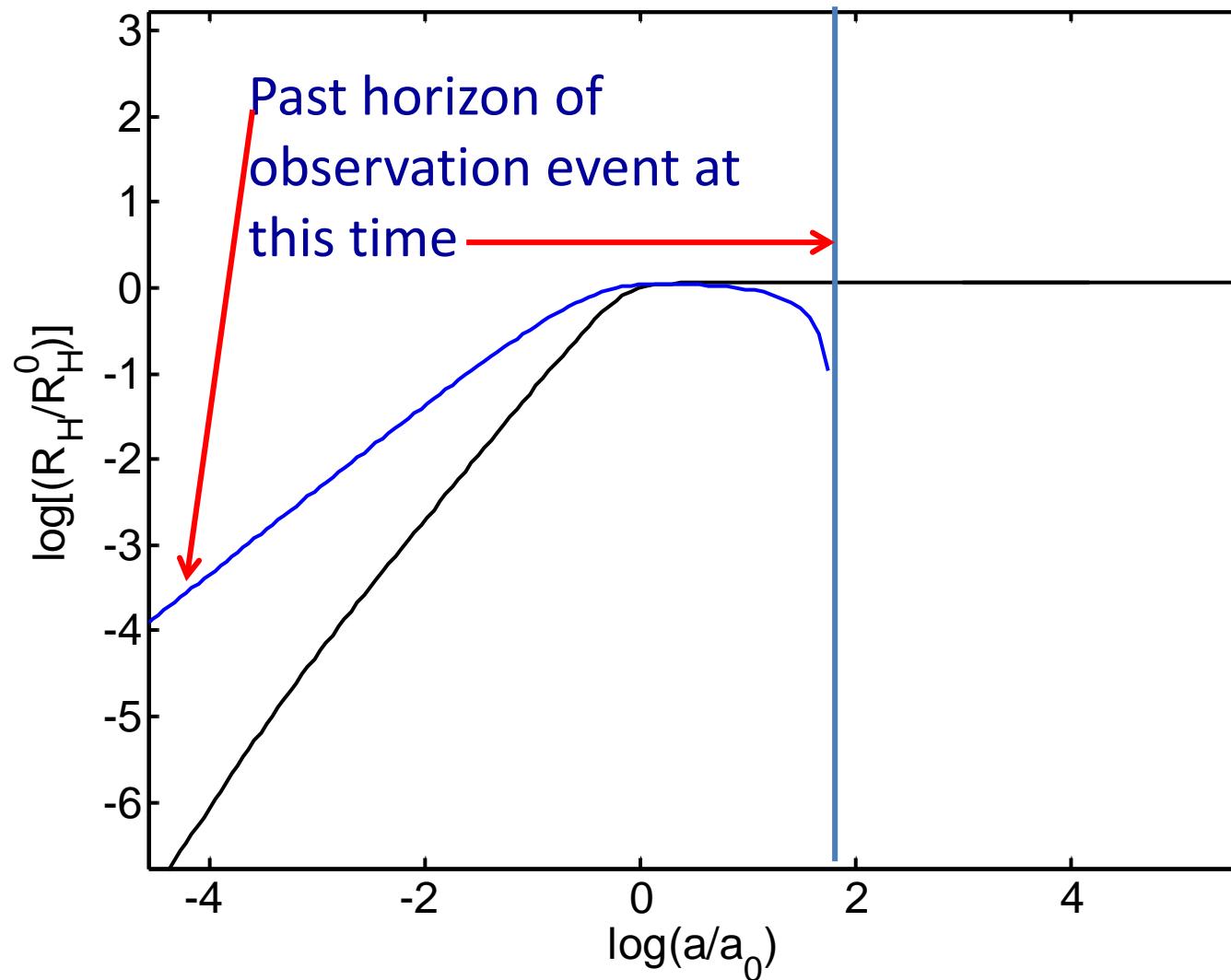
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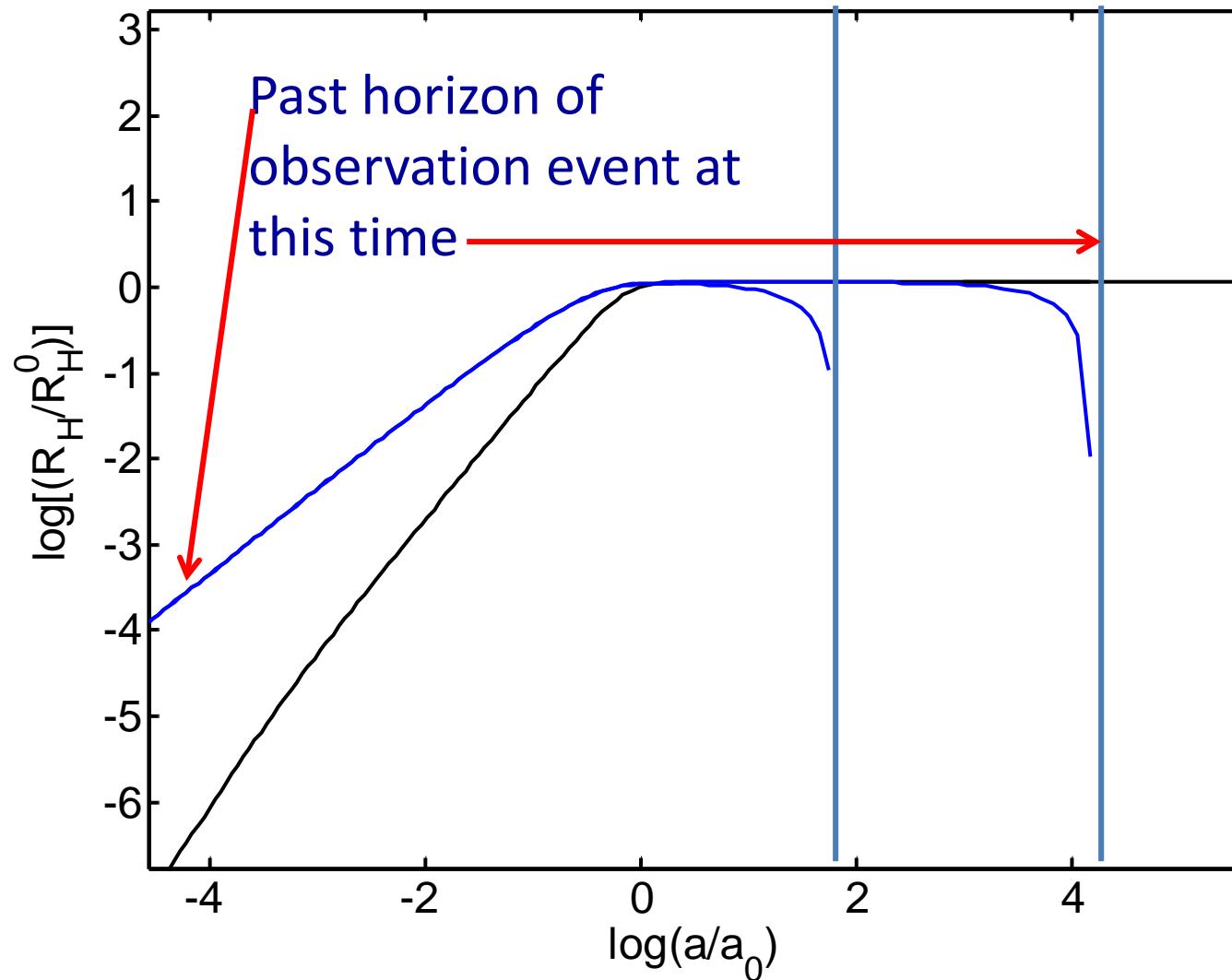


# The de Sitter horizon



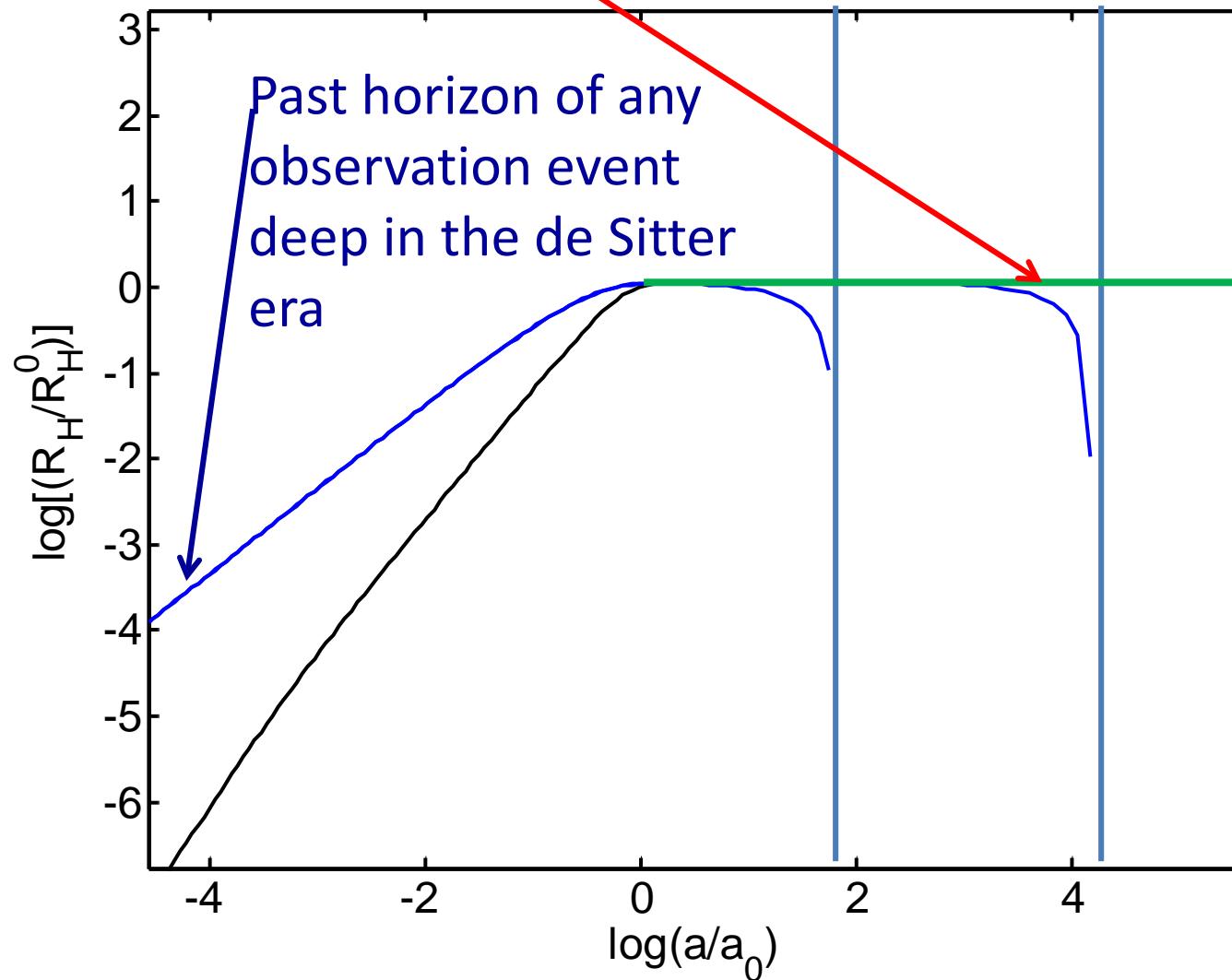
*Past Horizon: Physical distance from (comoving) observer of a photon that will reach the observer at the time of the observation.*

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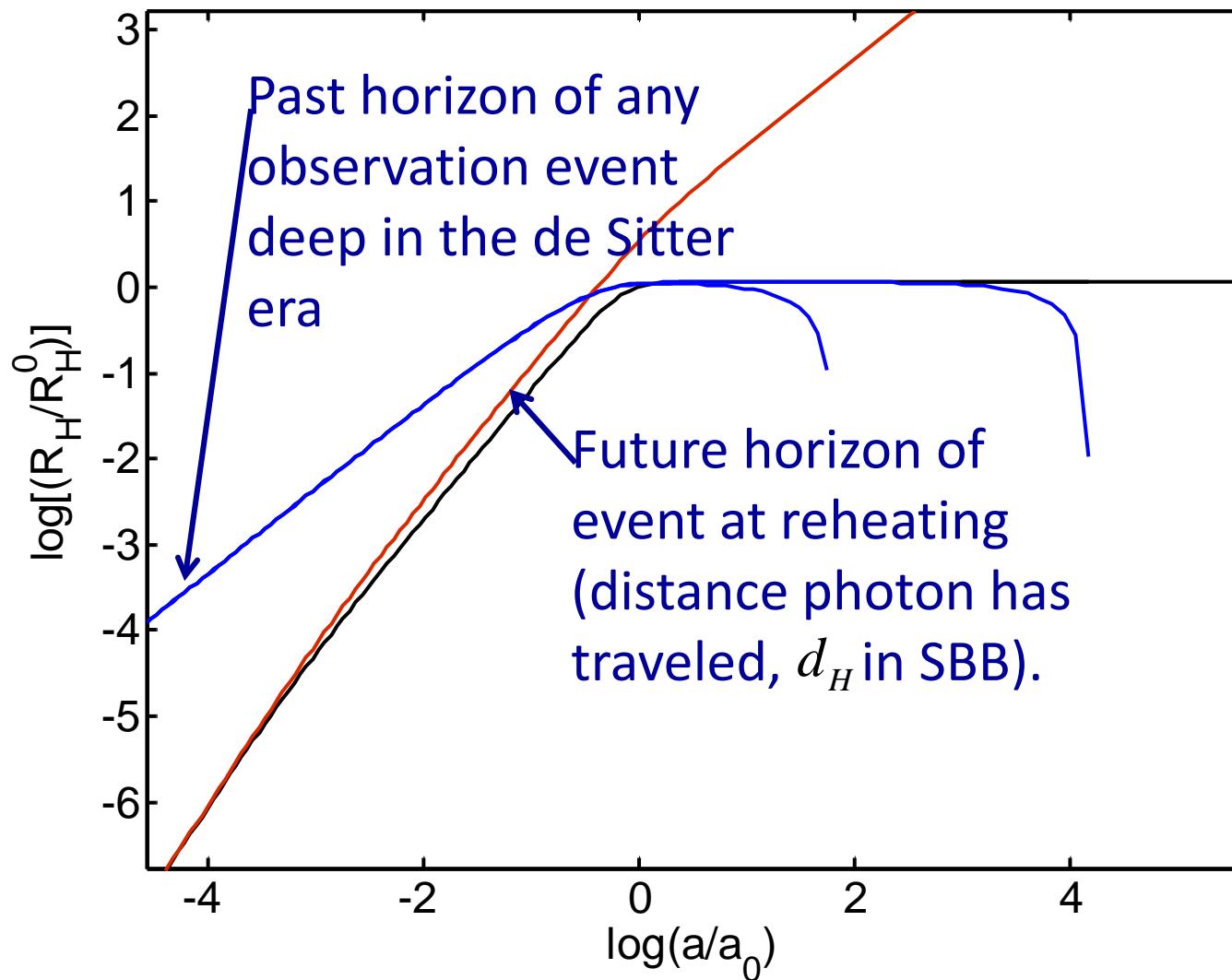


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# Implications of the de Sitter horizon

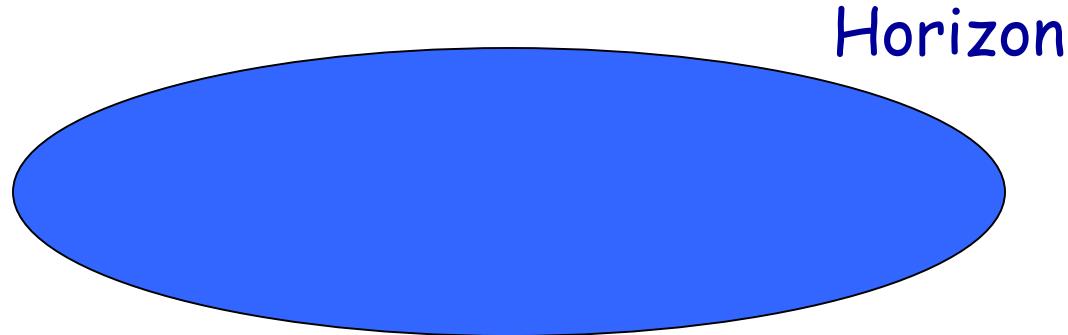
- Maximum entropy

$$S_{\Lambda} \propto A = H_{\Lambda}^{-2} = \left( \frac{\Lambda}{3} \right)^{-1}$$

- Gibbons-Hawking Temperature

$$T_{GH} = H_{\Lambda} = \sqrt{\frac{8\pi G}{3}} \rho_{\Lambda}$$

# "De Sitter Space: The ultimate equilibrium for the universe?



$$S \propto A = H^{-2} = \Lambda^{-1}$$

$$T_{GH} = H_\Lambda = \sqrt{\frac{8\pi G}{3}} \rho_\Lambda$$

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- Only a finite volume ever observed

- If  $\Lambda$  is truly constant: Cosmology as fluctuating Eqm.

- Maximum entropy  $\longrightarrow$  finite Hilbert space of dimension  $N = e^{S_\Lambda}$

Banks & Fischler & Dyson et al.

# Implications of the de Sitter horizon

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dSE cosmology

- ✓• If  $\Lambda$  is truly constant: Cosmology as fluctuating Eqm.?

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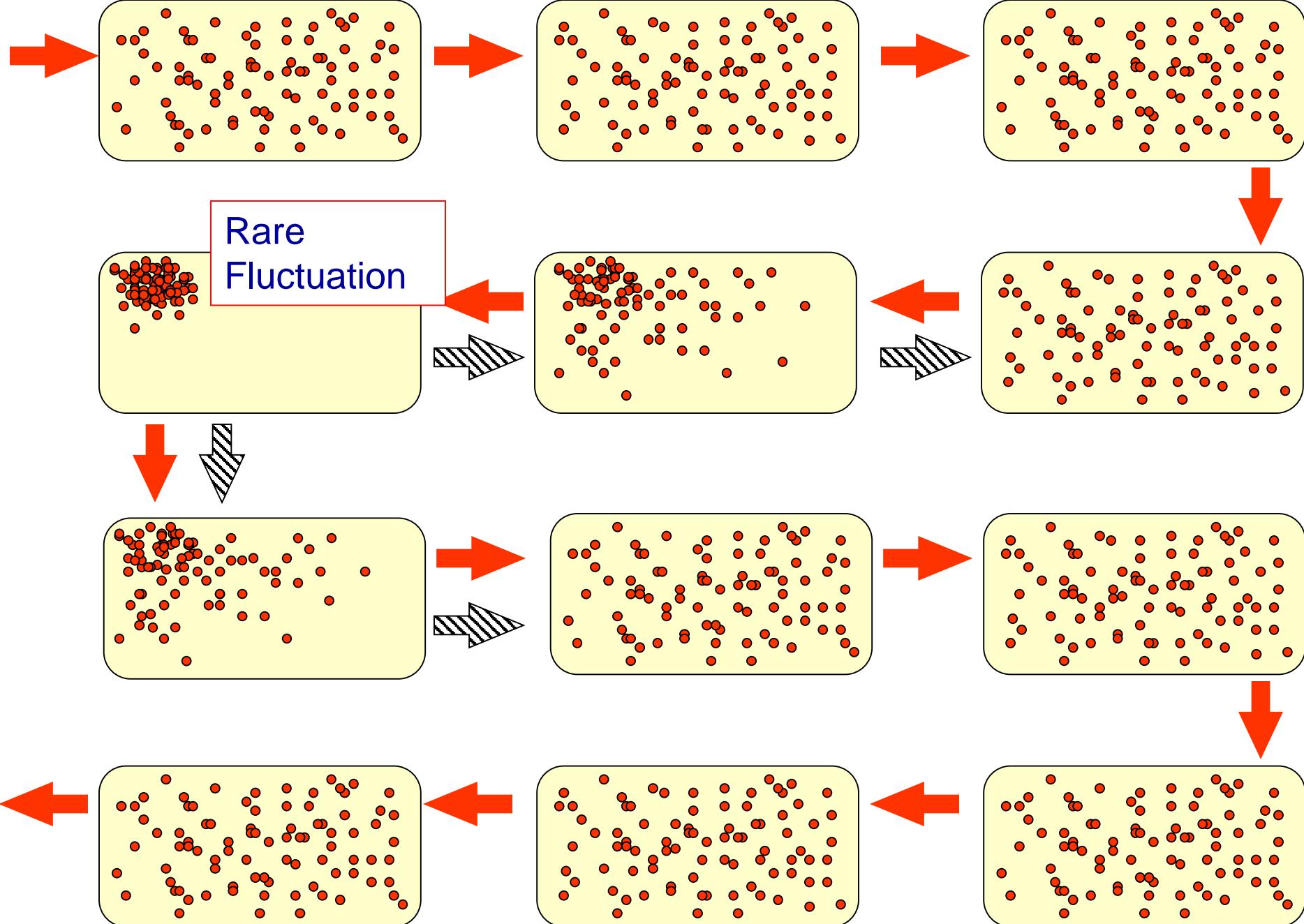
# Equilibrium Cosmology

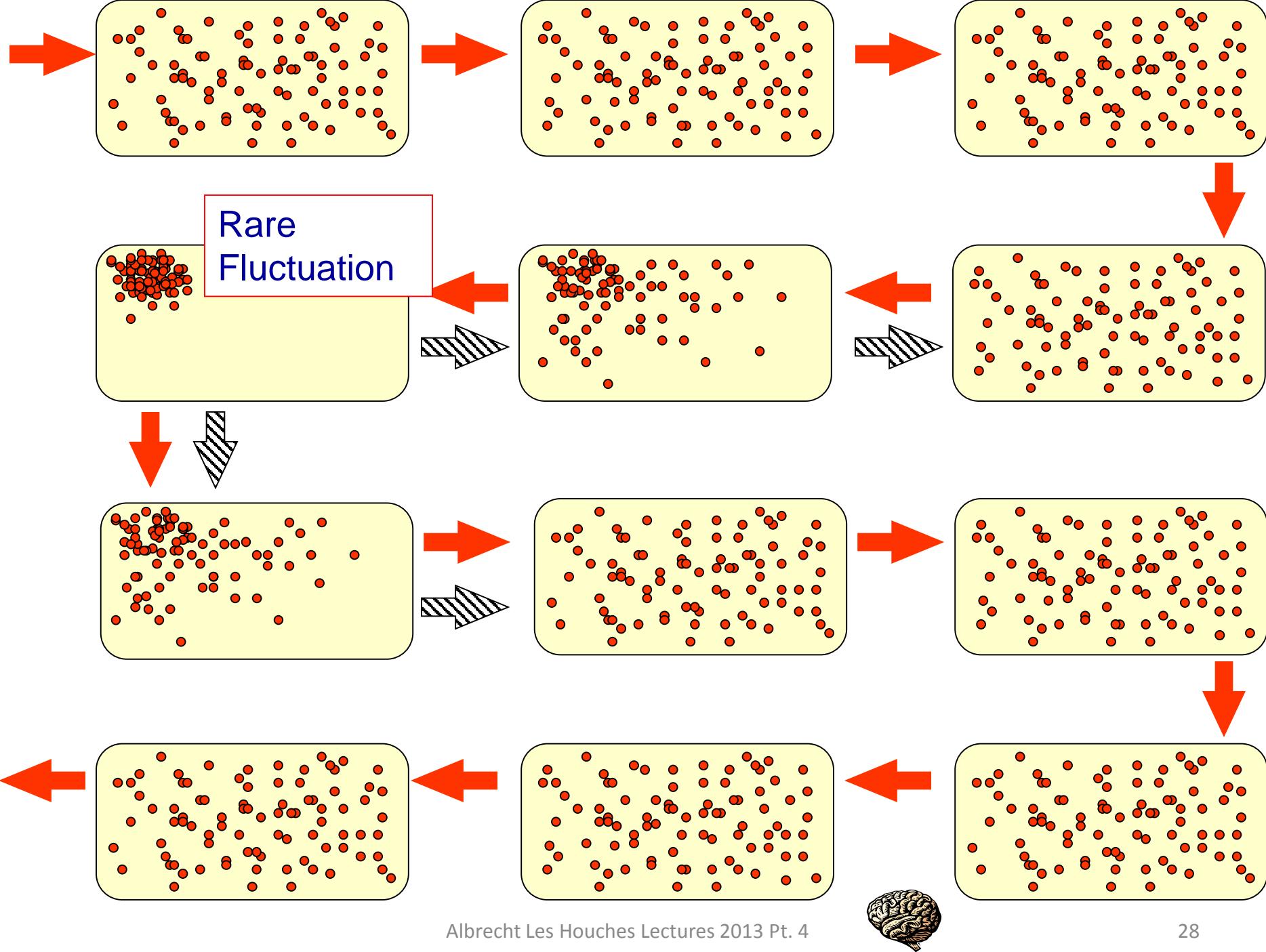
## Equilibrium Cosmology

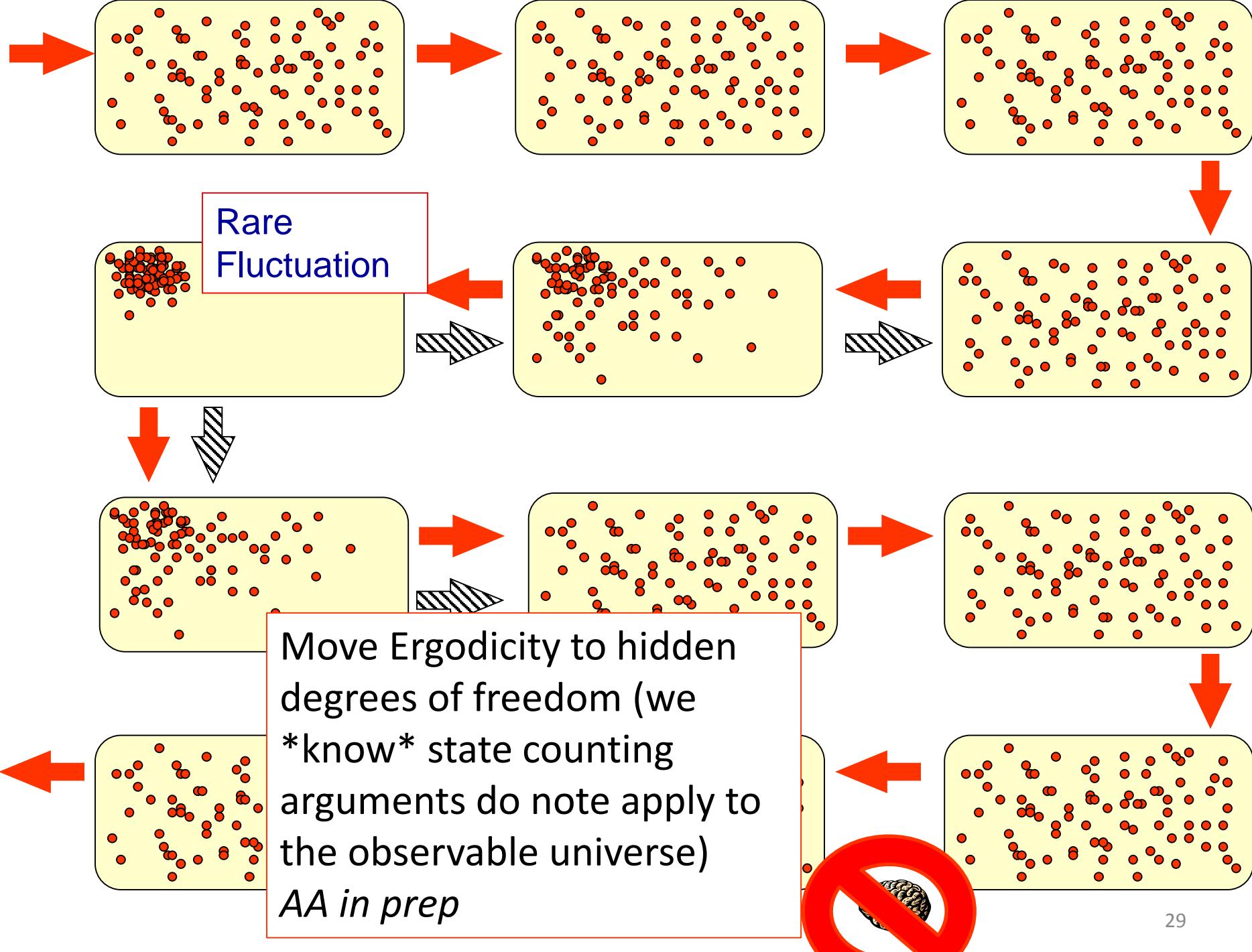
An eqm. theory does not require any theory of initial conditions. The probability of appearing in a given state is given entirely by stat mech, and is thus “given by the dynamics”.

If you know the Hamiltonian you know how to assign probabilities to different states without any special theory of initial conditions.

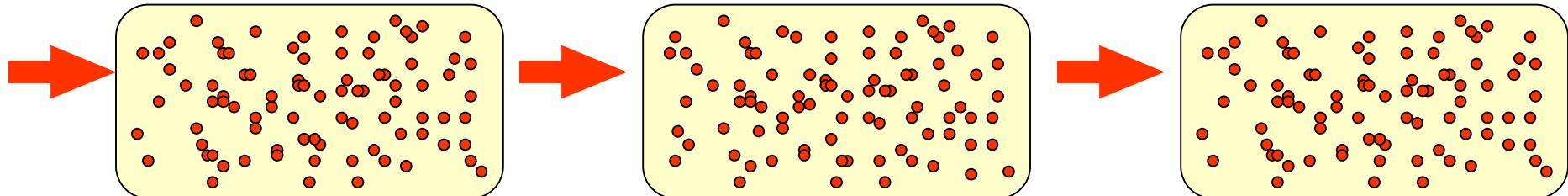
Dyson et al 2002



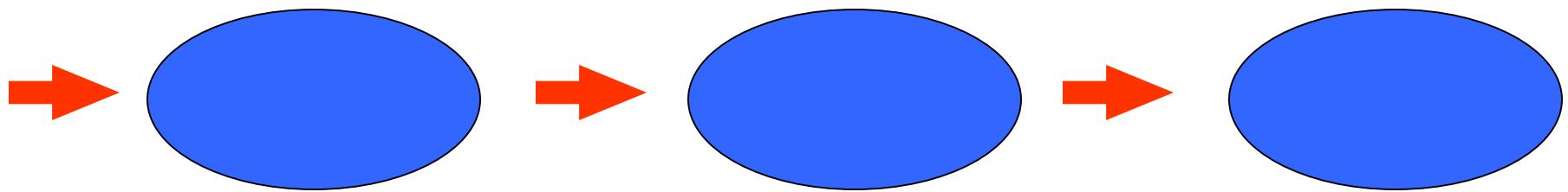




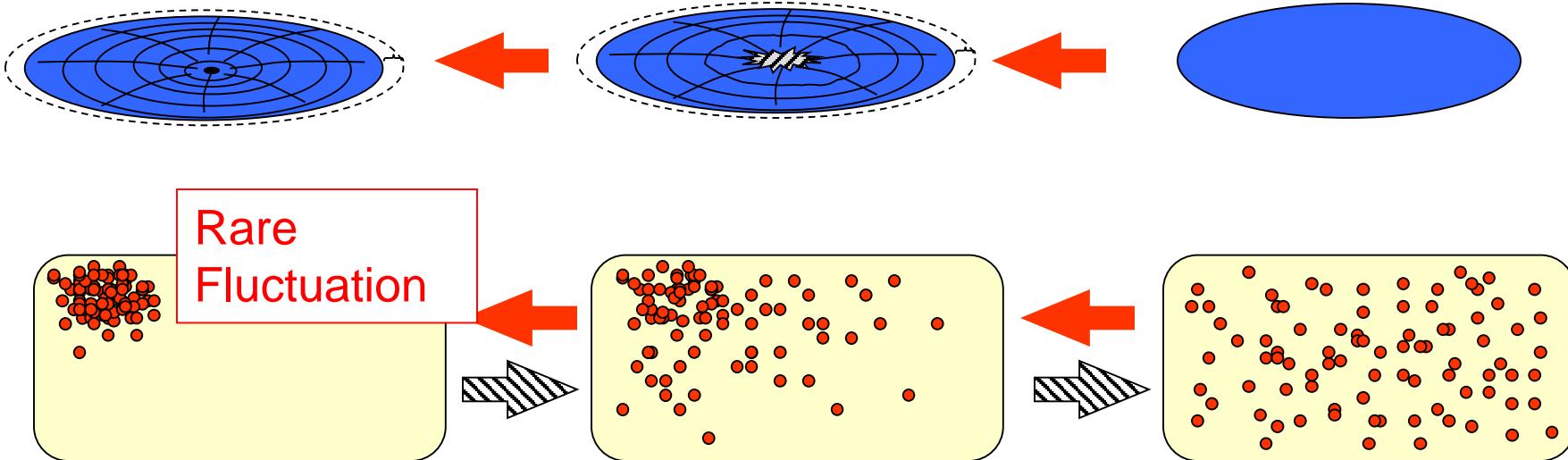
Concept:



Realization:



“de Sitter Space”



## Fluctuating from dSE to inflation:

- The process of an inflaton fluctuating from late time de Sitter to an inflating state is dominated by the “Farhi-Guth Guven” (FGG) process
- A “seed” is formed from the Gibbons-Hawking radiation that can then tunnel via the Guth-Farhi instanton.

- Rate is well approximated by the rate of seed formation:

$$\propto e^{-\frac{m_s}{T_{GH}}} = e^{-\frac{m_s}{H_\Lambda}}$$

- Seed mass:

$$m_s = \rho_I \left( c H_I^{-1} \right)^3 = 0.0013 \text{kg} \left( \frac{\left( 10^{16} \text{GeV} \right)^4}{\rho_I} \right)^{1/2}$$

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Small seed can produce an entire universe →  
Evade “Boltzmann Brain” problem

Albrecht Les Houches Lectures 2013 Pt. 4

## Fluctuating from dSE to inflation:

- The process of an inflaton fluctuating from late time de Sitter to an inflating space-time

**“Farhi-Guth Guven” (FGG)**

$M \rightarrow 0$  not a problem for G-F process (A. Ulvestad & AA 2012)

- A “seed” is formed from the Gibbons Hawking radiation that can then trigger an instanton.

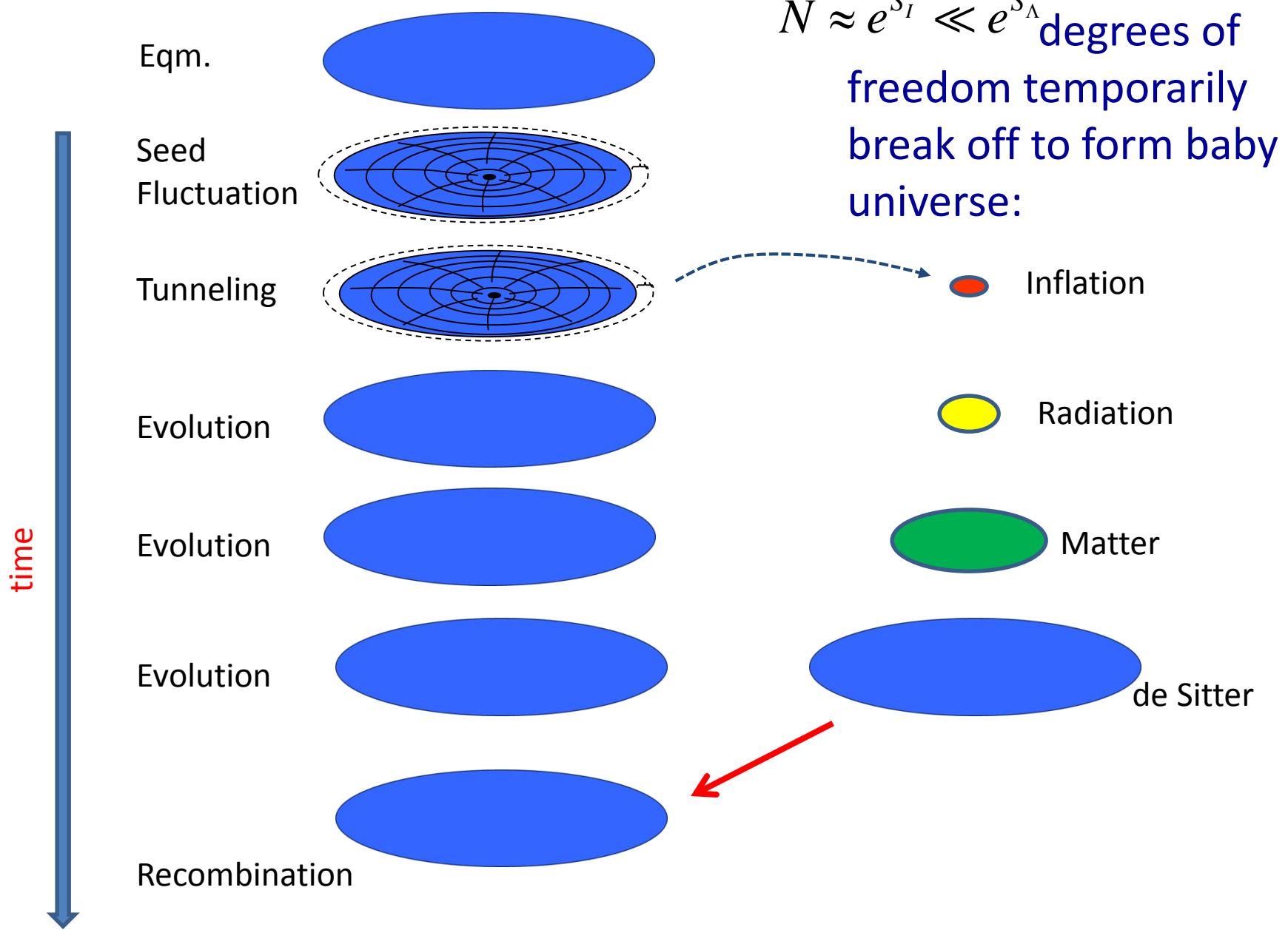
See also Freivogel et al 2006,  
Banks 2002

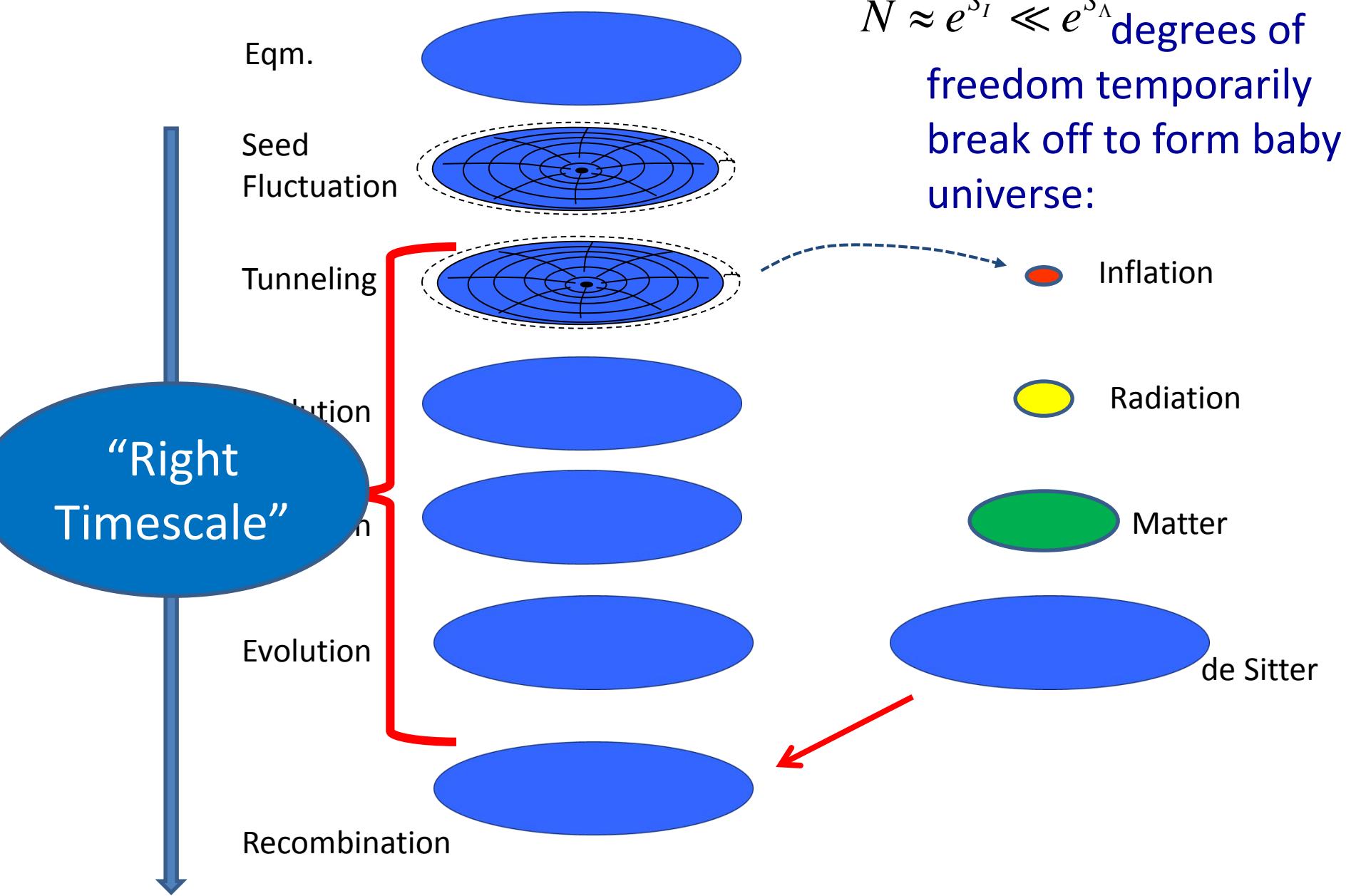
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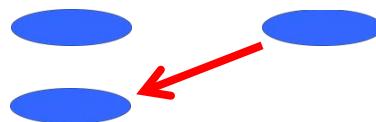


## Implications of finite Hilbert space $N = e^{S_\Lambda}$

- Recurrences
- Eqm.
- Breakdown of continuum field theory

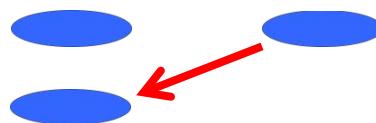
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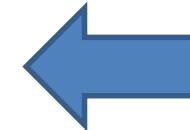
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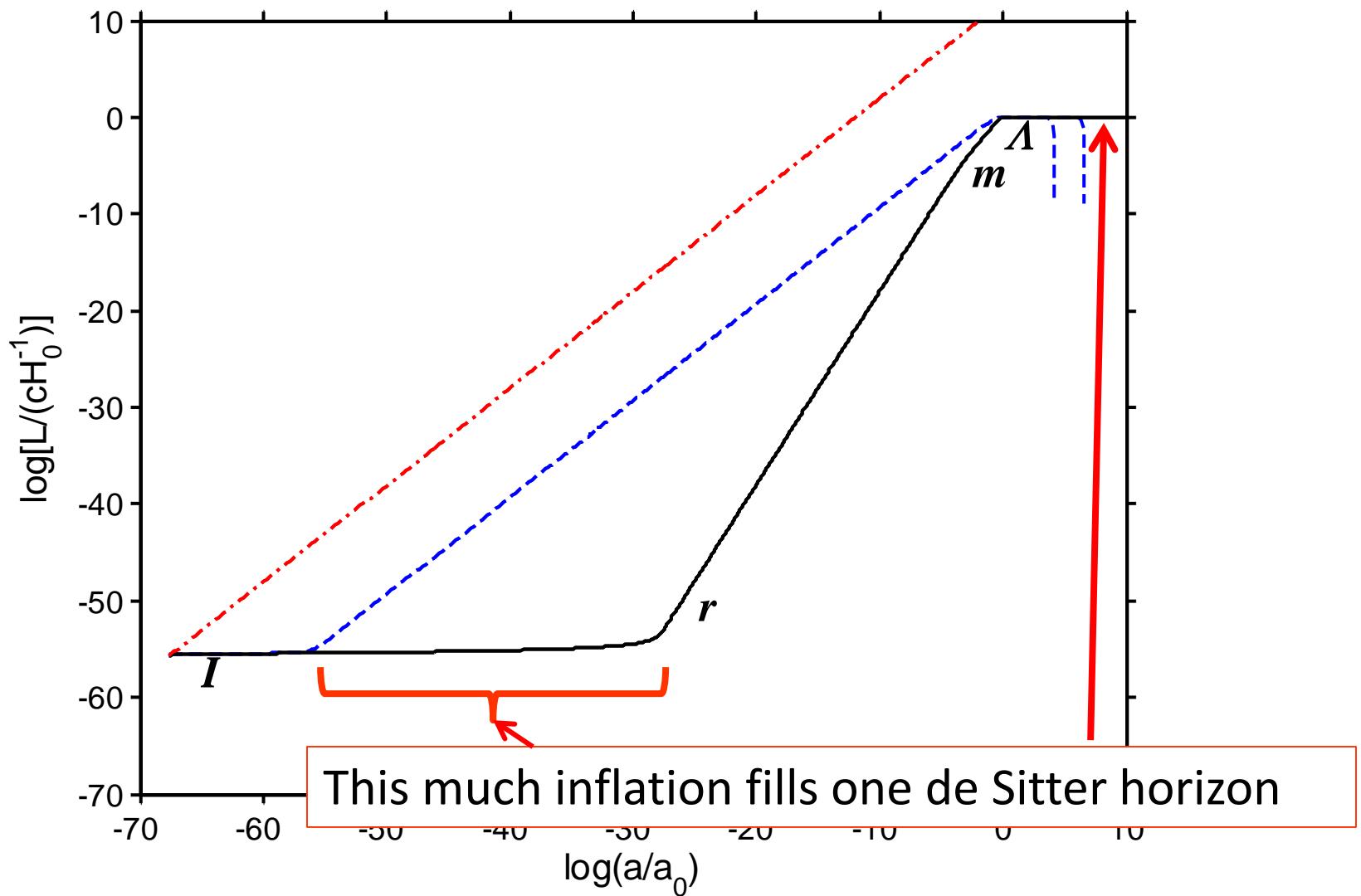
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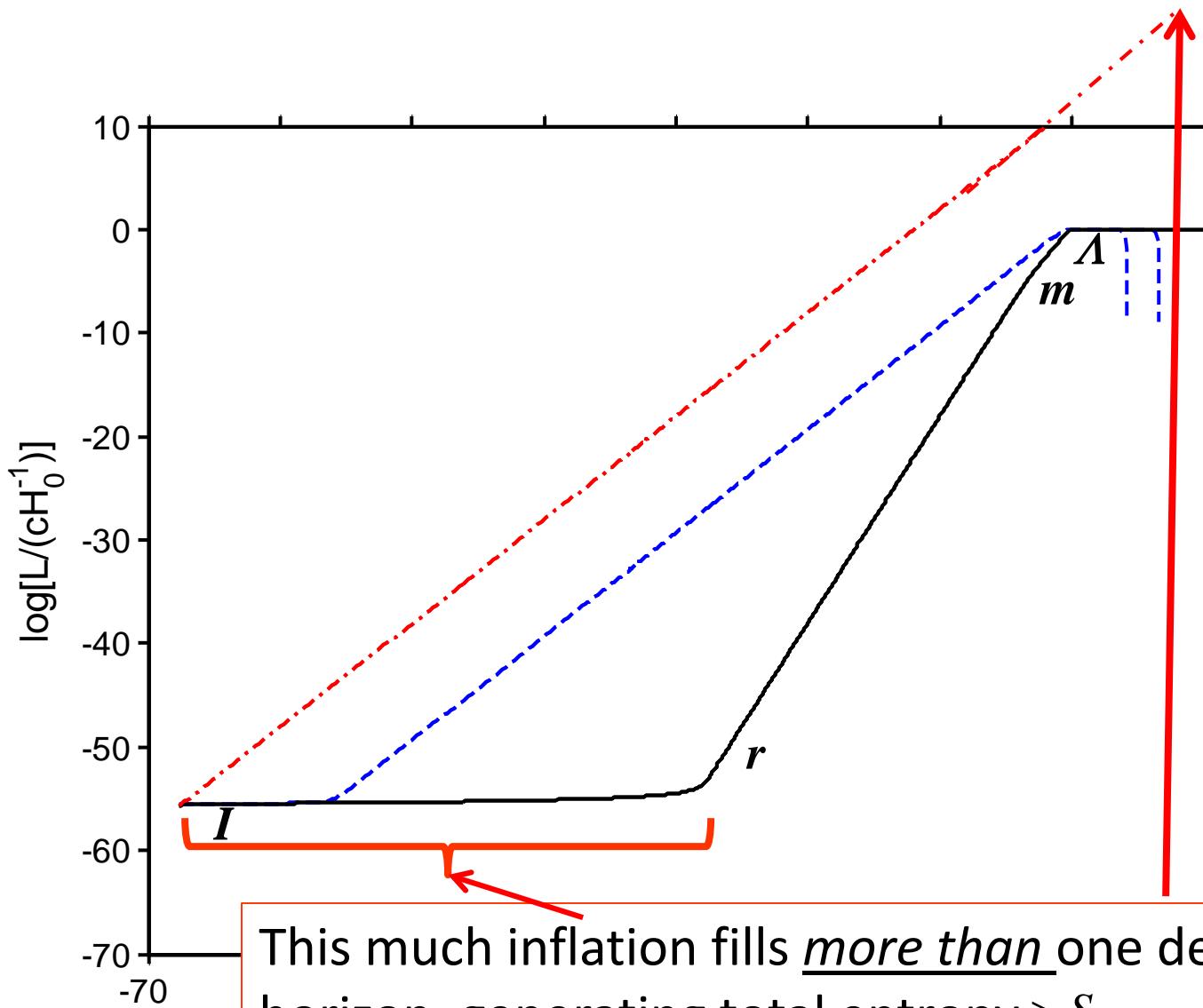


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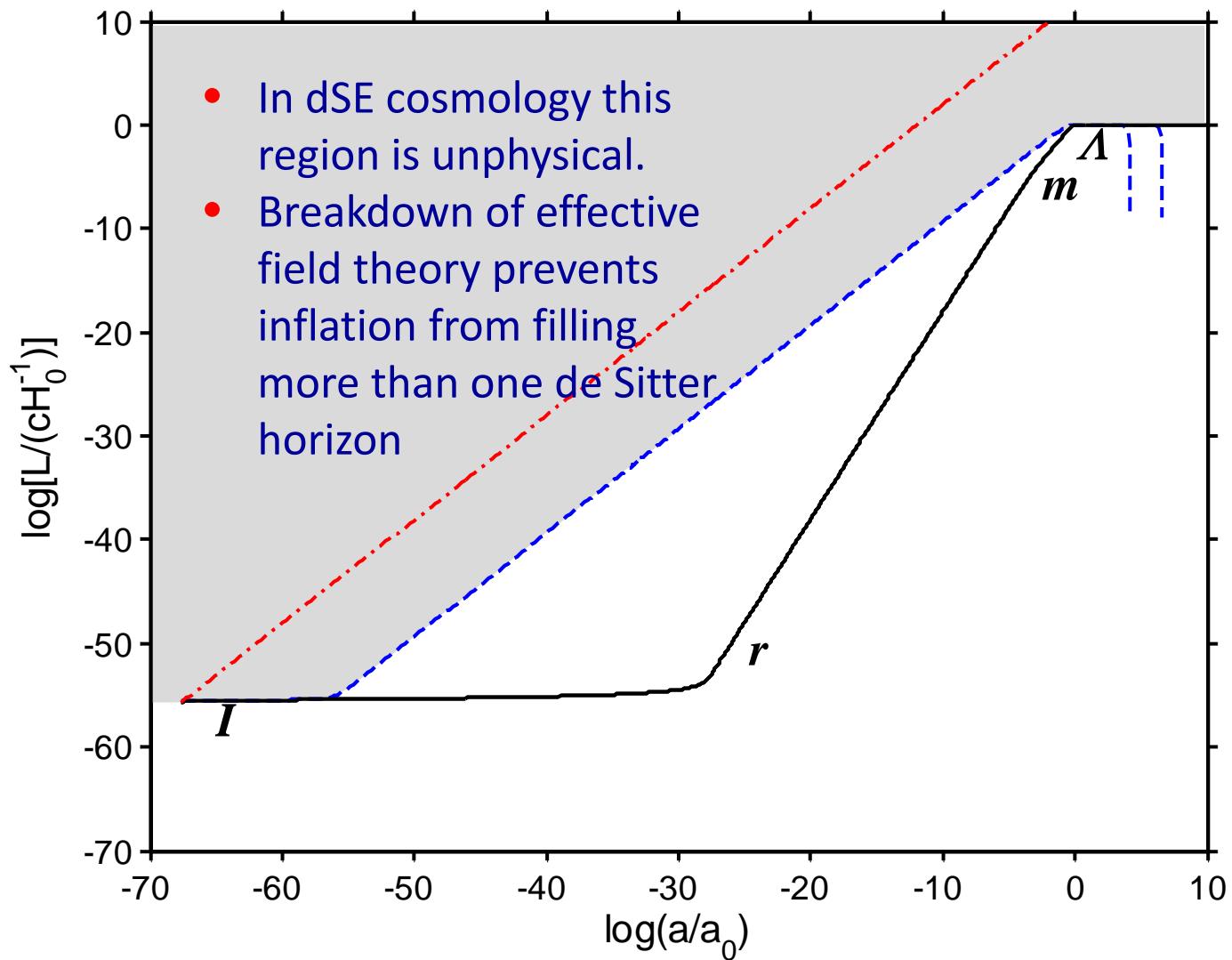
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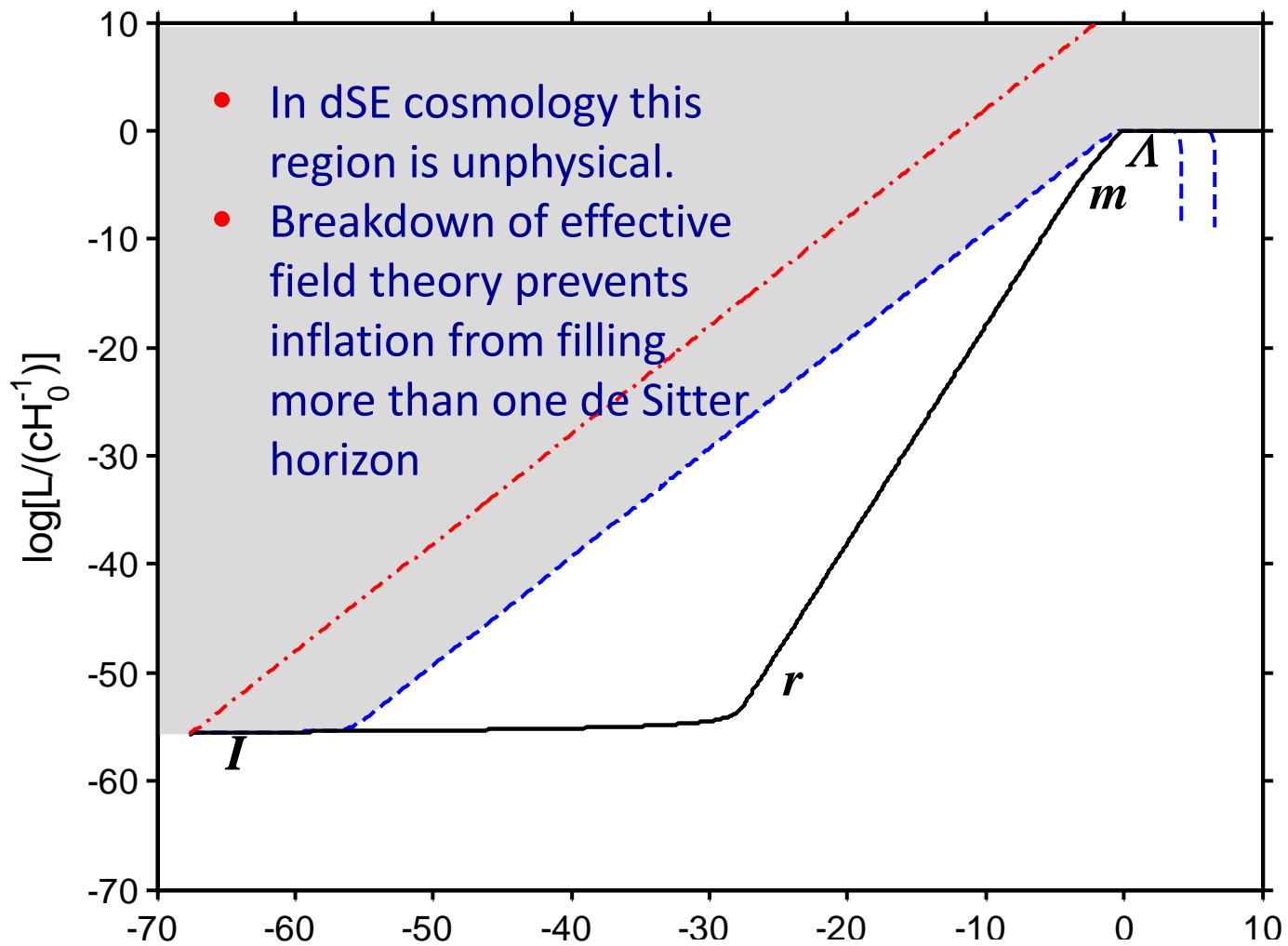




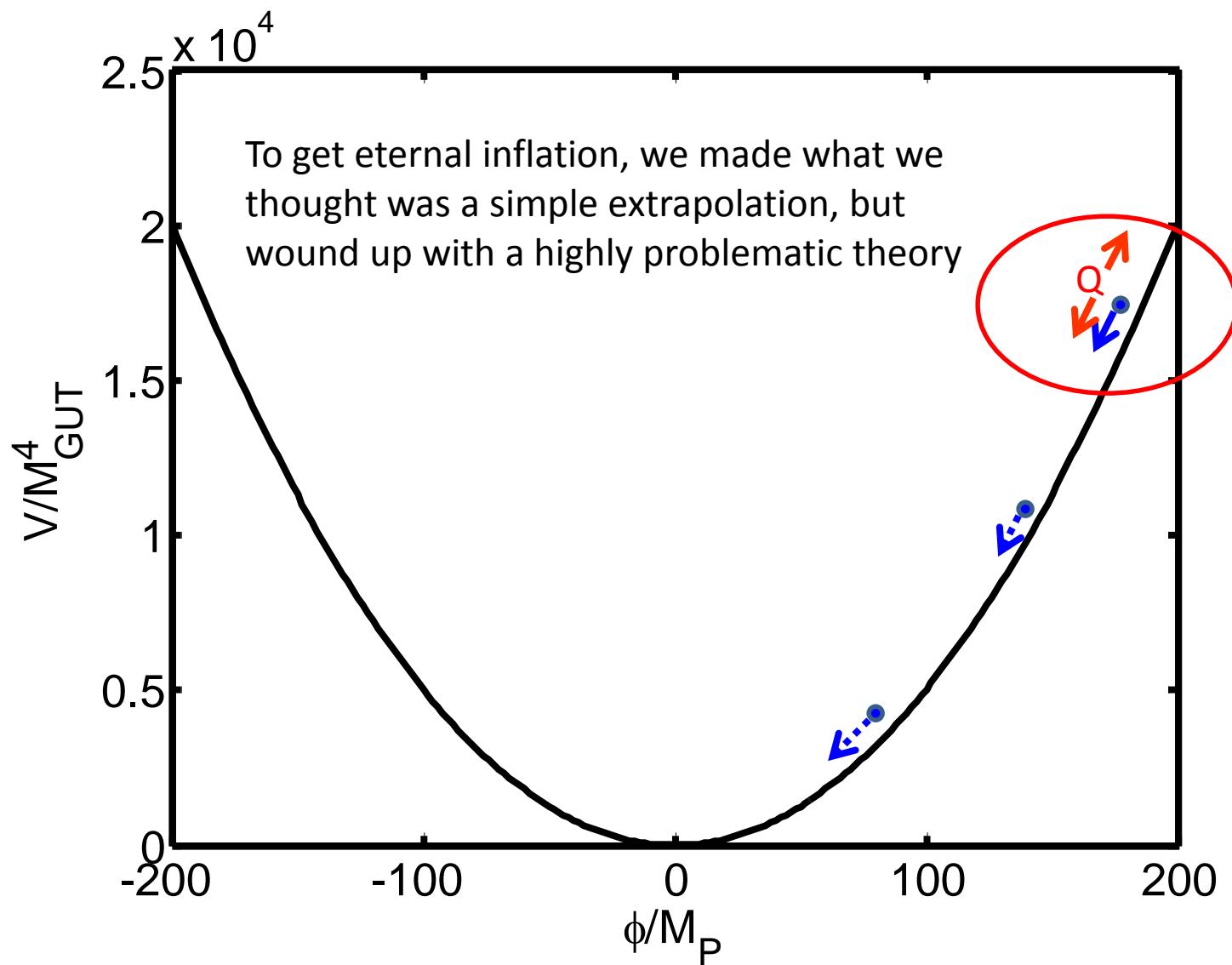


This much inflation fills more than one de Sitter horizon, generating total entropy  $> S_{Max} = S_\Lambda$  and affecting regions beyond the horizon of the observer



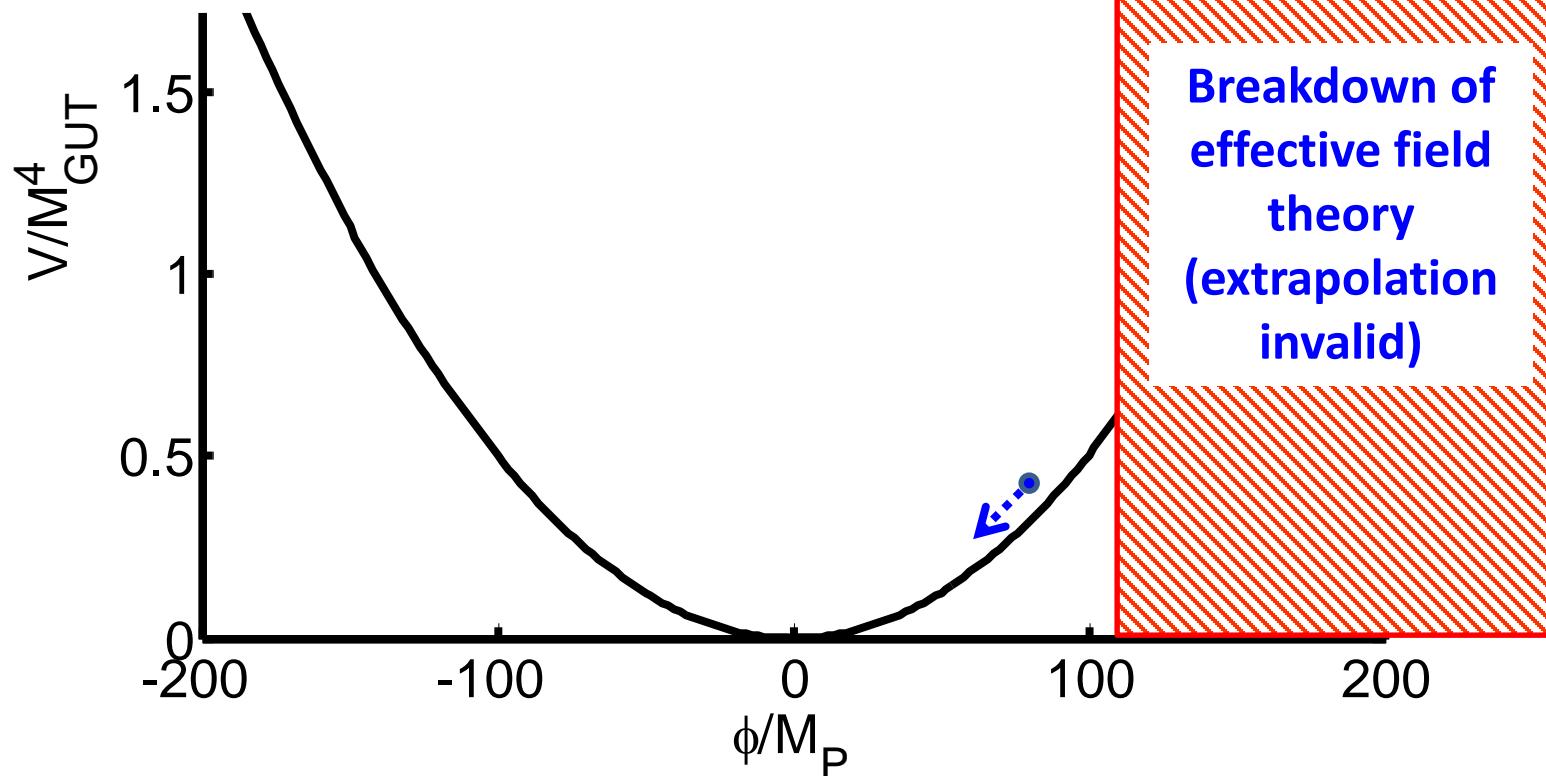


“Equivalent” to Banks-Fischler holographic constraint on number of e-foldings of inflation  
 (D Phillips & AA in prep)



$2.5 \times 10^4$

dSE: The extrapolation that leads to eternal inflation is naïve, in that it neglects the breakdown of effective field theory. dSE uses holographic arguments to estimate this breakdown.



## Fluctuating from dSE to inflation:

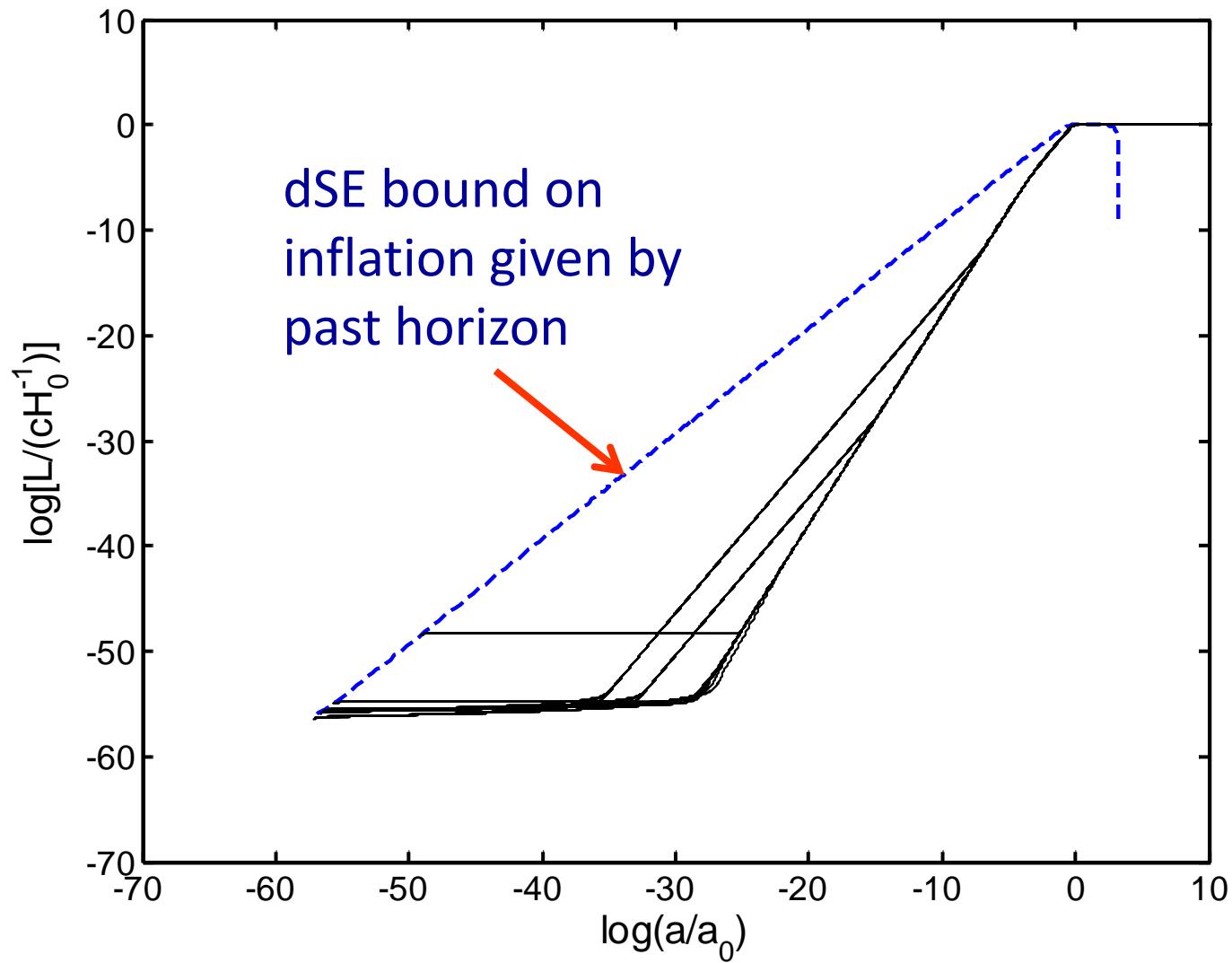
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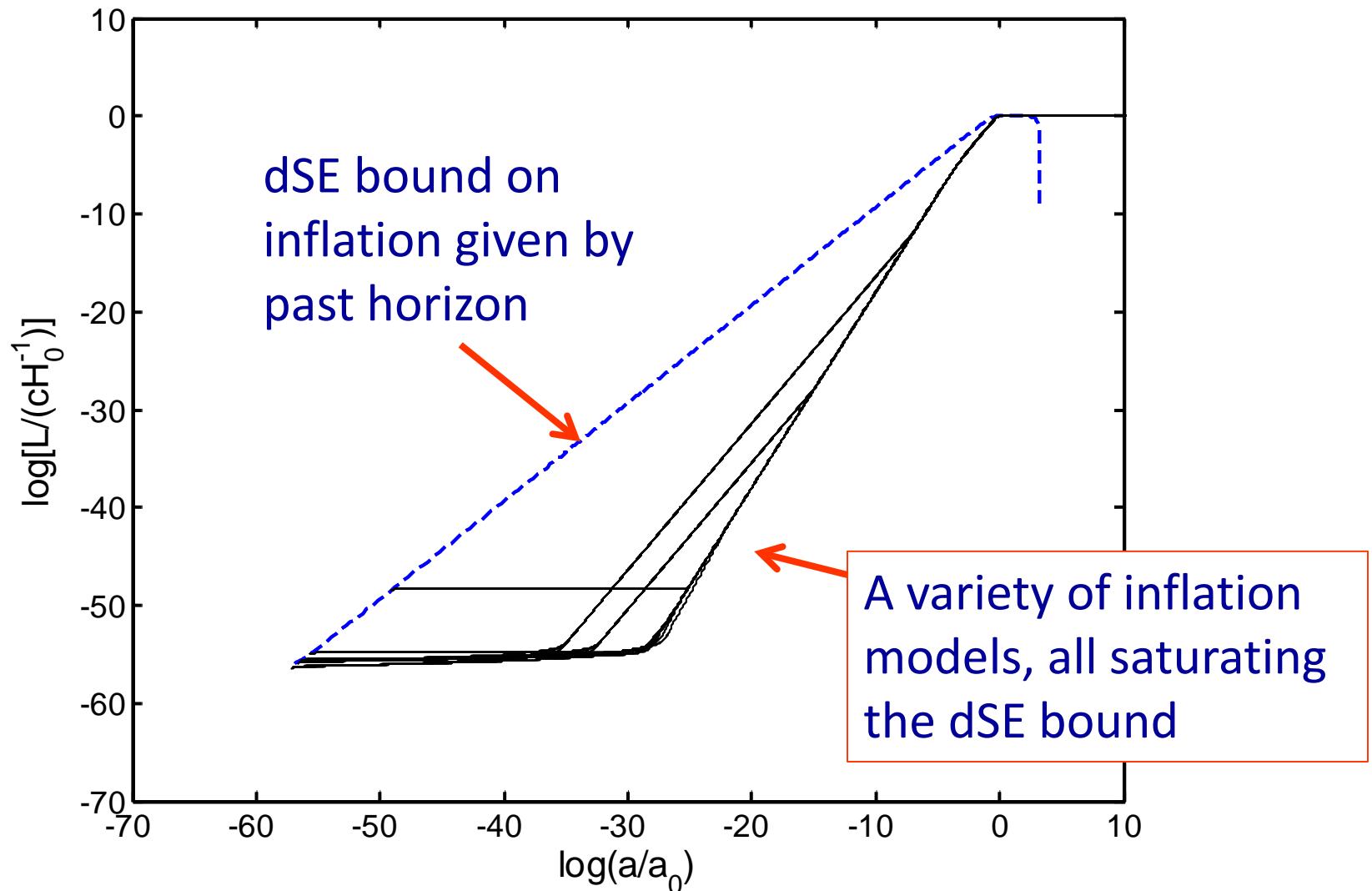
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Large  $\rho_I$  exponentially favored → saturation of dSE bound





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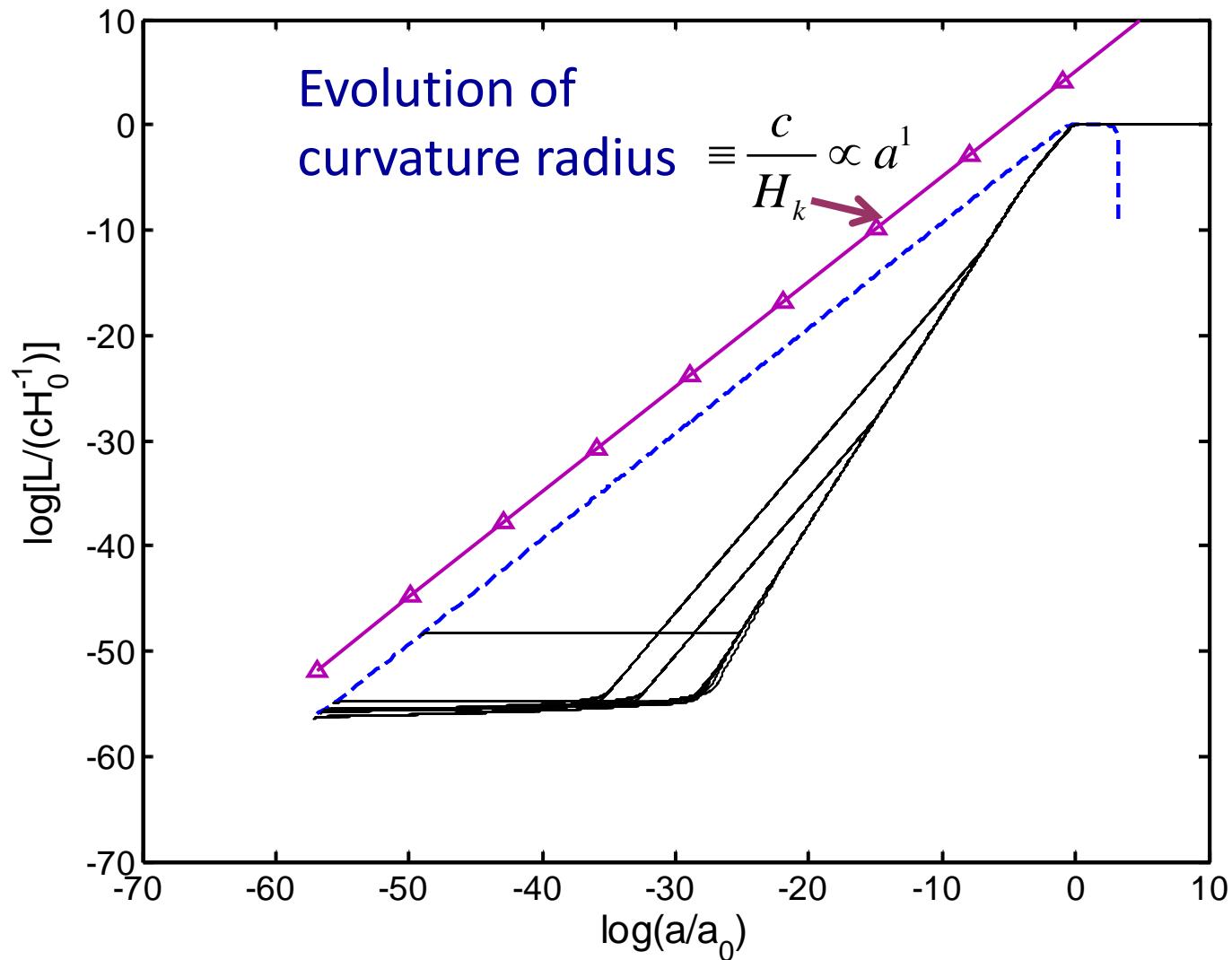


## Friedmann Eqn.

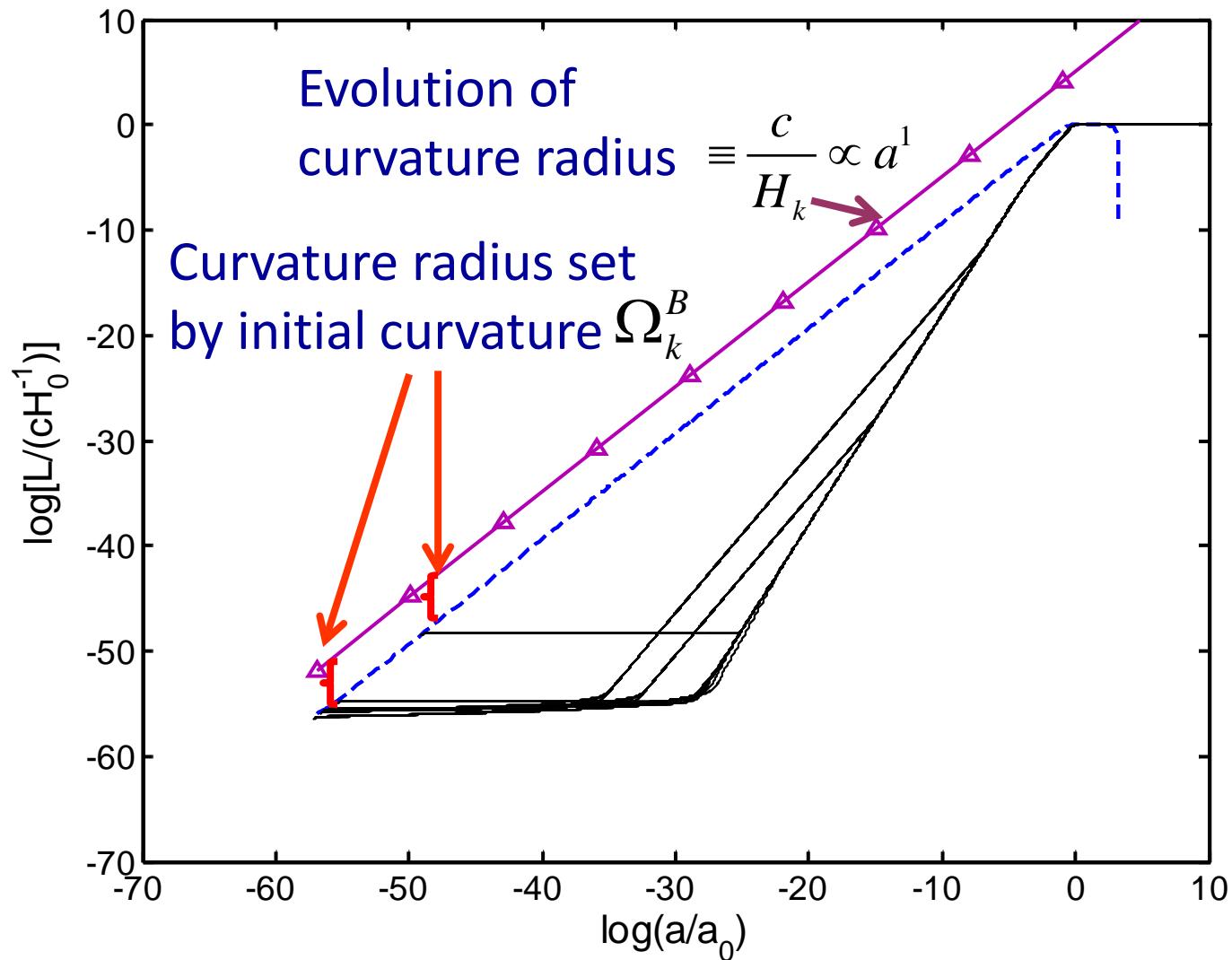
$$H^2 = \left( \frac{\dot{a}}{a} \right)^2 = \frac{8\pi}{3} G (\rho_I + \rho_k + \rho_r + \rho_m + \rho_{DE})$$

$\propto a^{-2}$

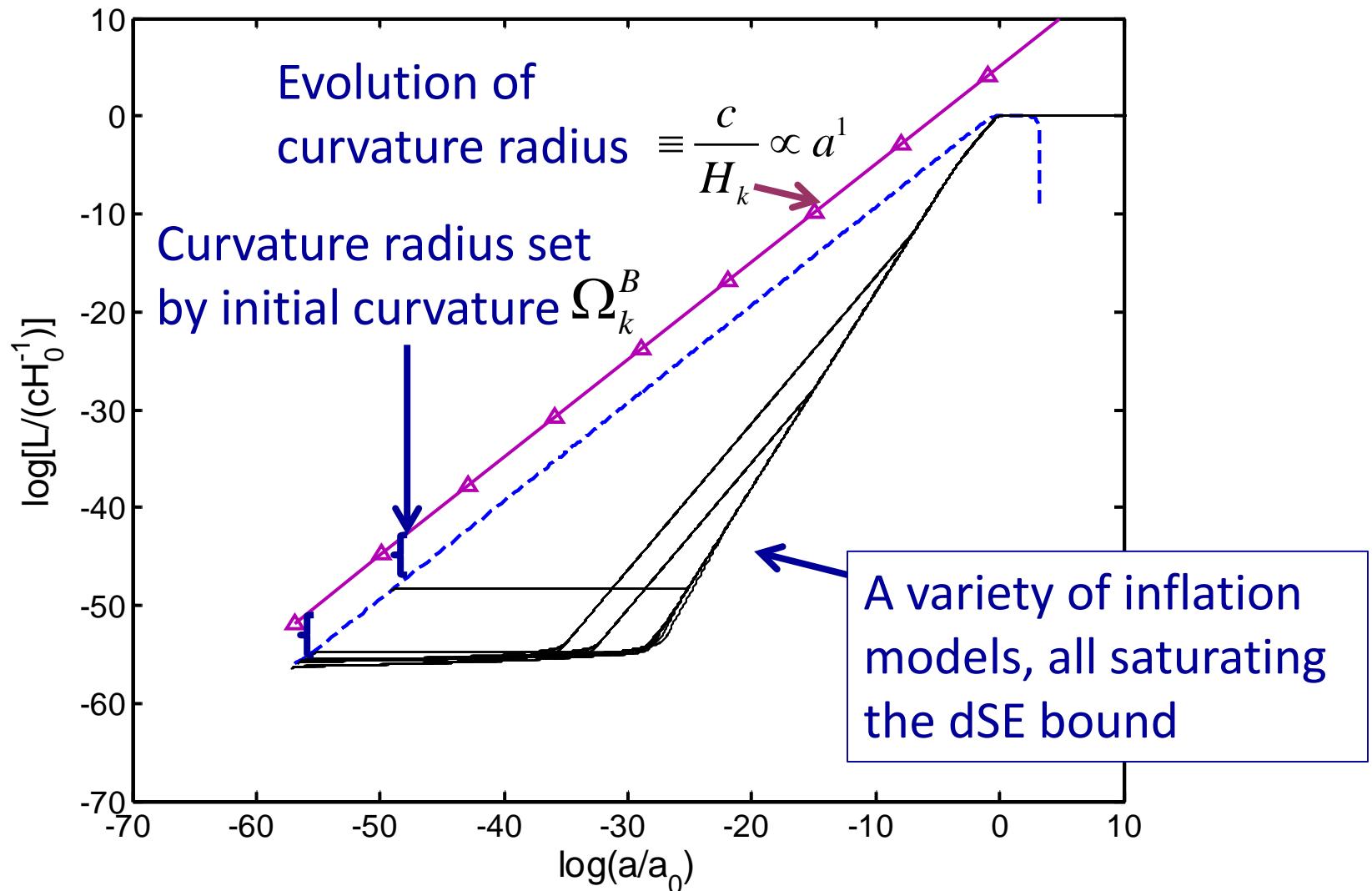
A red rectangular box surrounds the Friedmann equation. A red arrow points from the bottom-left corner of the box upwards and to the right, ending near the term  $\rho_I$ . A vertical red arrow points downwards from the top-right corner of the box, ending near the term  $a^{-2}$ .



Banks & Fischler & Dyson et al.

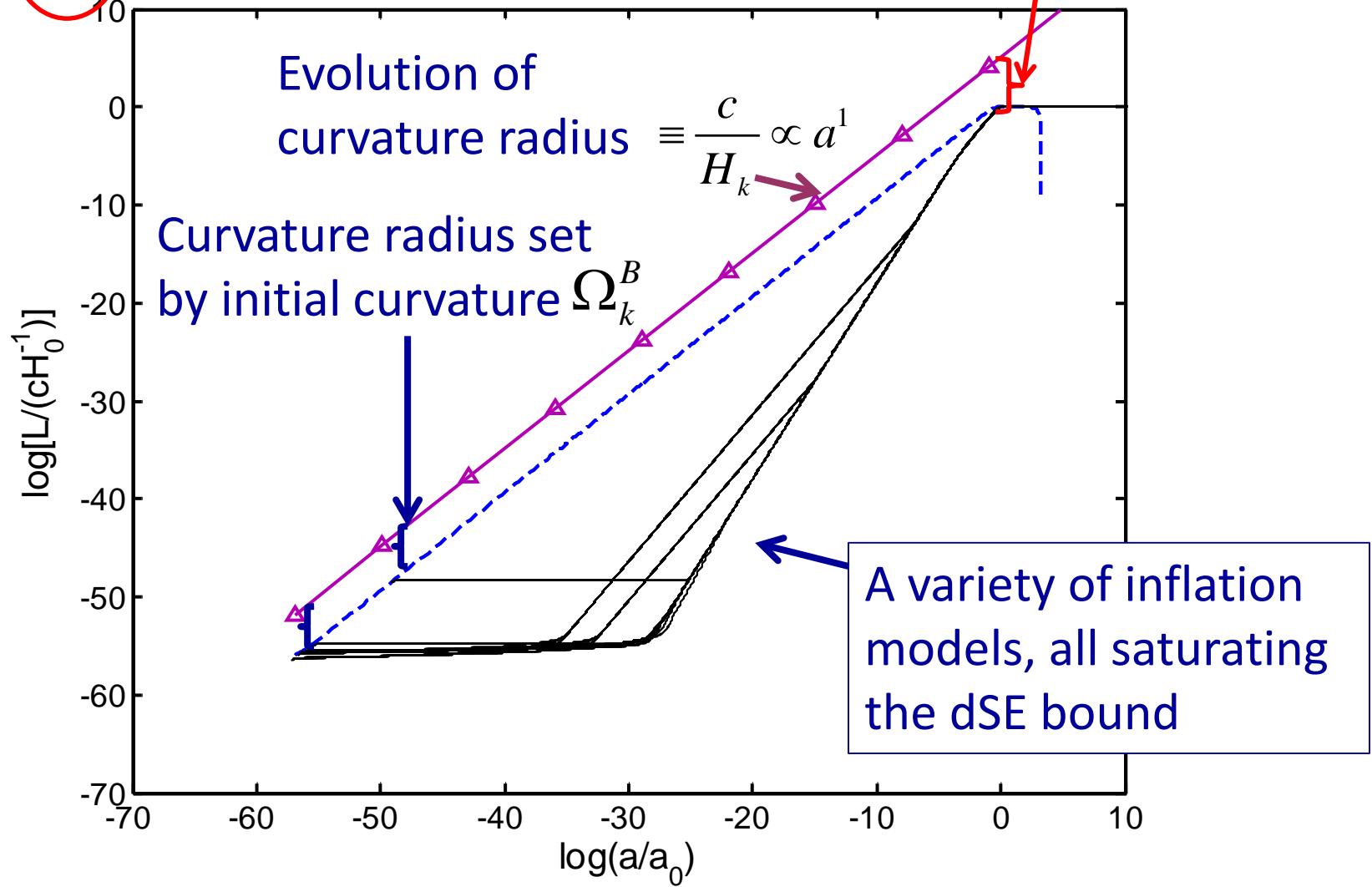


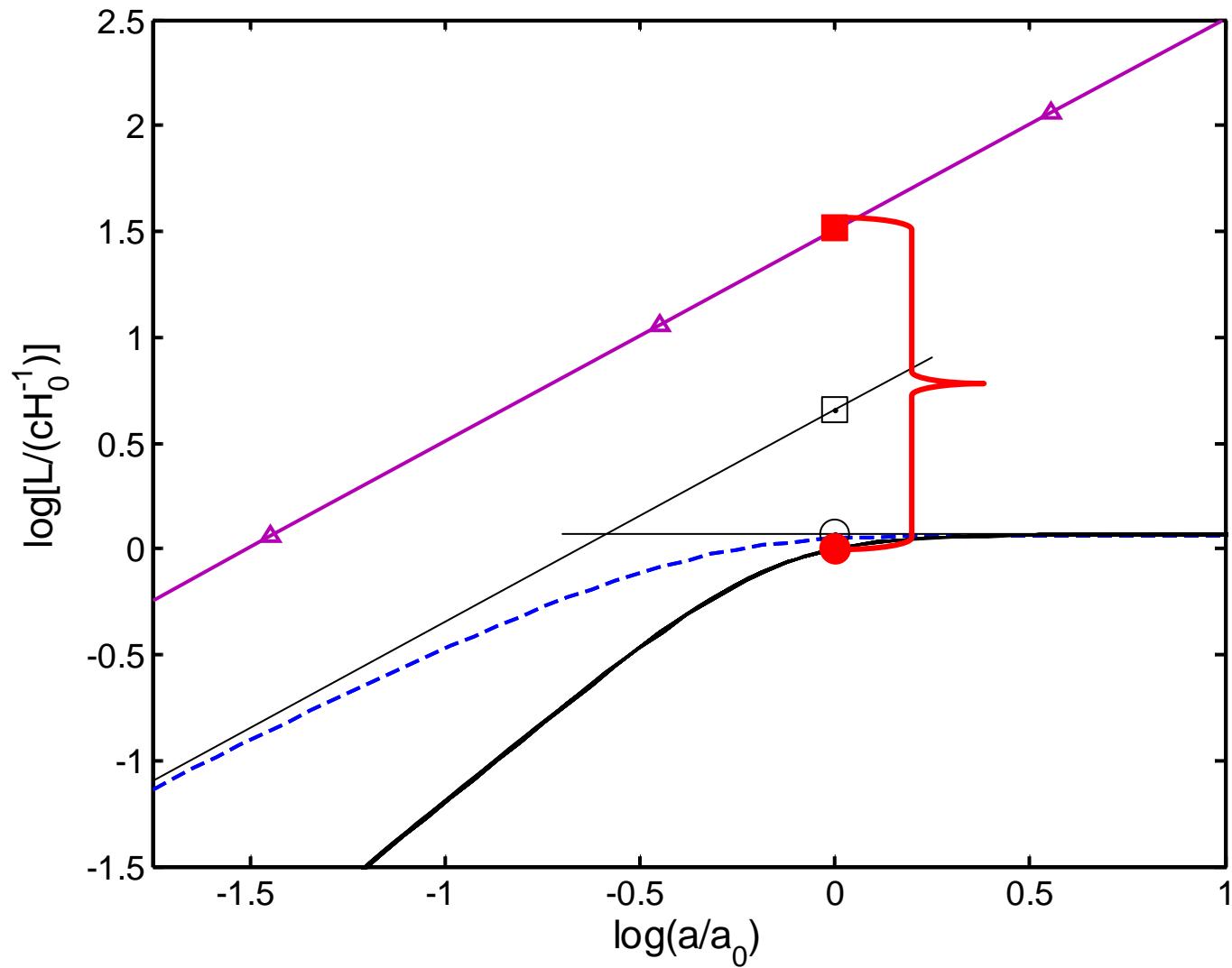
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$$\Omega_k \equiv \frac{\rho_k}{\rho_c} = \left( \frac{H_k}{H_0} \right)^2 \equiv \left( \frac{R_{H_0}}{R_k} \right)^2$$

is given by this gap





AA: arXiv:1104.3315

# dSE Cosmology and cosmic curvature

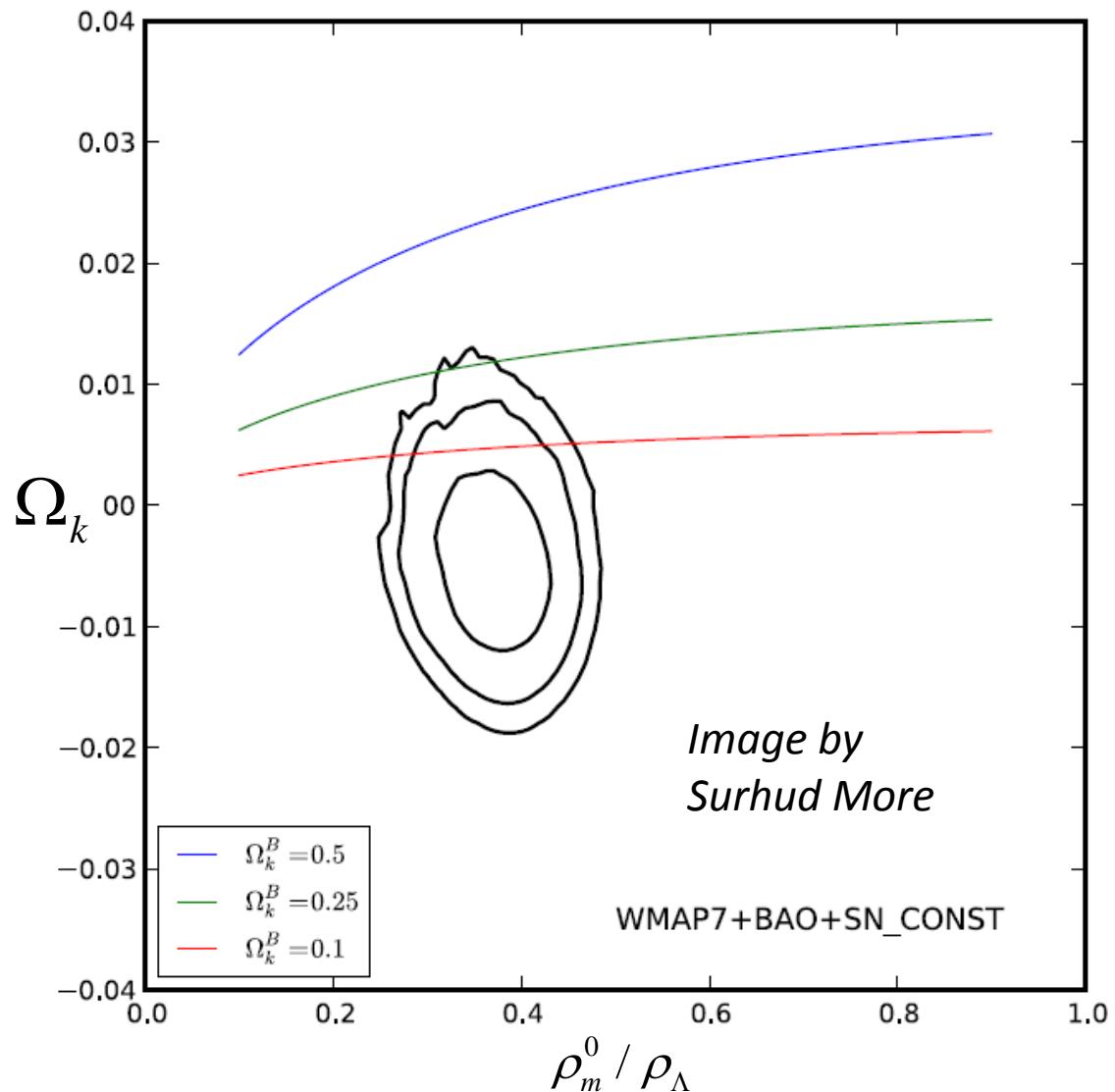
- The Guth-Farhi process starts inflation with an initial curvature set by the curvature of the Guth-Farhi bubble  $\Omega_k^B$
- Inflation dilutes the curvature, but dSE cosmology has a minimal amount of inflation

$$\Omega_k = \frac{1}{g^2} \frac{\Omega_k^B}{\left( \frac{\rho_m^0}{\rho_\Lambda} + \frac{\rho_k^0}{\rho_\Lambda} + 1 \right)}$$

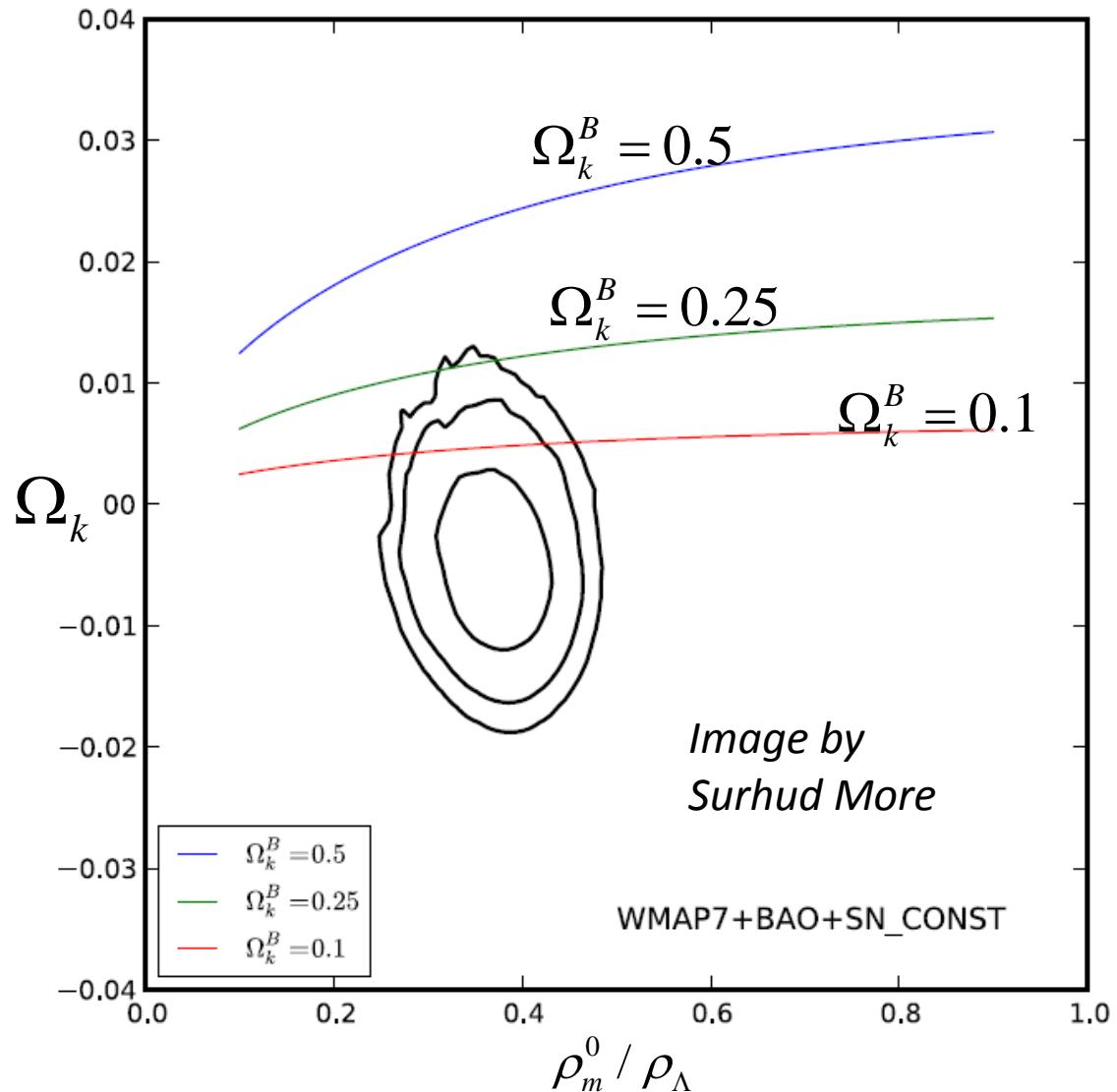
where

$$g \left( \frac{\rho_m^0}{\rho_\Lambda}, \frac{\rho_k^0}{\rho_\Lambda} \right) \equiv \int_0^\infty \frac{dx}{x^2 \sqrt{x^{-3} \frac{\rho_m^0}{\rho_\Lambda} + x^{-2} \frac{\rho_k^0}{\rho_\Lambda} + 1}}$$

- Predicted  $\Omega_k$   
from dSE cosmology is:
- Independent of almost all details of the cosmology
  - Just consistent with current observations
  - Will easily be detected by future observations

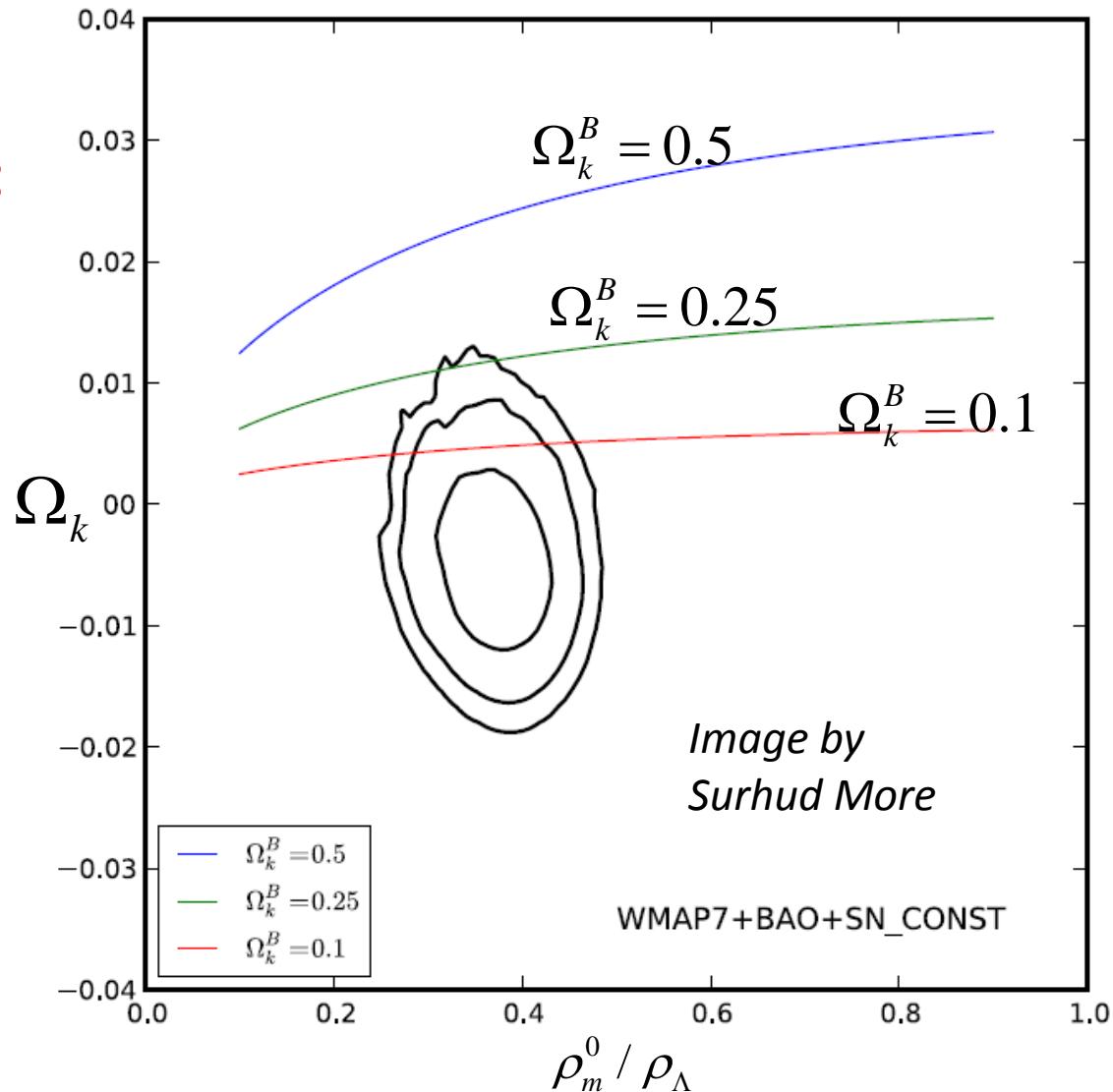


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Work in progress on expected values  
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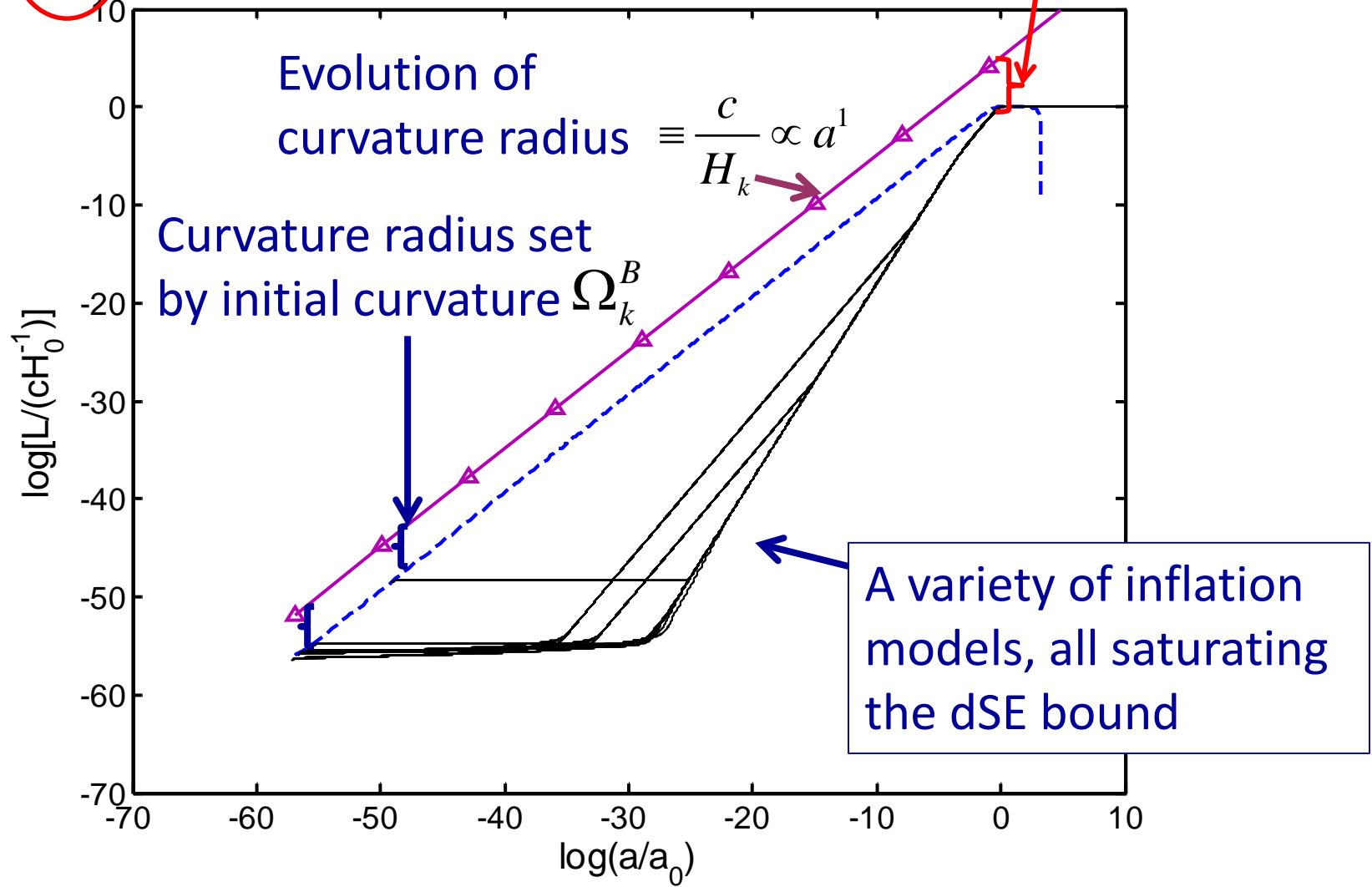
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## Conclusions (Part 4)

- The search for a “big picture” of the Universe that explains why the region we observe should take this form has proven challenging, but has generated exciting ideas.
- We know we can do science with the Universe
- It appears that there is something right about cosmic inflation
- dSE cosmology offers a finite alternative to the extravagant (and problematic) infinities of eternal inflation (plus, no initial conditions problem)
- Predictions of observable levels of cosmic curvature from dSE cosmology will give an important future test

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- dSE cosmology offers a finite alternative to the extravagant (and problematic) infinities of eternal inflation (plus, no initial conditions problem)
- Predictions of observable levels of cosmic curvature from dSE cosmology will give an important future test (also, alternative to bubble start)



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- 3) Classical and quantum probabilities in the multiverse
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End Part 4

Andreas Albrecht; UC Davis  
Les Houches Lectures; July-Aug 2013

# Some parting thoughts

- We are *\*so\** fortunate to be doing cosmology in these times
- Fantastic successes (slow roll inflation & structure)
- Great puzzles and challenges (the Guth dream, dark energy)
- Which ideas to import from everyday physics, which to throw out?
- Which ideas are too radical, which not radical enough?
- Let us all make the most of these amazing opportunities, and navigate these difficult issues as bravely and energetically as possible!